

Single proton knock-out at JLab, MAMI, and MIT-Bates

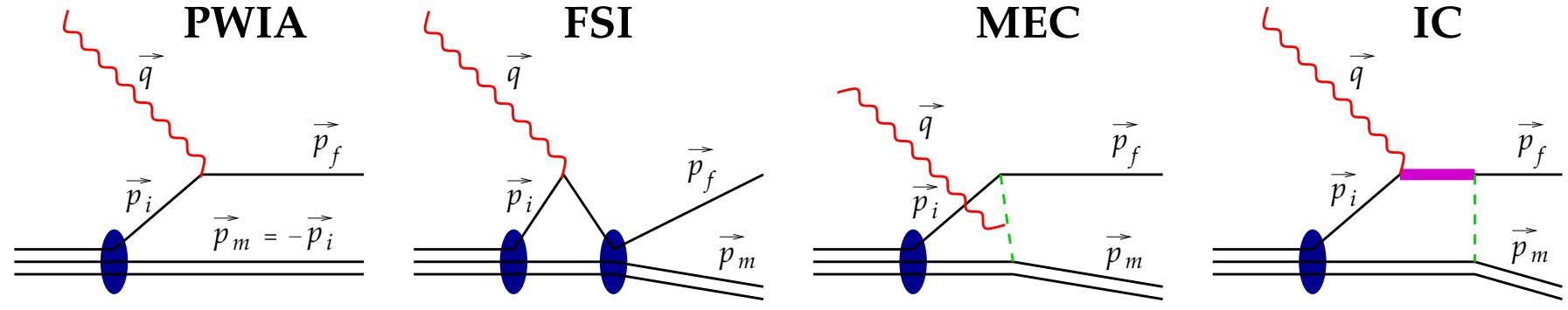
S. Širca, U. of Ljubljana, Slovenia

Milos, Greece, 10 Sep 2007

Topics [the $(e, e'p)$ ‘hierarchy’]

- ▷ Complex nuclei
 - Limits of the impulse approximation
 - Dynamical relativistic effects
 - Nuclear transparency
 - Short-range correlations → J. Watson
 - Connection to neutrino probes
 - Medium modification of FFs → S. Strauch, O. Buss
- ▷ ${}^4\text{He}$, ${}^3\text{He}$
 - Testing ground of few-body theories
 - Benchmark experiments
 - Breakdown of factorization at high p_m
 - 3-body mechanisms
 - ${}^3\text{He}$ ground-state WF components
 - Triple polarization
- ▷ Deuteron
 - The nucleus we don’t know well enough
 - Severe discrepancies even at low E

The A(e, e'N)A-1 alphabet



missing energy $E_m = \omega - T_N - T_{A-1}$

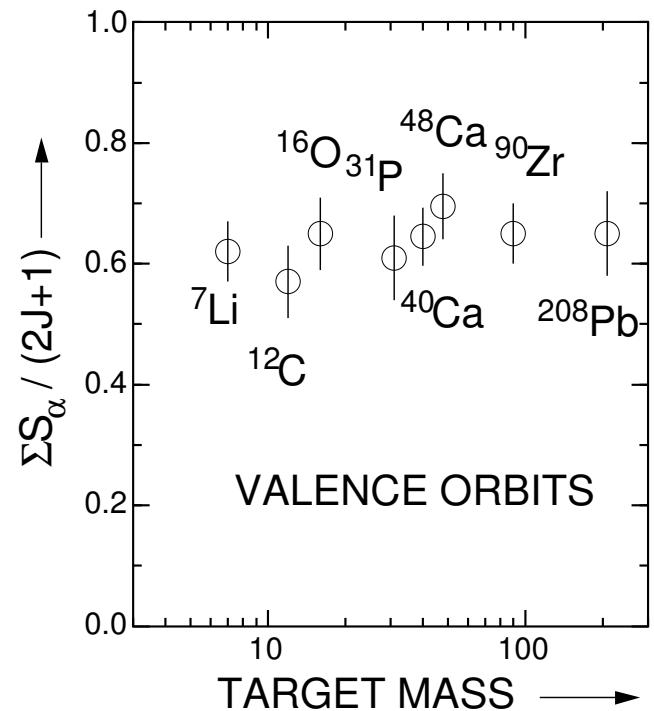
missing momentum $\vec{p}_m = \vec{q} - \vec{p}_f$

PWIA
$$\frac{d^6\sigma}{d\omega d\Omega_e dp_f d\Omega_N} = K\sigma_{eN} S(E_m, \vec{p}_m)$$

$S(E_m, \vec{p}_m) = \sum_{\alpha} S_{\alpha} |\Phi_{\alpha}(-\vec{p}_m)|^2 \delta(E_{\alpha} - E_m)$

DWIA
$$\frac{d^6\sigma}{d\omega d\Omega_e dp_f d\Omega_N} = K\sigma_{eN} D(E_m, \vec{p}_m, \vec{p}_f)$$

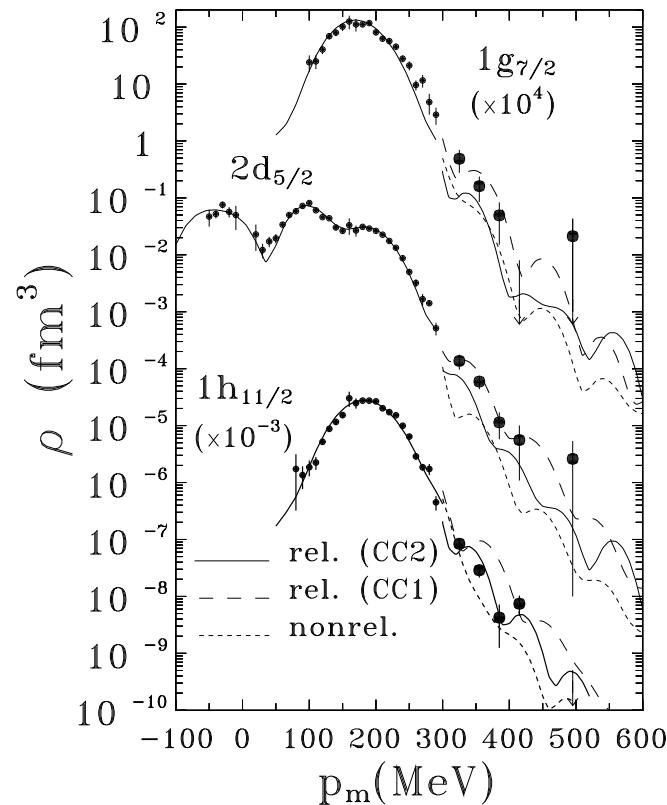
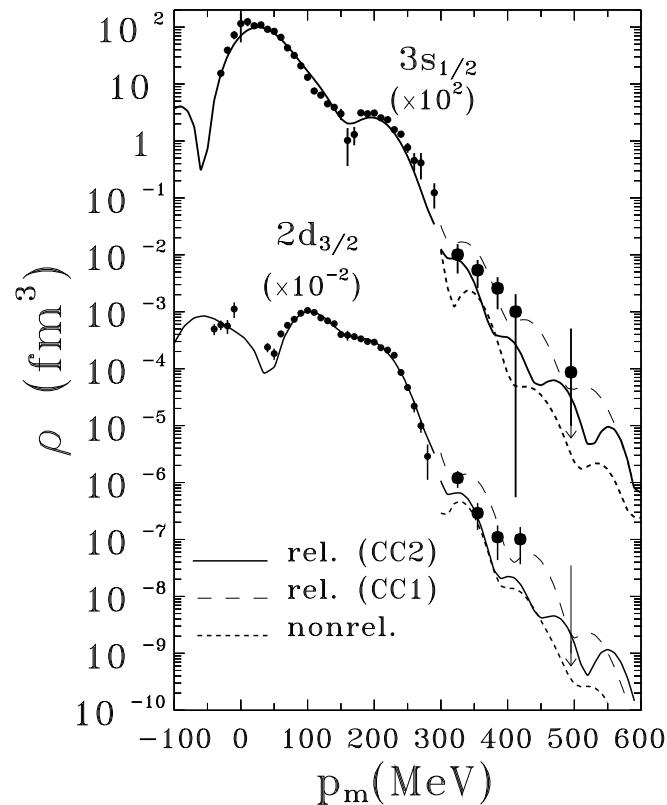
reduced XS $\rho \sim d^6\sigma / K\sigma_{eN}$



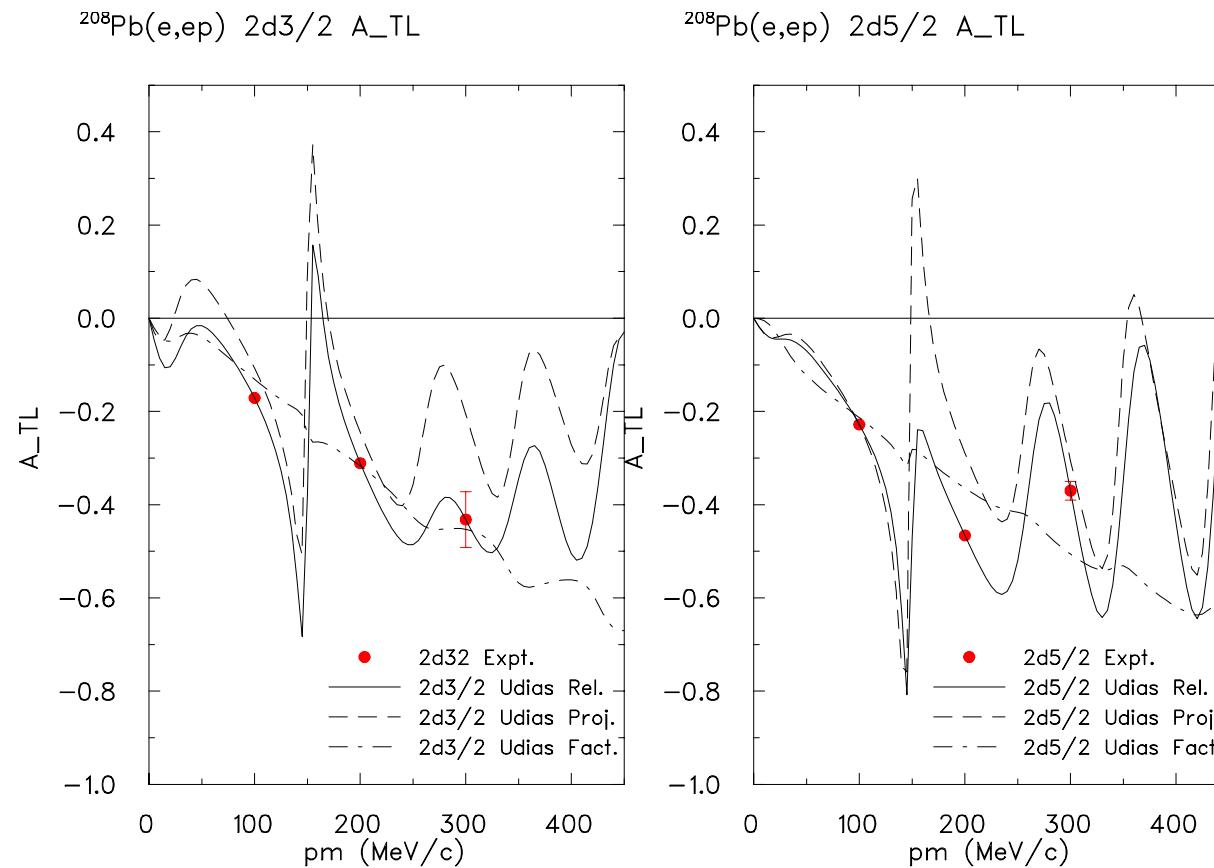
Complex nuclei

$^{208}\text{Pb}(\text{e}, \text{e}'\text{p})^{207}\text{Tl}$

- Successes and limitations of experimental studies and theory
- ^{208}Pb : textbook nucleus for MF approaches
- Exclusive ($\text{e}, \text{e}'\text{p}$) process probes different regions of interior ($0 \leq l \leq 5$)
 $E_m = 0$ ($3s_{1/2}$), 0.351 ($2d_{3/2}$), 1.348 ($1h_{11/2}$), 1.683 ($2d_{5/2}$), 3.470 MeV ($1g_{7/2}$)



- Cross-sections to $p_m \approx 500 \text{ MeV/c}$ and A_{LT} up to $p_m \approx 300 \text{ MeV/c}$
- First time in true QE kinematics at fixed $(\omega, \vec{q}) = (430, 1000)$
- Origin of excess strength at high p_m — LRC or relativity?
- Compare spectroscopic factors to those at lower Q^2



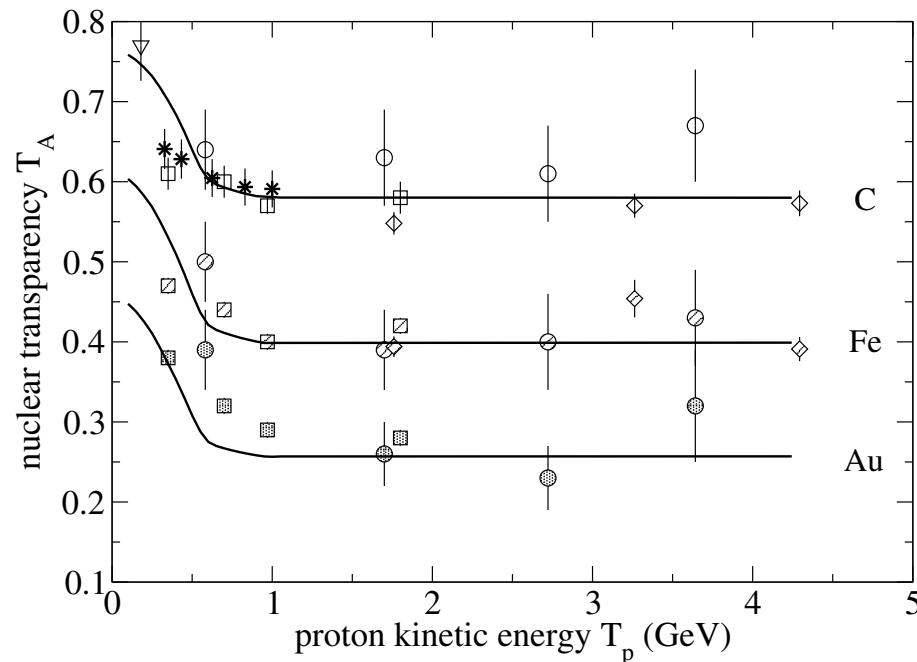
- (Final-state) reduction of proton flux in $(e, e'p)$ at high T_p

$$T_A(Q^2) = \frac{1}{\varepsilon^{\text{SRC}}(A)} \frac{\int_V d^3 p_m dE_m N^{\text{exp}}(E_m, \vec{p}_m)}{\int_V d^3 p_m dE_m N^{\text{calc}}(E_m, \vec{p}_m)}$$

$$N^{\text{calc}} = N^{\text{PWIA}}, \quad \varepsilon^{\text{SRC}}(A) \neq 1$$

$$\text{better: } N^{\text{calc}} = (1 - x)N^{\text{IPSM}} + xN^{\text{SRC}}, \quad \varepsilon^{\text{SRC}}(A) \equiv 1$$

- Validity of optical-potential vs. Glauber approaches for FSI



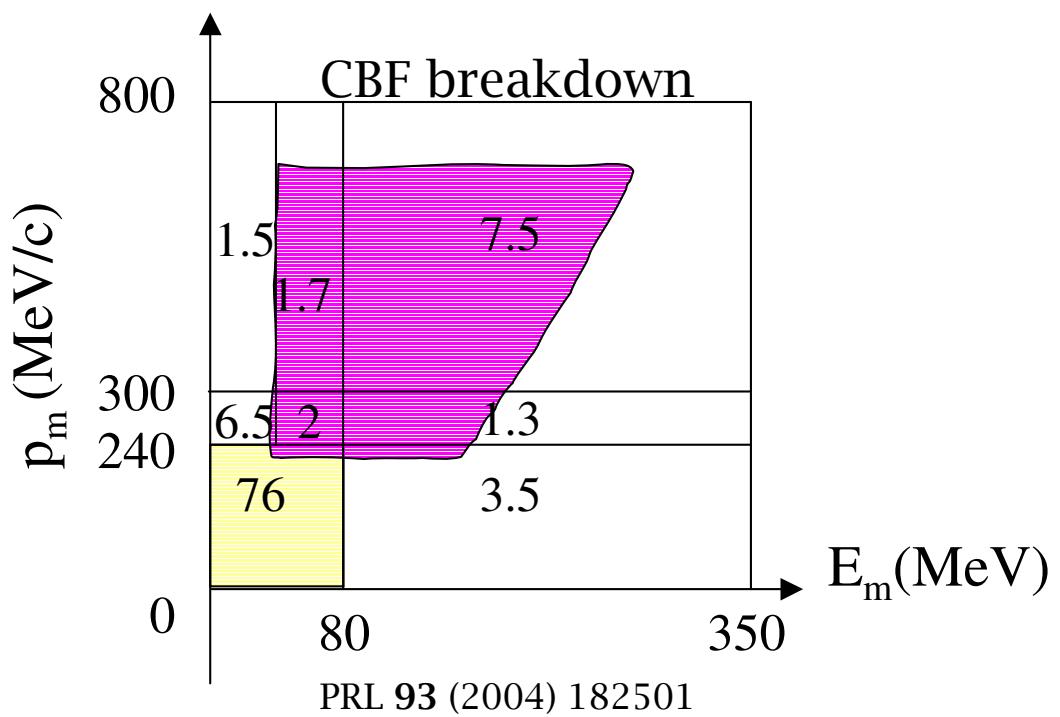
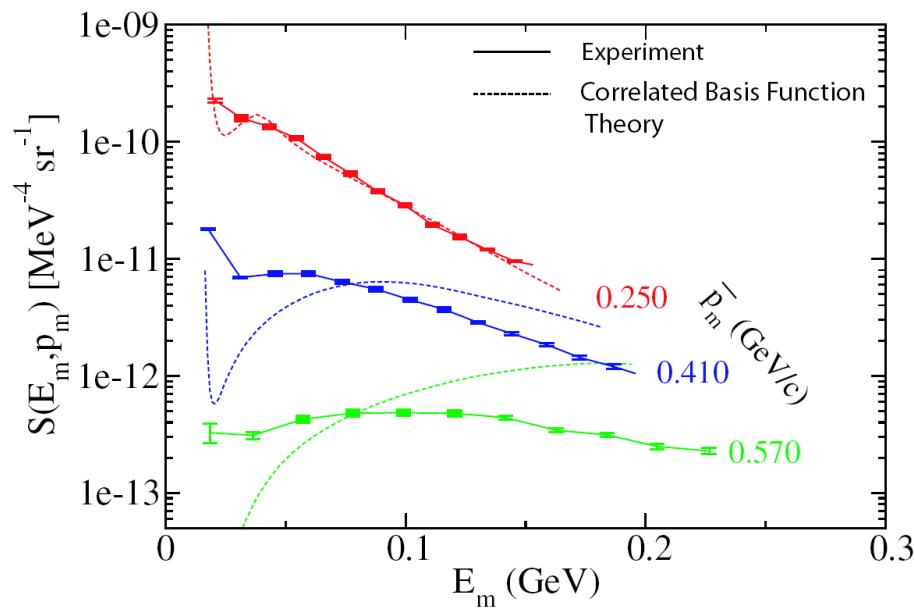
PRC 72 (2005) 054602

PRC 68 (2003) 064603

Correlated strength in $S(E_m, p_m)$

(MODEL) SEARCH FOR SRC

- correlations depopulate IP strength at low E, k ($= E_m, p_m$ in PWIA)
- but also directly identifiable at high E *and* k , quasi- \parallel kinematics best bet
- $\sim 80\%$ of IP motion (CBF, NPA 505 (1989) 267)
- correlated strength 0.61 ± 0.06 (exp), 0.64 (CBF) in measured region
(in terms of # protons)
 1.32 (CBF) total

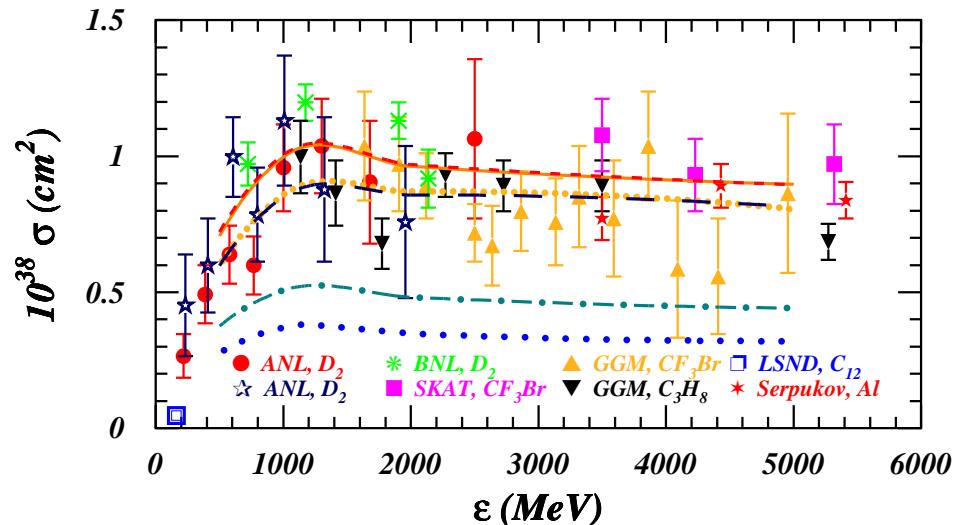
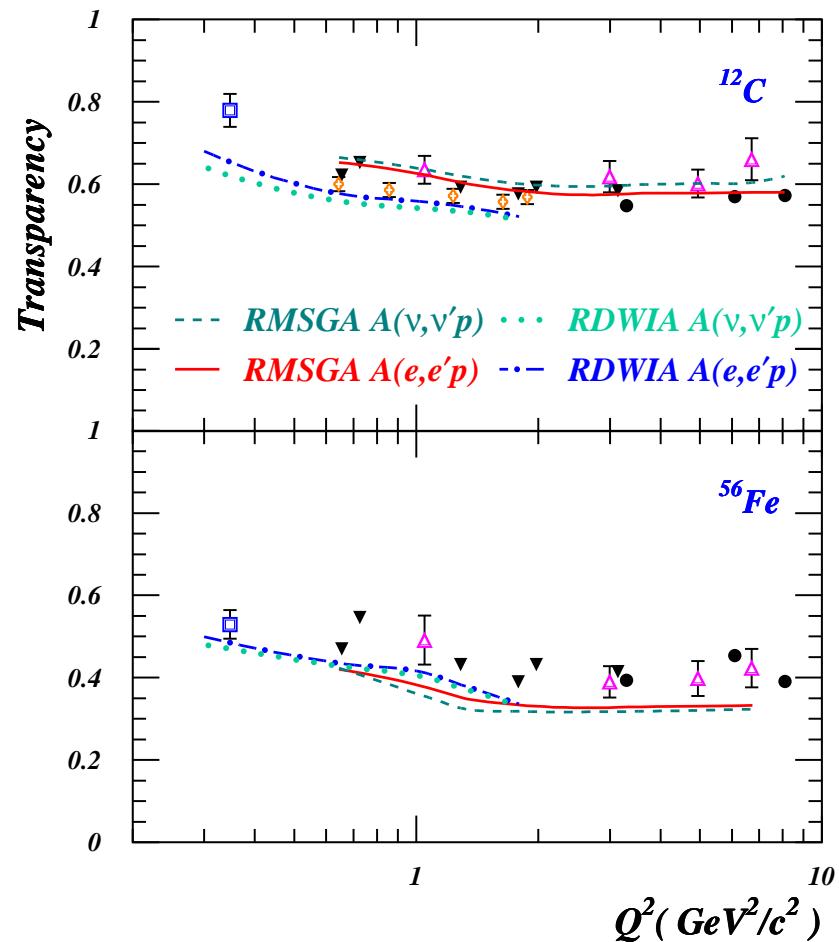


- minor issue: measured and calculated $S(E, k)$ differ

From $A(e, e'p)$ to $A(\nu, \nu'p)$

MADRID-SEVILLE + GHENT

- RDWIA: complex optical potentials for p – A FSI
- RMSGA: Glauber multiple-scattering extension of eikonal + frozen approx
- goal: obtain FSI estimates for $A(\nu, \nu'p)$ from transparency in $A(e, e'p)$



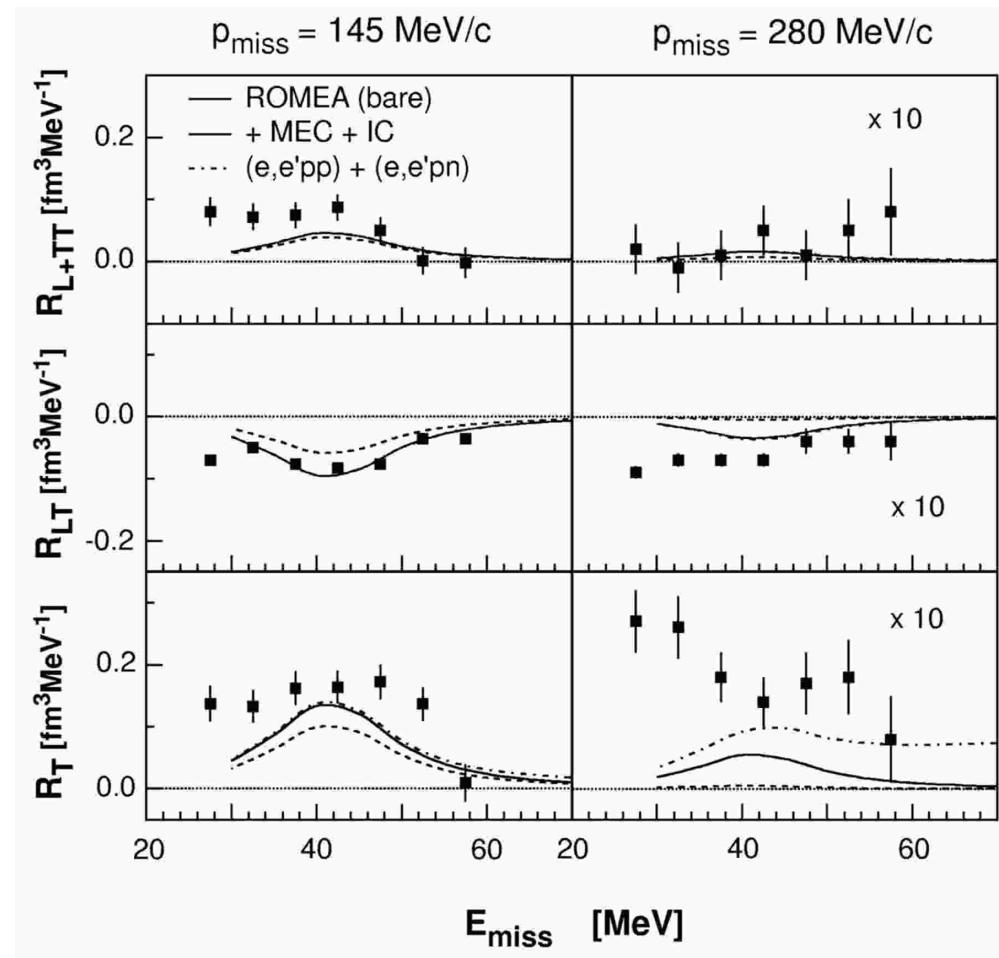
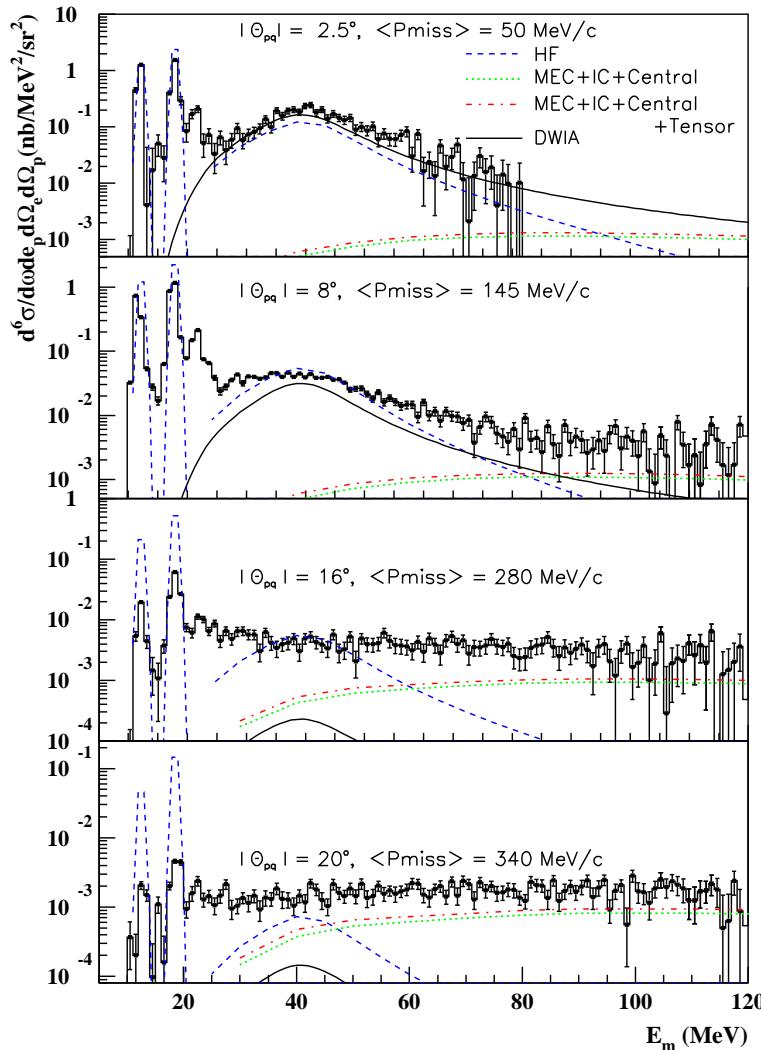
- T independent of leptonic probe
- RDWIA, RMSGA, RPWIA $\times T$ all agree
- $\sim 50\%$ of measured (ν_μ, μ^-) strength is single-nucleon knockout
- π production (Δ): work in progress

PRC 73 (2006) 024607

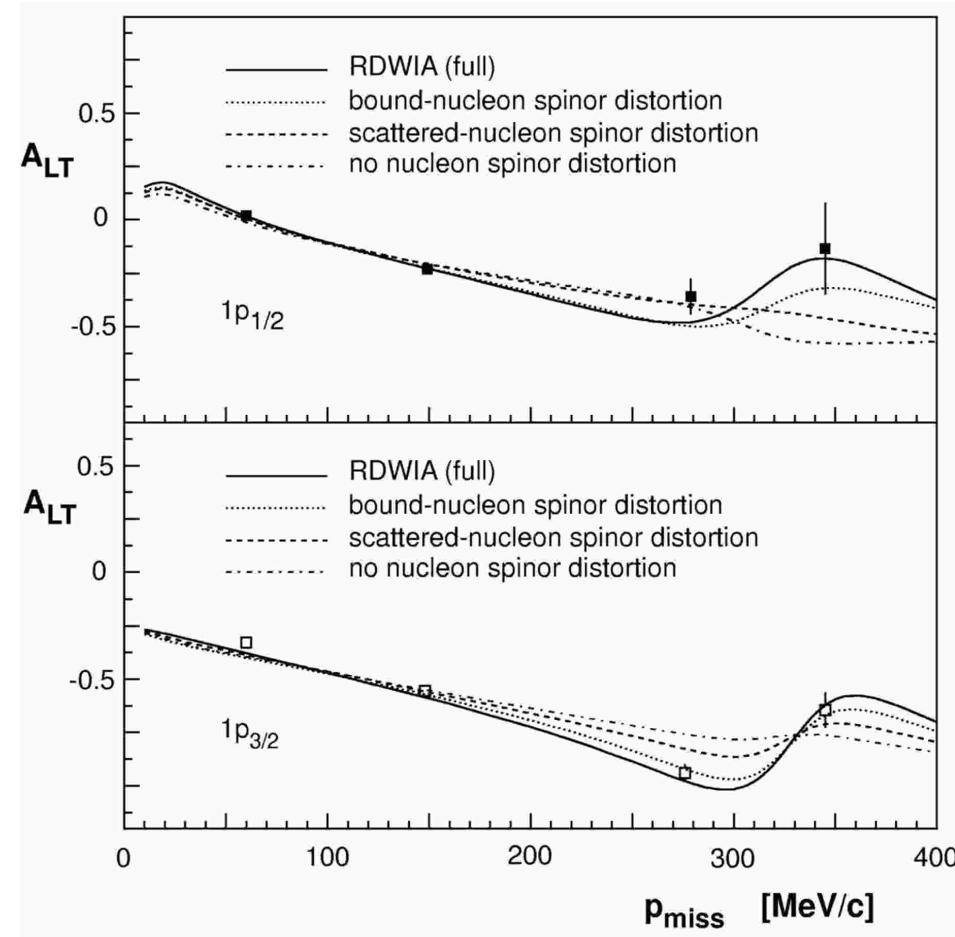
$^{16}\text{O}(\text{e}, \text{e}'\text{p})\text{n}$: continuum

JLAB E89-003

- $d^6\sigma, R_T, R_{L+TT}, R_{LT}$ measured for $25 \leq E_m \leq 120 \text{ MeV}$ and $p_m \leq 340 \text{ MeV}/c$

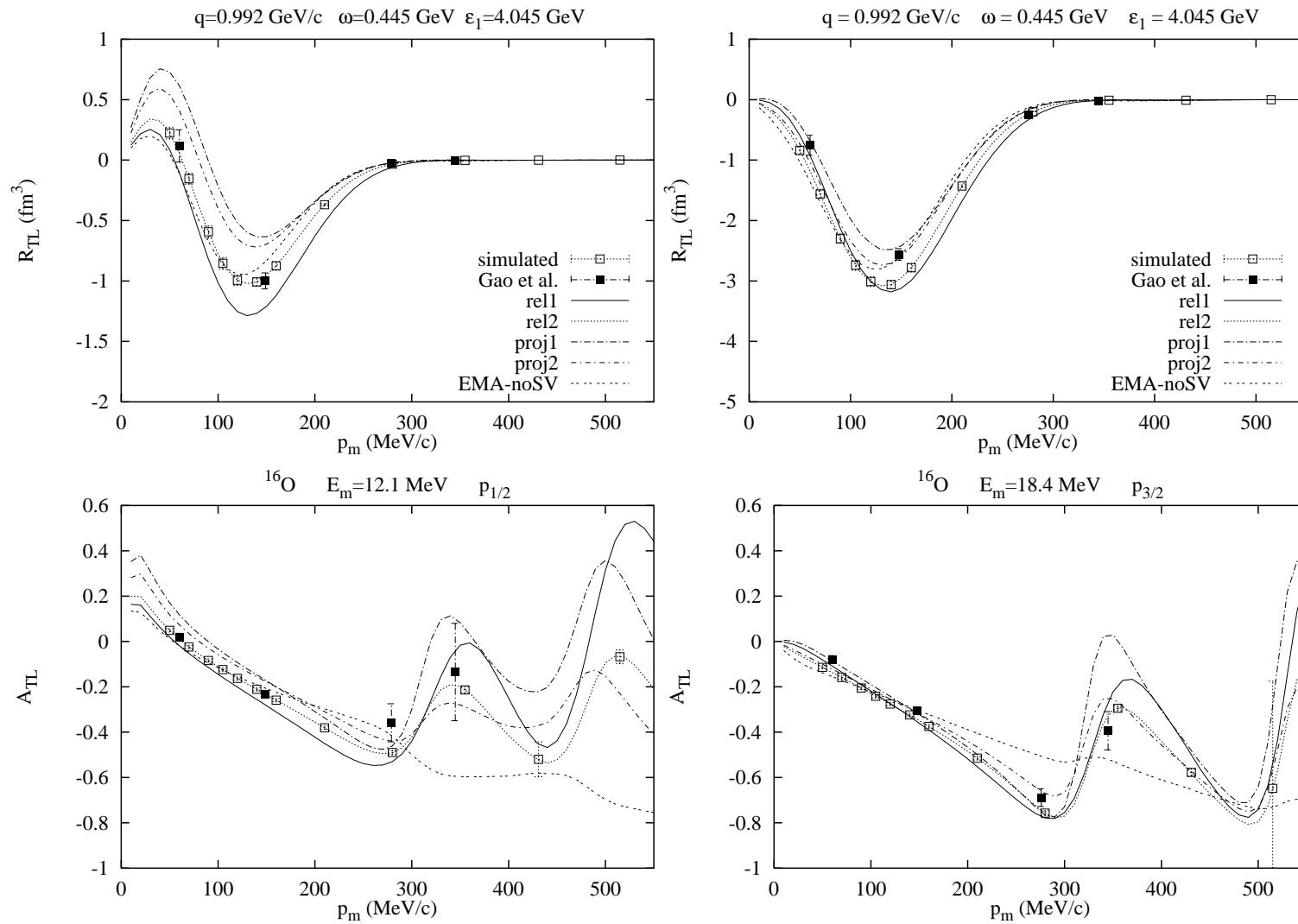


PRC 70 (2004) 034606, PRL 86 (2001) 5670



Enhancement of lower components of bound (and scattered) Dirac spinors
crucial in description of A_{LT} , R_{LT}

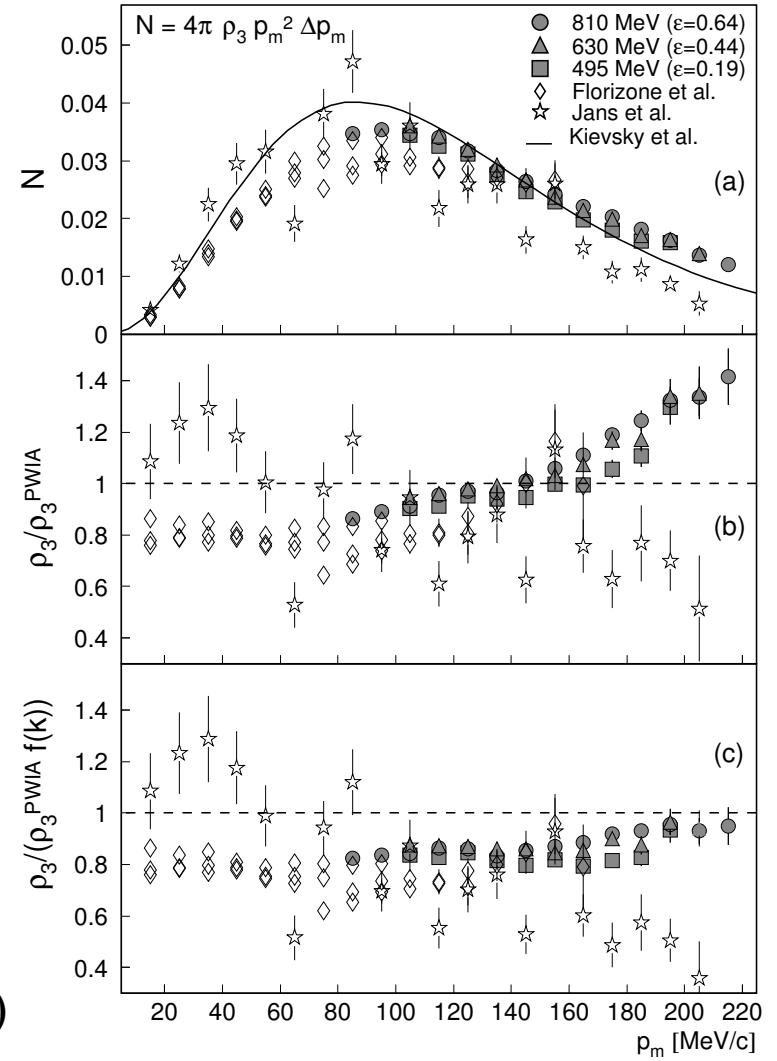
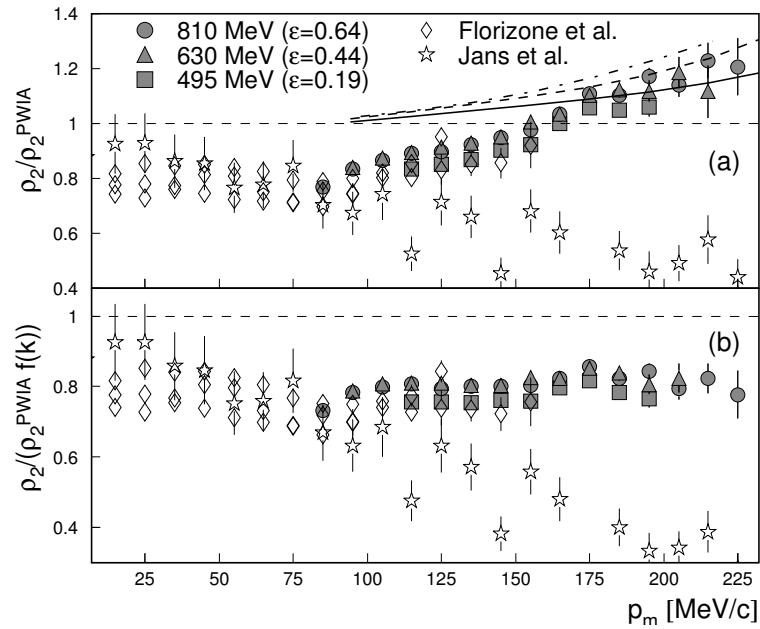
PRC **70** (2004) 034606, PRL **84** (2000) 3265



... work in progress

${}^3\text{He}(e, e'p)\text{pn}$ and ${}^3\text{He}(e, e'p)\text{d}$ below QE peak

MAMI/A1



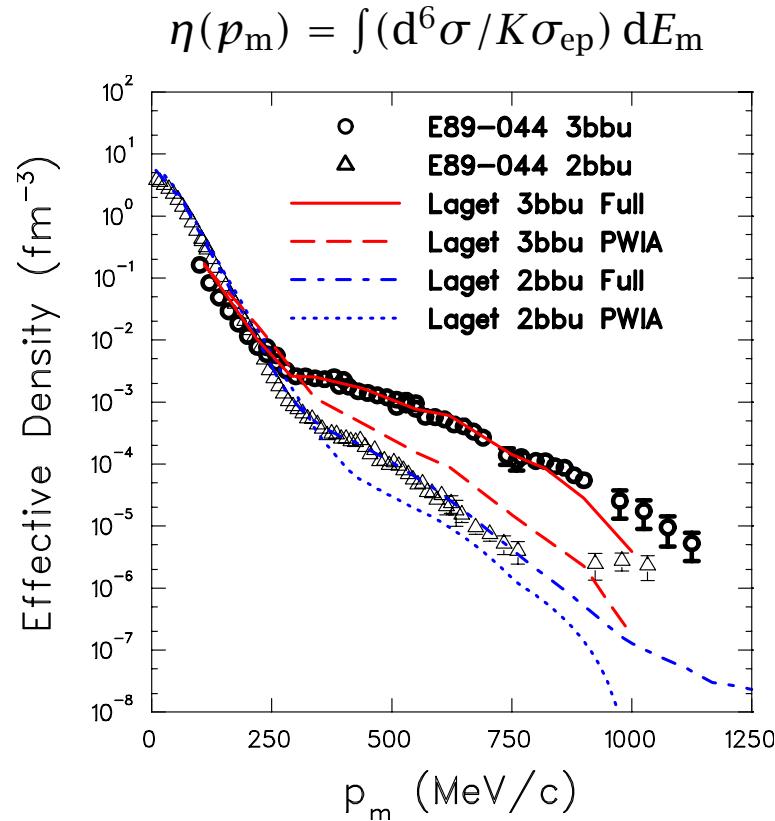
- **low p_m :** data $\approx 20\%$ below full Faddeev ($f(k)$) roughly accounts for FSI + MEC)
... now improving
- **high p_m :** \approx correct p_m dependence
- (?) relativistic dynamics ($q \approx 700 \text{ MeV}/c$)

PRL 93 (2004) 132301

$^3\text{He}(e, e'p)\text{pn}$ on the QE peak

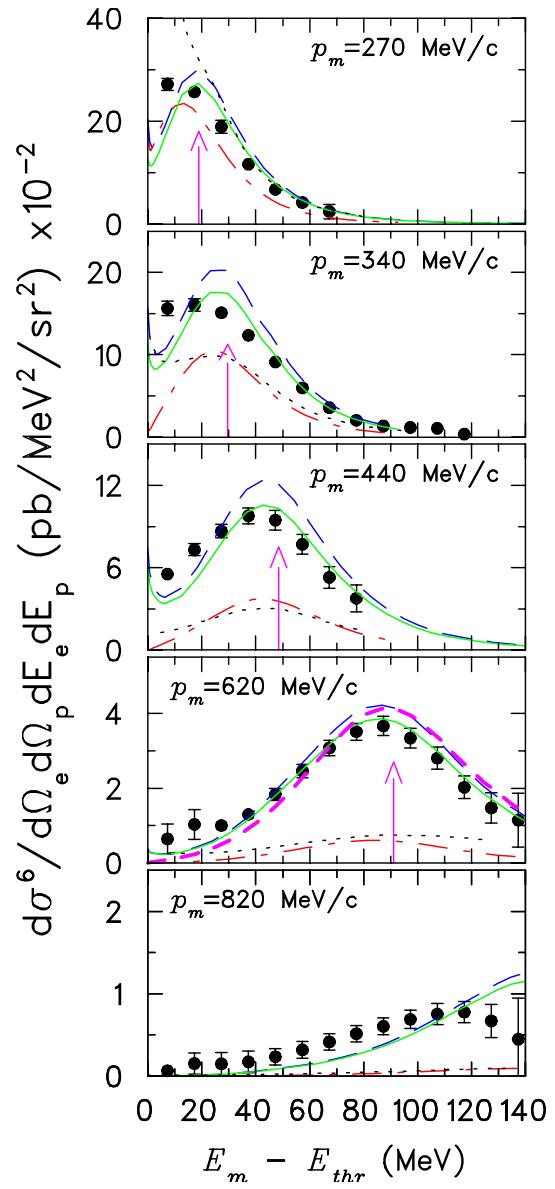
E89-044

- fixed $q = 1500 \text{ MeV}/c$, $\omega = 840 \text{ MeV}$
- p_m up to $1 \text{ GeV}/c$, E_m up to π threshold
- benchmark data (both channels)



- role of correlations (3bbu)

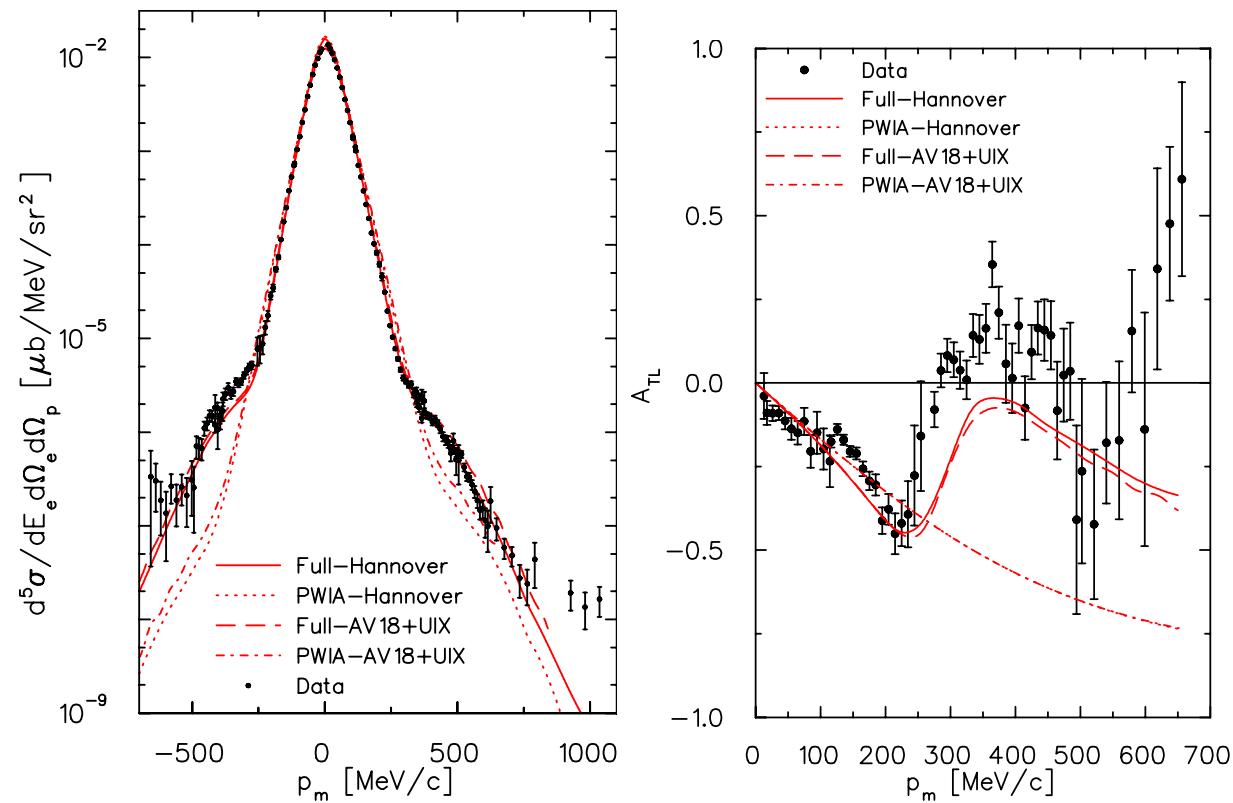
PRL 94 (2005) 082305



$^3\text{He}(e, e' p)^2\text{H}$ on the QE peak

E89-044

- [same e^- kinematics] strong FSI for $150 \leq p_m \leq 750 \text{ MeV}/c$
- large discrepancies wrt theory near $p_m = 1000 \text{ MeV}/c$
- A_{LT} shows breakdown of factorization (IF of PWIA & rescattering amps)

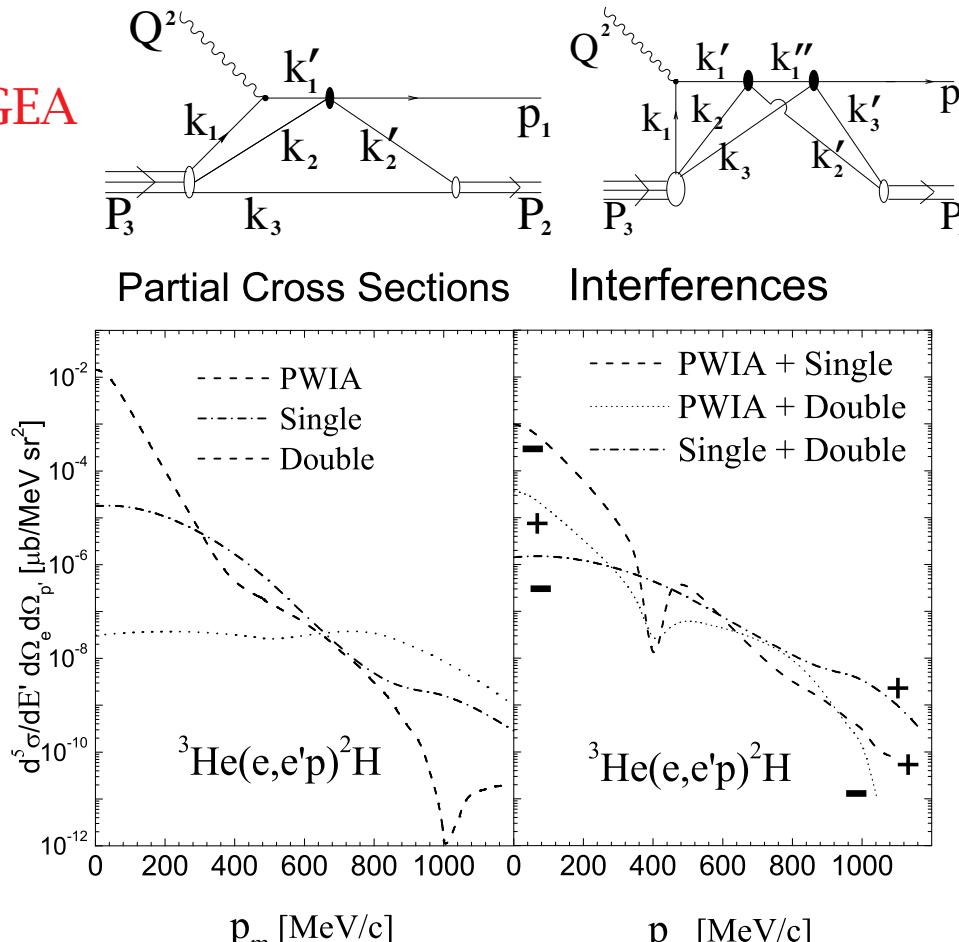
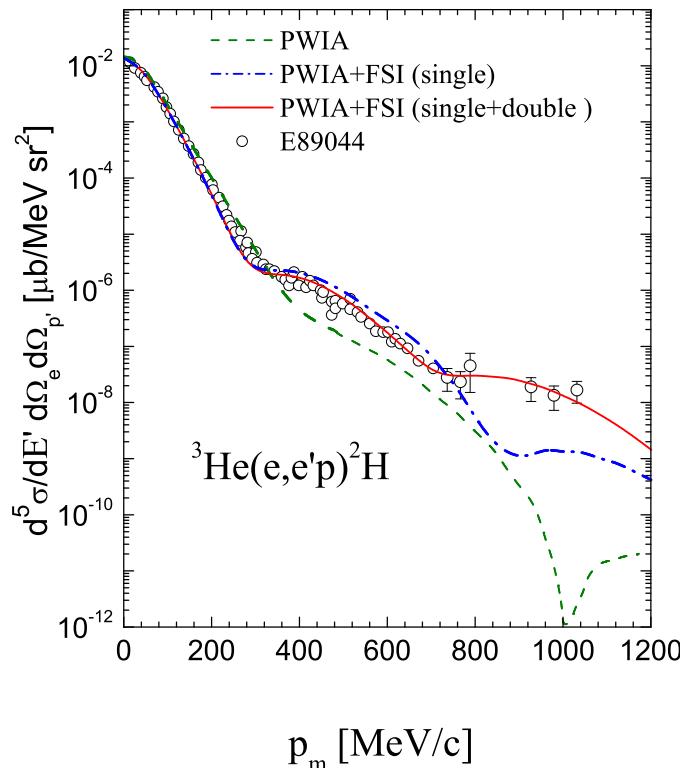


PRL 94 (2005) 192302

★ ★ ★ see also nucl-th/0705.3951 (**unfactorized calculations !**)

FSI in ${}^3\text{He}(\text{e}, \text{e}'\text{p}){}^2\text{H}$ at high p_m

- GEA: generalized eikonal approx (++rescatterings, $A - 1$ excitation) \leftarrow FSI
- Pisa WFs (AV18) \leftarrow initial-state correlations
- no MEC (small), no IC (small)
- **2bbu:** excellent agreement w/ GEA

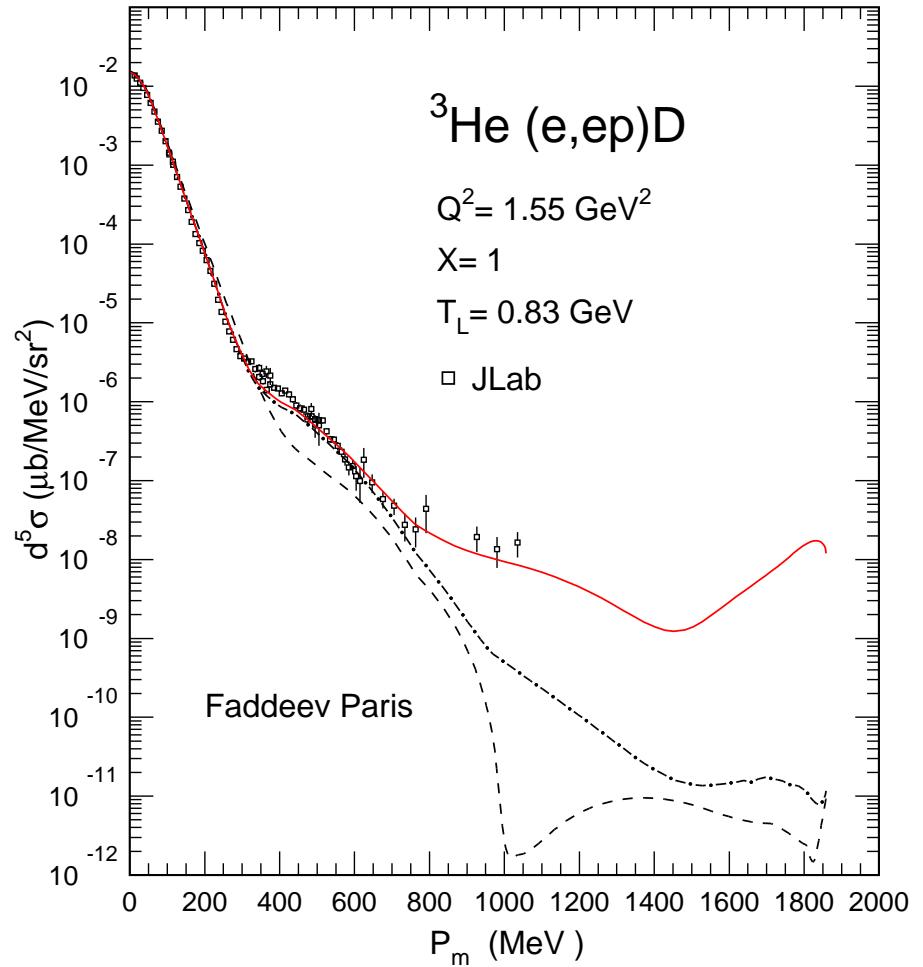
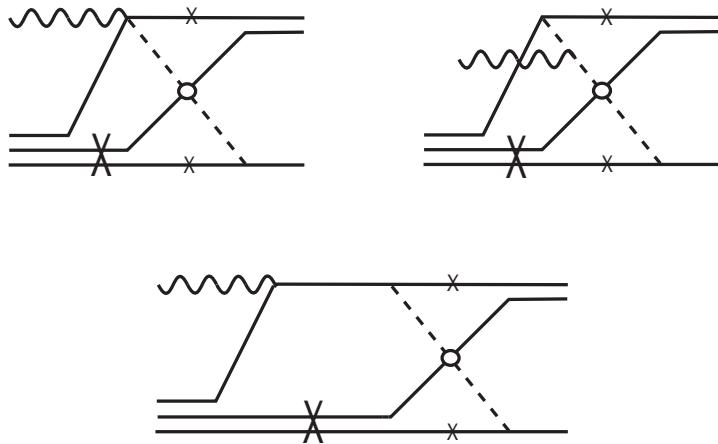


- **3bbu:** data underestimated by $\approx 20\%$ at $p_m \gtrsim 800 \text{ MeV}/c$ PRL 95 (2005) 052502

... however

$^3\text{He}(\text{e}, \text{e}'\text{p})^2\text{H}$ at high p_m

- particular 3-body mechanism at $p_m \gtrsim 700 \text{ MeV}/c$
- maximal in QE kinematics ($x = 1$)
- connects to pd-elastic at $\theta = \pi$

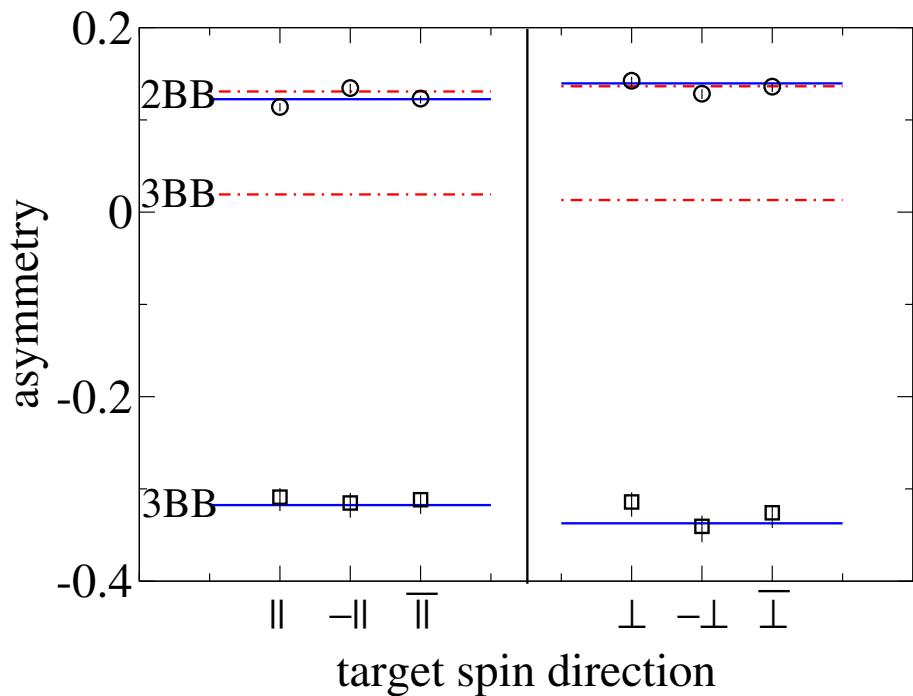
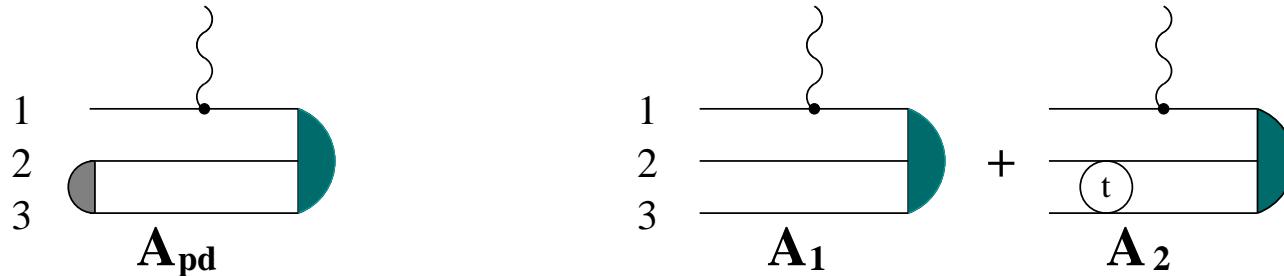


PRC 72 (2005) 024001, PRC 72 (2005) 064003

${}^3\vec{\text{He}}(\vec{e}, e' p) {}^2\text{H}$ and ${}^3\vec{\text{He}}(\vec{e}, e' p)\text{pn}$

MAMI/A1

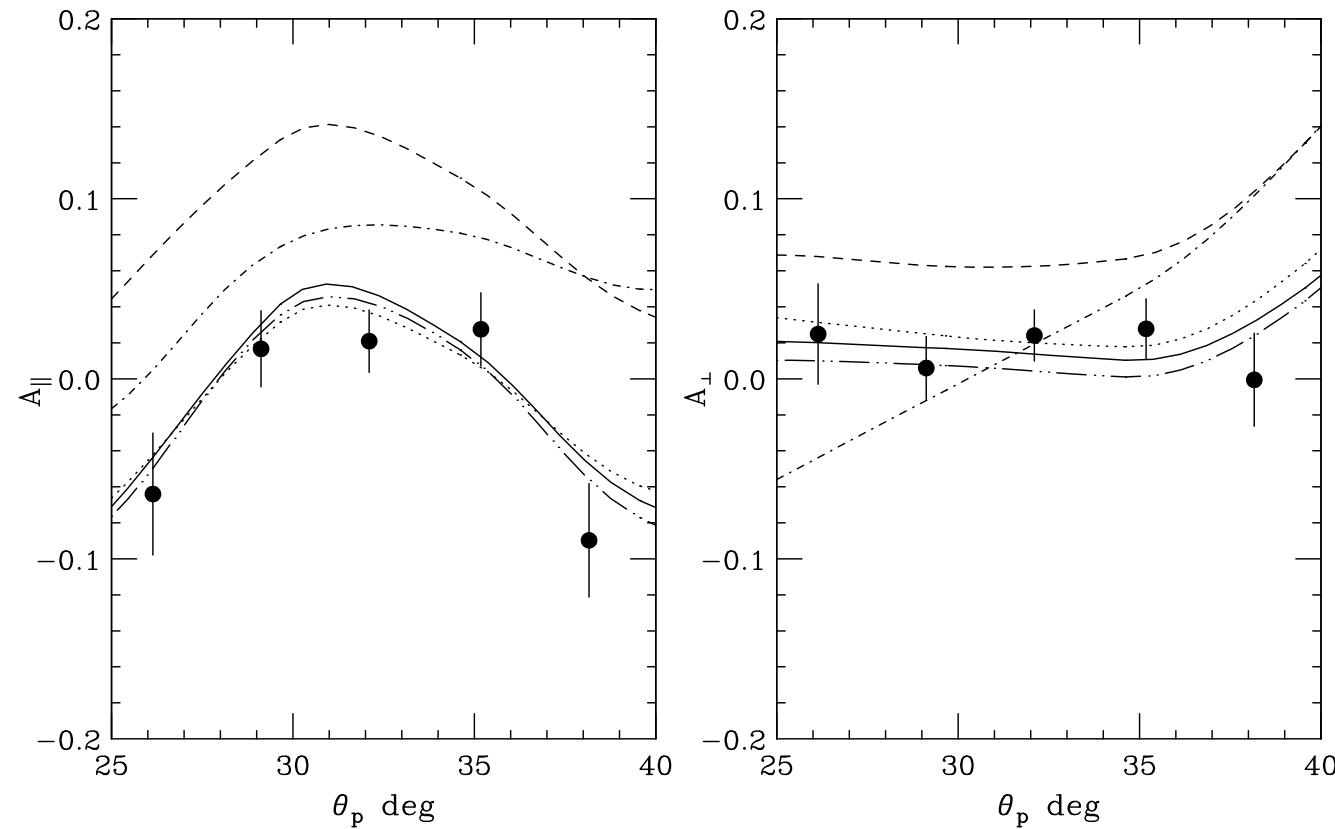
- QE ($Q^2 = 0.31$, $\omega = 135$, $q = 570$ central)
- 3NF, MEC negligible, FSI small in 2bbu, large in 3bbu



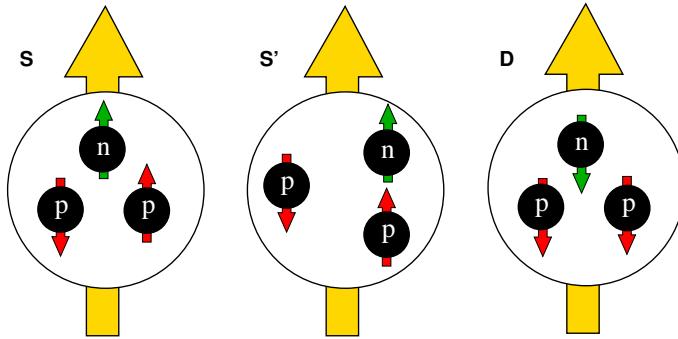
- ▷ **2bbu**
 $A_{\text{PWIA}} \approx A_{\text{PWIA+FSI}}$
 || kinematics + small p_d
 \Rightarrow polarized p target, $P_p \approx -\frac{1}{3}P_{\text{He}}$
- ▷ **3bbu**
 $A_{\text{PWIA}} \approx 0$ ($p \uparrow p \downarrow$)
 $A_{\text{PWIA+FSI}}$ large & negative
 not a polarized p target

PRC 72 (2005) 054005, EPJA 25 (2005) 177

- $Q^2 = 0.67$ at QE peak and on low- ω side
- non-relativistic calc no longer applicable
- relativistic kinematics + approximate FSI (A_2 -term) sufficient



- better understand ${}^3\text{He}$ as opposed to using it as effective n target
- any polarized ${}^3\text{He}$ exp depends on this to some extent



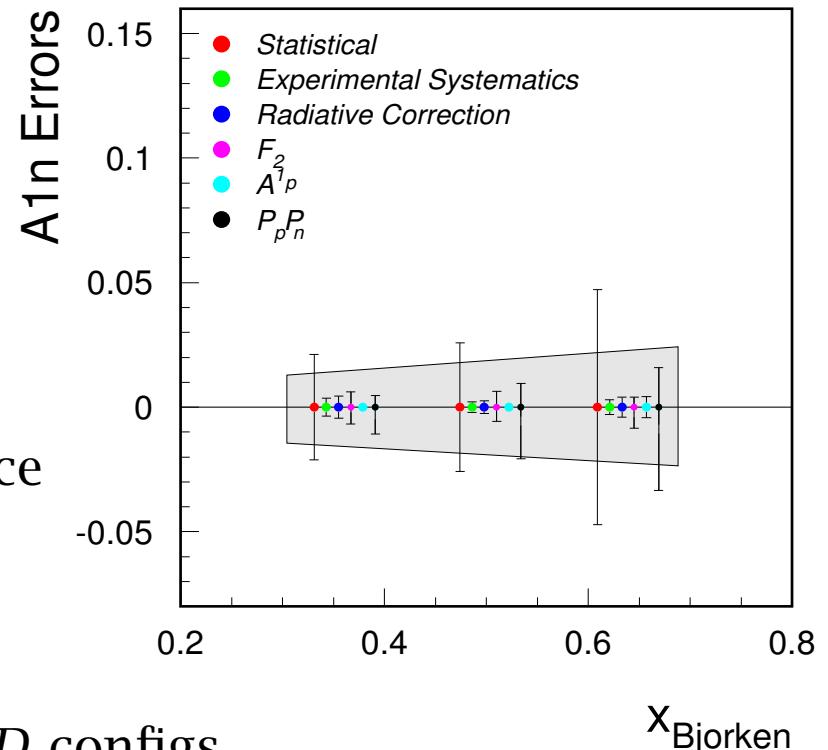
S (90% WF) space symmetric, $p \uparrow p \downarrow$

D (8.5% WF) tensor component of NN force

S' (1.5% WF) mixed symmetry config
(spin-isospin)-space correlations

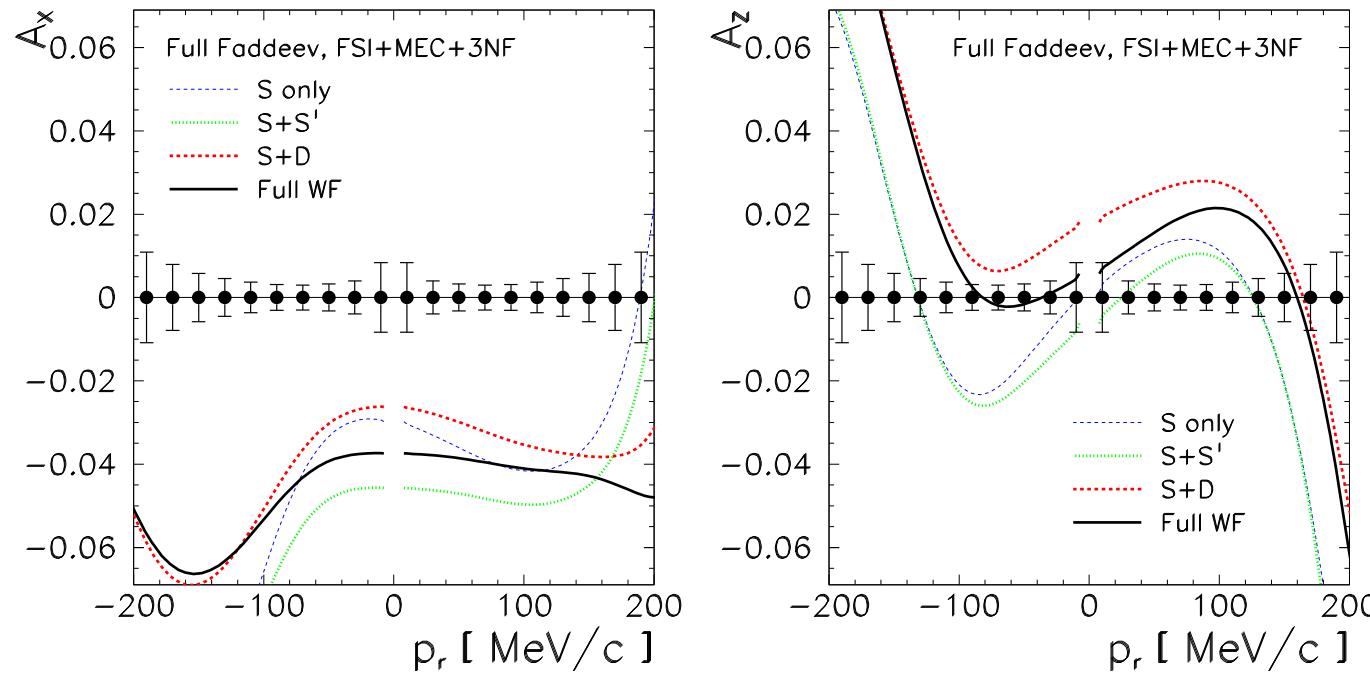
× protons partly polarized due to S' and D configs

- (iso)spin dependence of reaction mechanism (MEC, IC)



$$\frac{d\sigma(h, \vec{S})}{d\Omega_e dE_e d\Omega_d dp_d} = \frac{d\sigma_0}{\dots} \left[1 + \vec{S} \cdot \vec{A}^0 + h(\mathcal{A}_e + \vec{S} \cdot \vec{A}) \right]$$

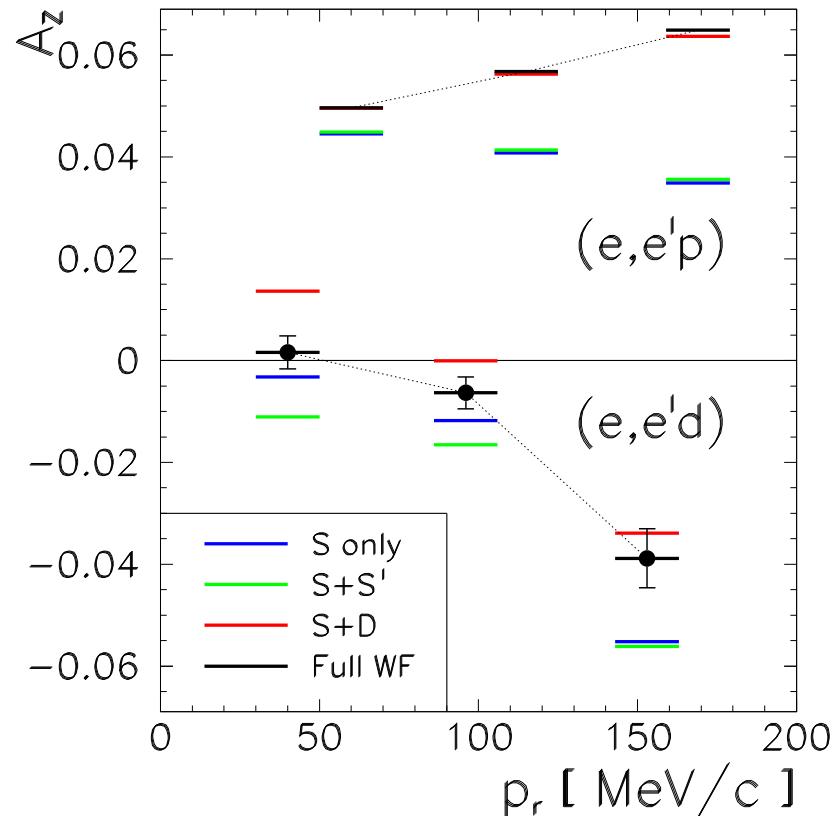
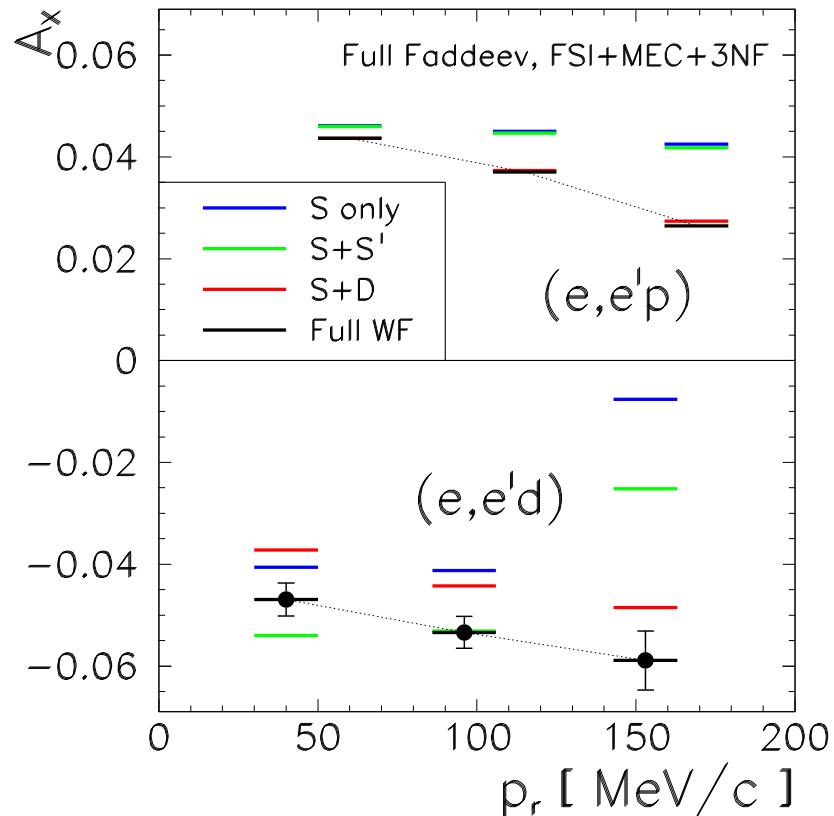
$$A_{x,z} = \frac{[d\sigma_{++} + d\sigma_{--}] - [d\sigma_{+-} + d\sigma_{-+}]}{[d\sigma_{++} + d\sigma_{--}] + [d\sigma_{+-} + d\sigma_{-+}]}$$



- sensitivity to small-WF components

${}^3\vec{\text{He}}(\vec{e}, e'd)p$ vs. ${}^3\vec{\text{He}}(\vec{e}, e'p)d$

E05-102



- no sensitivity to S' in proton channel

Triple polarization ${}^3\text{He}(\vec{e}, e' \vec{p})$

MAMI/A1

- spin-dependent momentum distributions of $\vec{p}\vec{d}$ clusters in polarized ${}^3\text{He}$

$$N_\mu = \langle \Psi_{pd}^{(-)} M_d m | \hat{j}_\mu(\vec{q}) | \Psi M \rangle$$

$$\gamma \left(M = \frac{1}{2}, M_d = 0, m = +\frac{1}{2} \right) \propto \left| N_{-1}^{\text{spin PWIA}} \left(\frac{1}{2}, 0, -\frac{1}{2} \right) \right|^2$$

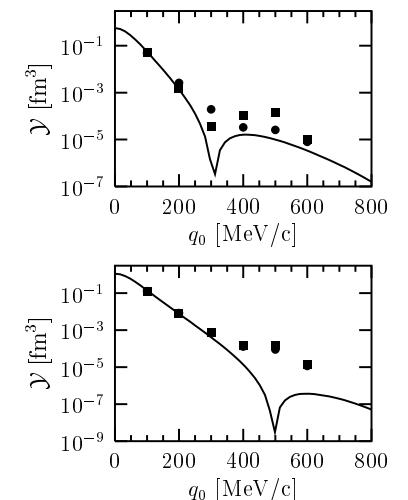
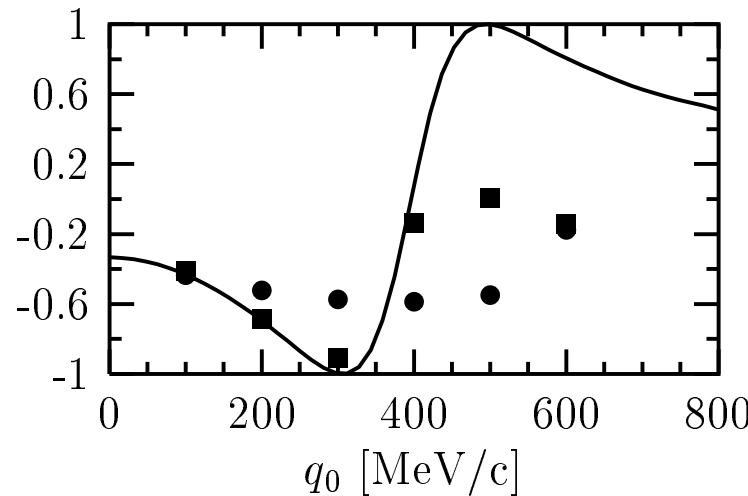
$$\gamma \left(M = \frac{1}{2}, M_d = 1, m = -\frac{1}{2} \right) \propto \left| N_{+1}^{\text{spin PWIA}} \left(\frac{1}{2}, 1, +\frac{1}{2} \right) \right|^2$$

$$A = \frac{\gamma(1/2, 0, 1/2) - \gamma(1/2, 1, -1/2)}{\gamma(1/2, 0, 1/2) + \gamma(1/2, 1, -1/2)}$$

$$\sigma_L \propto |N_0|^2$$

$$\sigma_T \propto |N_{+1}|^2 + |N_{-1}|^2$$

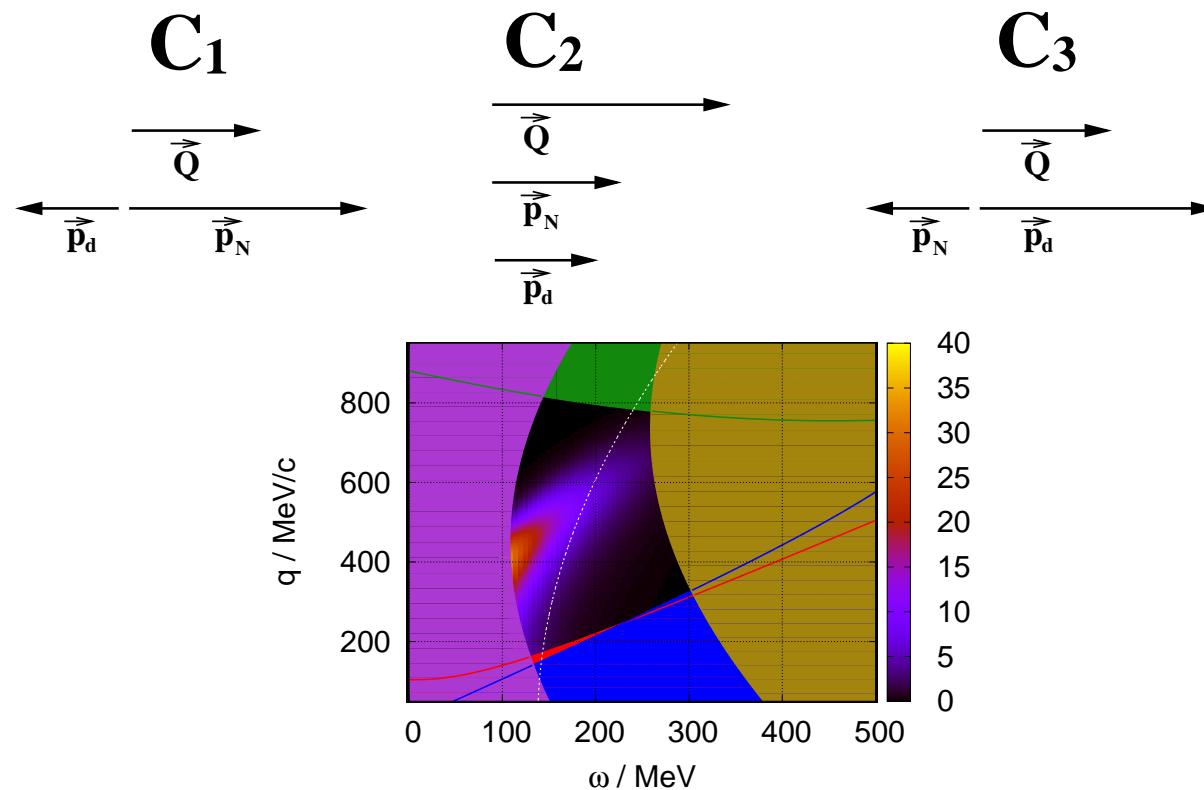
$$\sigma_{T'} \propto |N_{+1}|^2 - |N_{-1}|^2$$



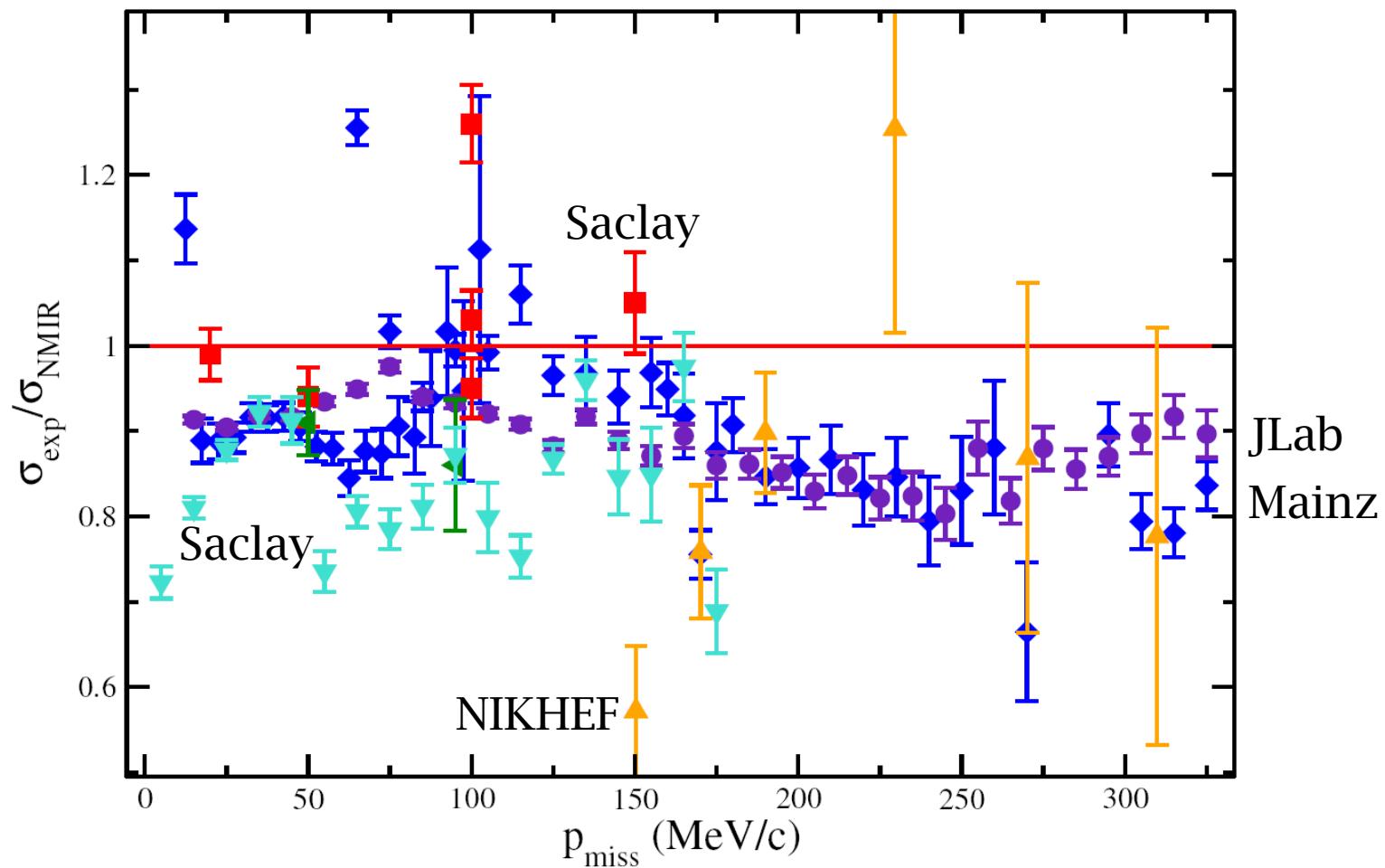
Triple polarization ${}^3\vec{\text{He}}(\vec{e}, e' \vec{p})$

MAMI/A1

- PWIA: σ_L , σ_T , $\sigma_{T'}$ yield spin-dependent momentum distribution
- FSI, MEC preclude direct access except at $p_d \lesssim 2 \text{ fm}^{-1}$
- rich interplay \triangleright final-state symmetrization: large effect in C_3
 - \triangleright FSI: largest in C_2
 - \triangleright MEC: most prominent in C_1

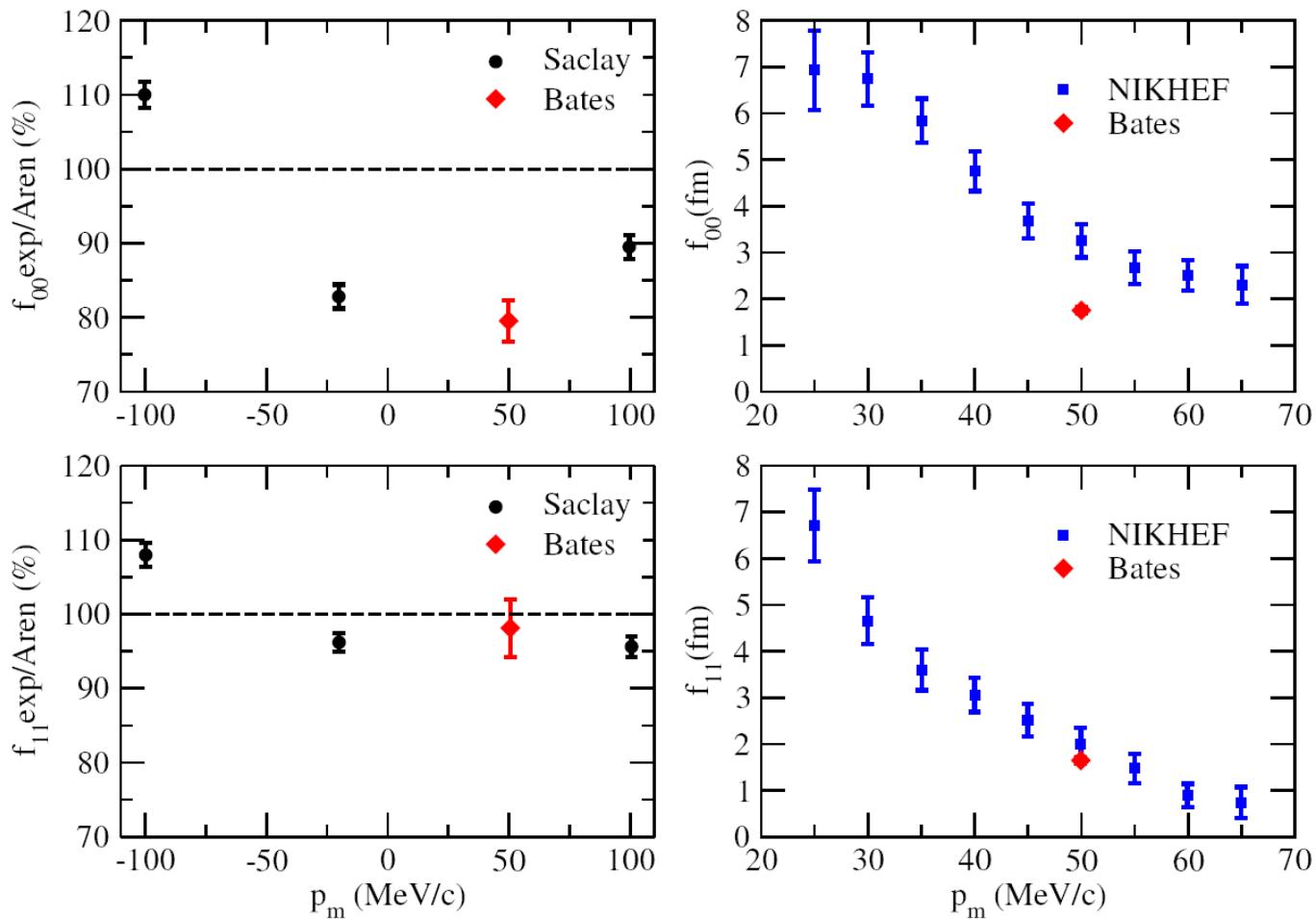


$d(e, e'p)$



- experiments \sim self-consistent
- all are $\sim 10\%$ below theory

L/T separation in $d(e, e'p)$

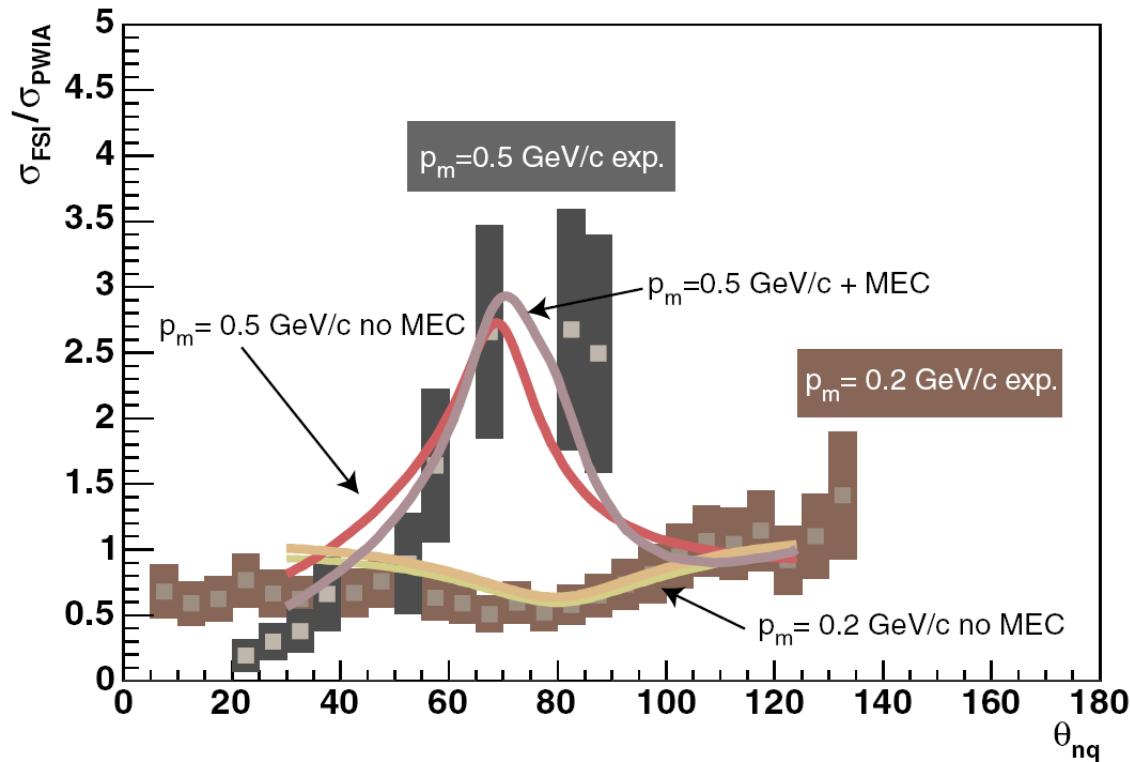


- also L/T data from Mainz but experimental error at low p_m
- discrepancies in L and T and no experimental program to address them

$d(e, e'p)$ at high Q^2

E01-020

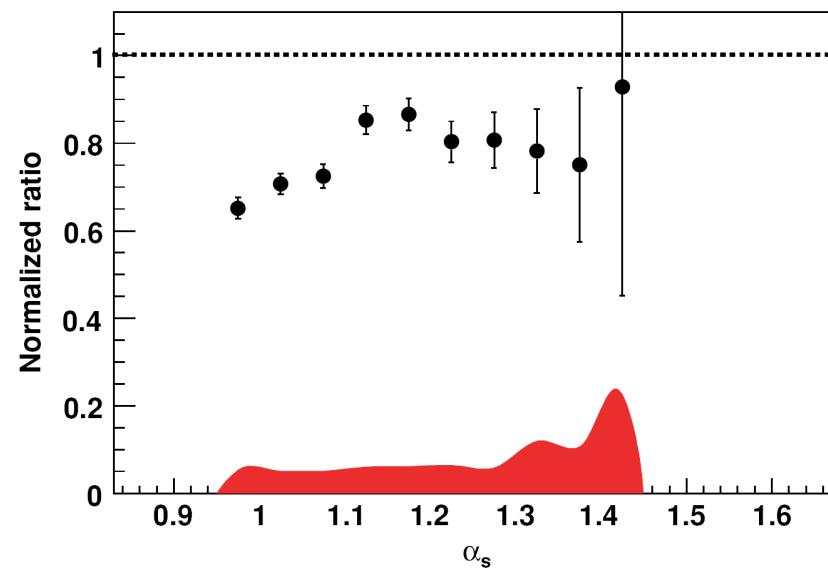
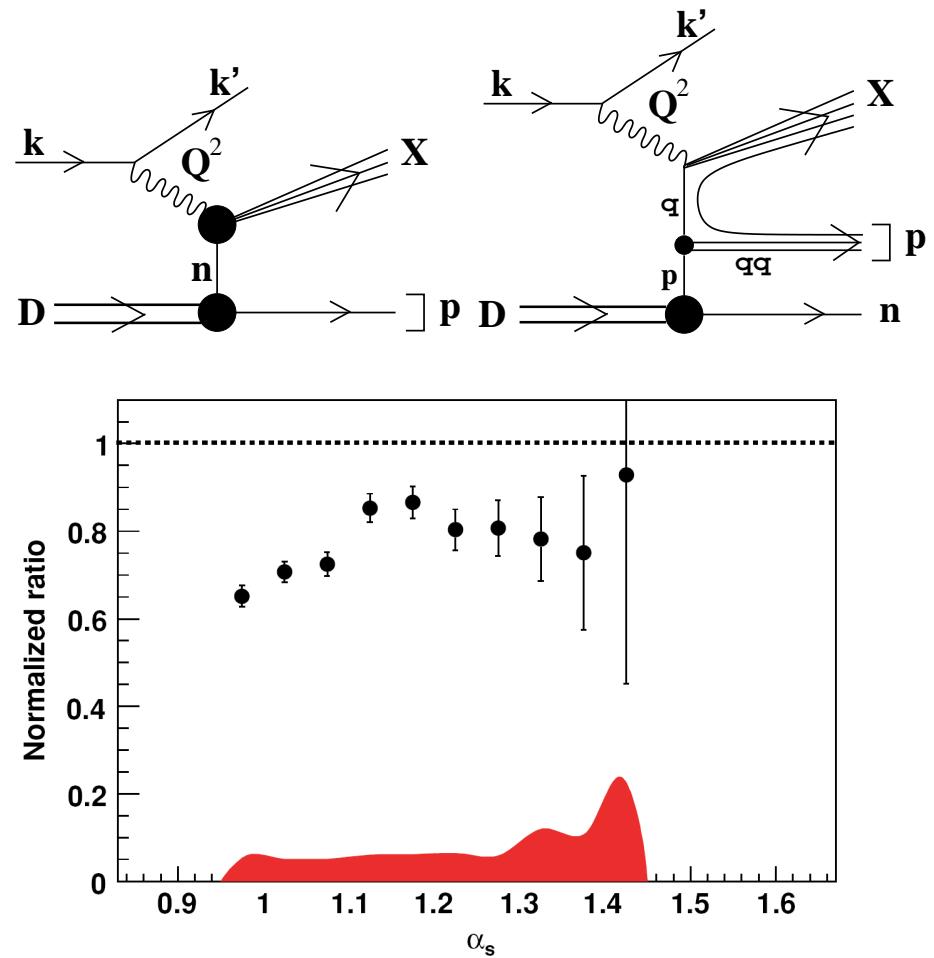
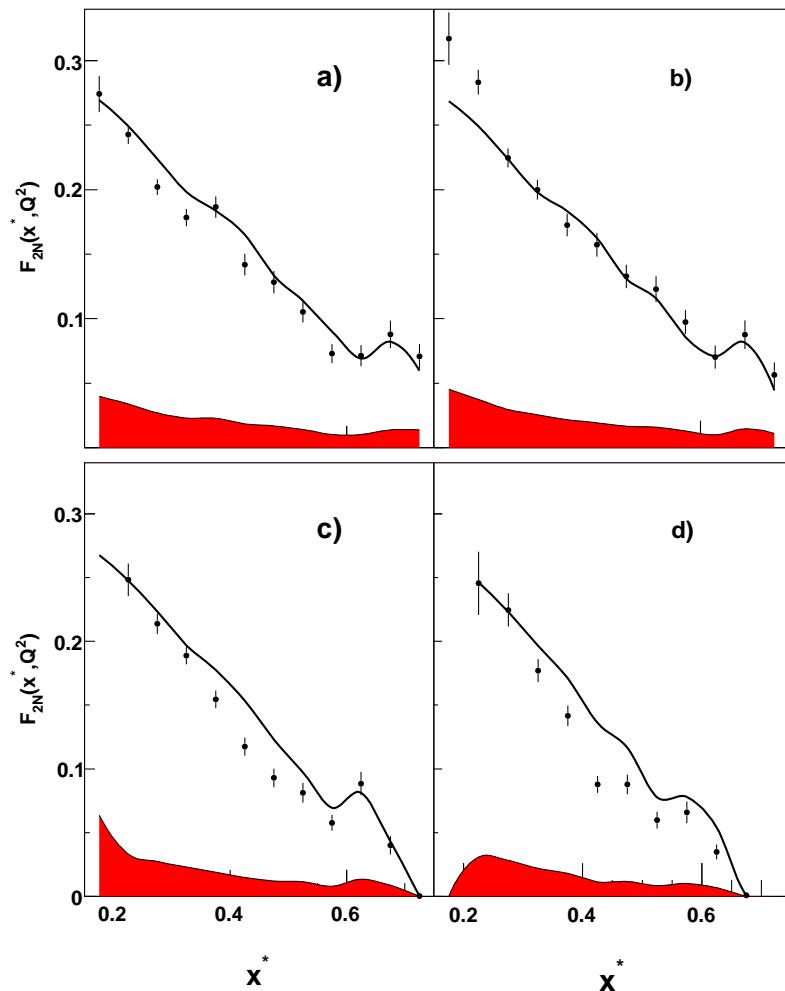
- $Q^2 = 0.8, 2.1, 3.5 \text{ GeV}^2$ for $0 \lesssim p_m \lesssim 500 \text{ MeV}/c$
- test Generalized Eikonal Approximation of FSI
- characteristic θ_{nq} -dependence with predicted max. at 80°



PRELIMINARY EPJA 28 s01 (2006) 19
see also CLAS E94-019

Semi-inclusive deep-inelastic ${}^2\text{H}(\text{e}, \text{e}'\text{p}_S)\text{X}$

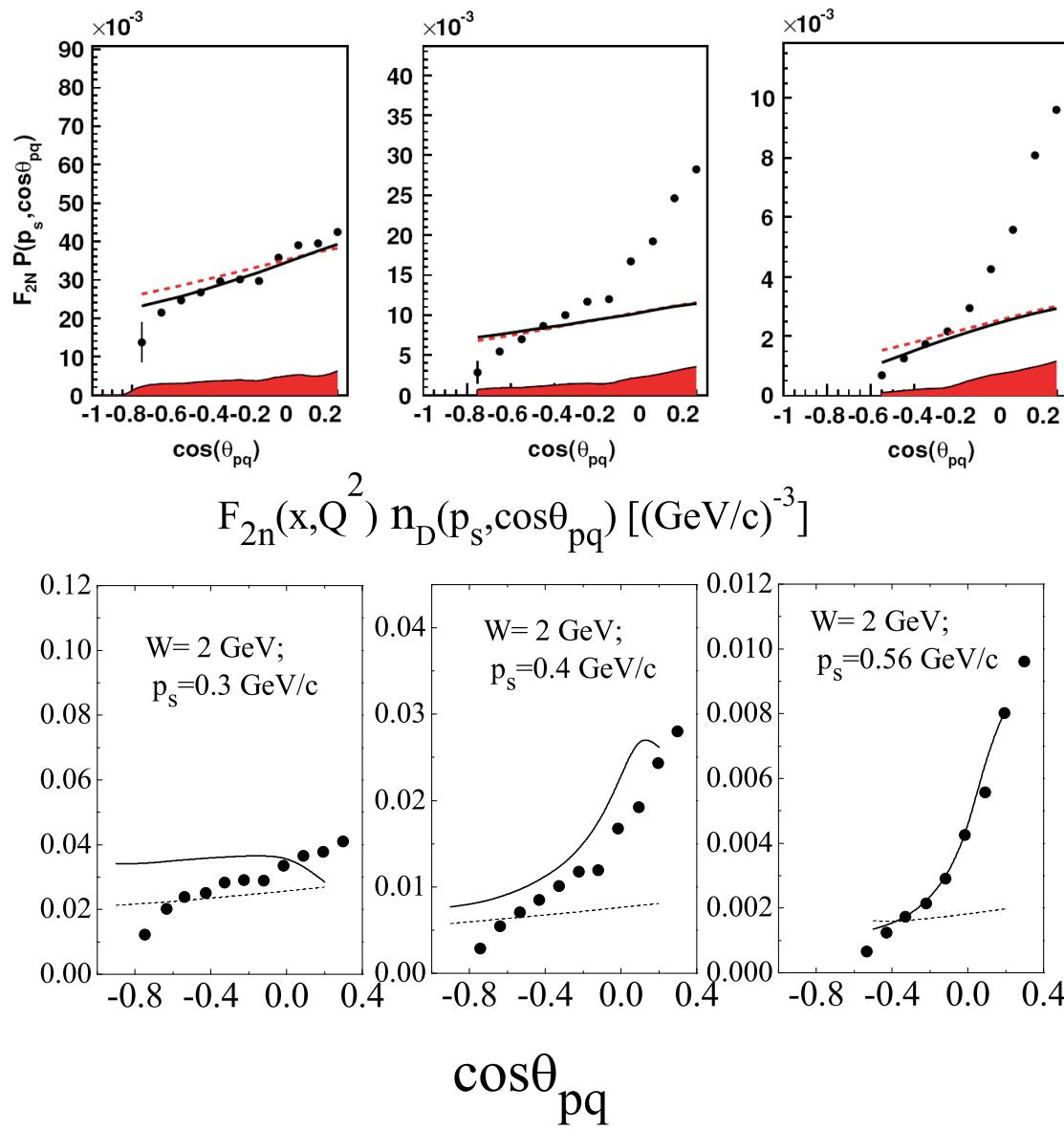
- regions where FSI minimized \Rightarrow investigation of DIS structure functions
- regions where FSI maximized \Rightarrow study hadronization mechanisms



PRC 73 (2006) 035212

FSI in $^2\text{H}(\text{e}, \text{e}'\text{p}_s)\text{X}$

$W = 2 \text{ GeV}$

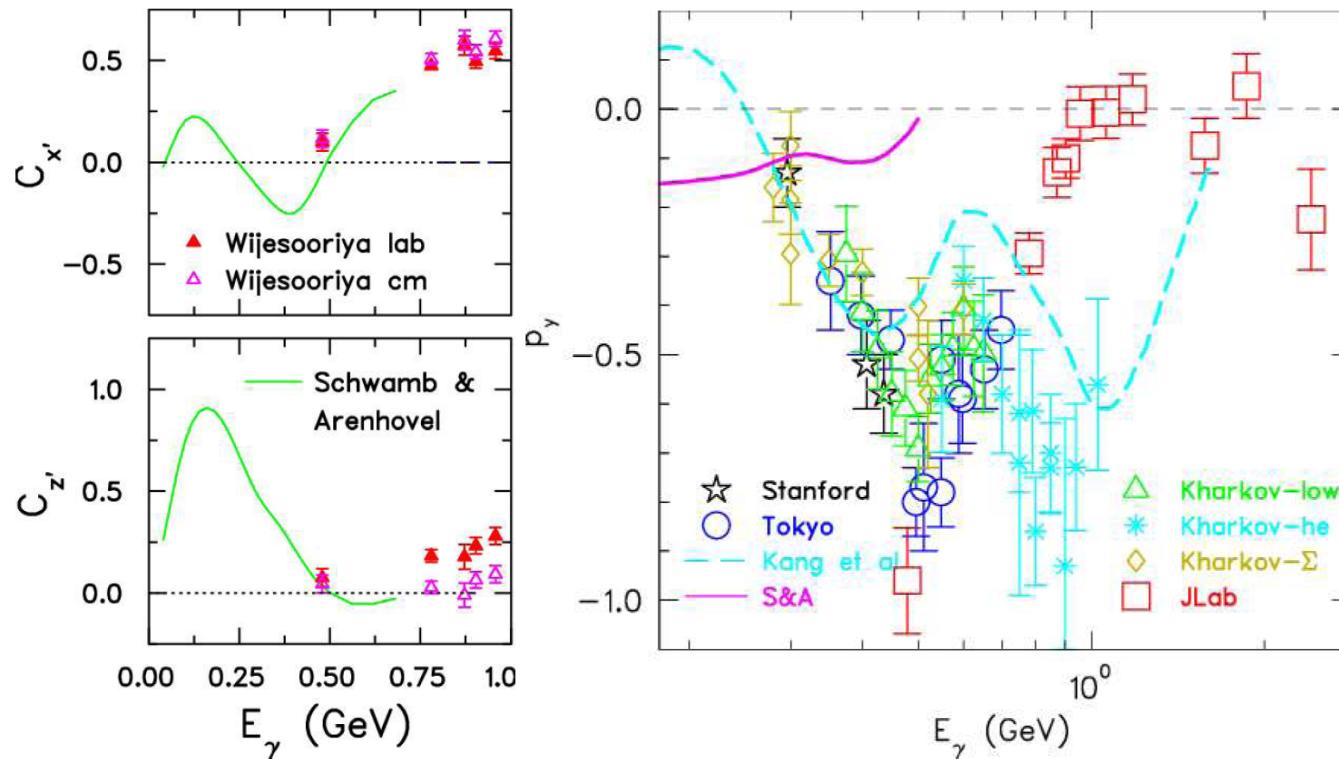


nucl-th/0705.3617

Polarization observables in $\vec{y}d \rightarrow \vec{p}n$

E05-103

- agreement w.r.t. theory for C'_x and C'_z but not for p_y PRL 98 (2007) 182302
- but $\sigma(\theta)C'_x = 2 \operatorname{Re} \sum_{i=1}^3 [F_{i,+}^* F_{i+3,-} + F_{i,-} F_{i+3,+}^*]$ PRL 86 (2001) 2975
- $\sigma(\theta)p_y = 2 \operatorname{Im} \sum_{i=1}^3 [F_{i,+}^* F_{i+3,-} + F_{i,-} F_{i+3,+}^*]$

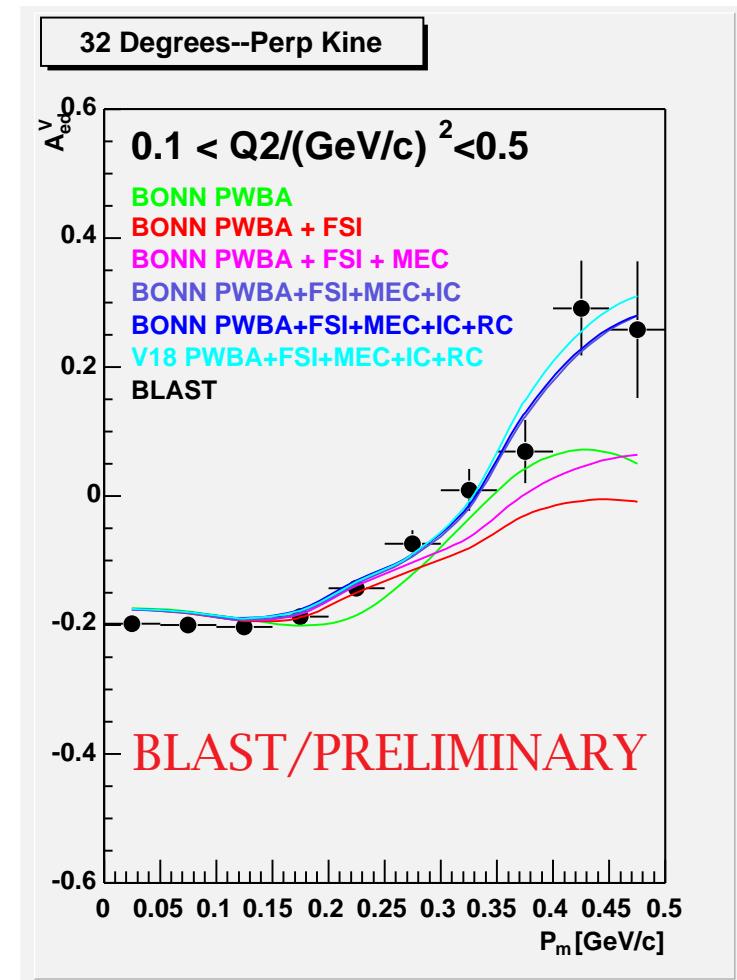
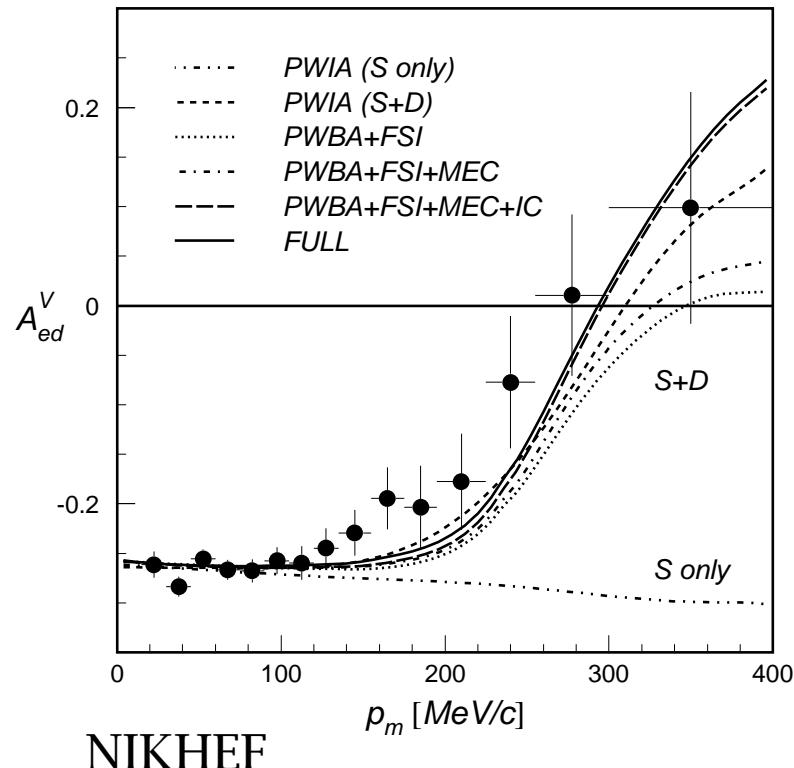


- clean up the mess in existing data, in particular on p_y
- input to (state-of-the-art) theory beyond $E_\gamma \sim 300$ MeV

Spin-correlation parameter A_{ed}^V in ${}^2\vec{\text{H}}(\vec{e}, e' p)$

BLAST

$$\sigma = \sigma_0 \left\{ 1 + P_1^d A_{\text{d}}^V + P_2^d A_{\text{d}}^T + h \left(A_{\text{e}} + P_1^d A_{\text{ed}}^V + P_2^d A_{\text{ed}}^T \right) \right\}$$

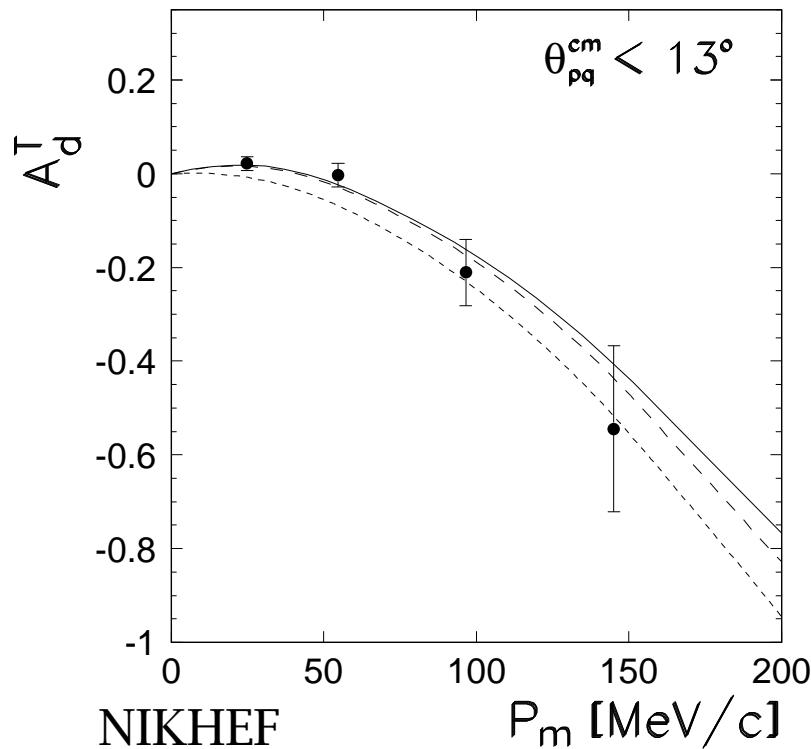


Tensor analyzing power A_d^T in ${}^2\vec{H}(e, e' p)$

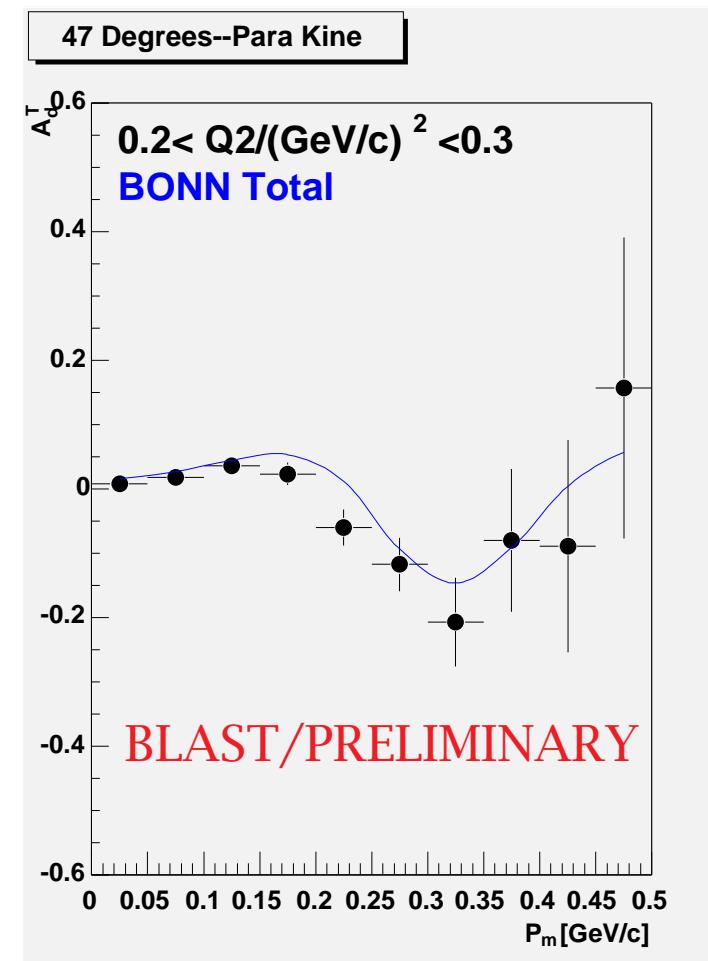
BLAST

$$\sigma = \sigma_0 \left\{ 1 + P_1^d A_d^V + P_2^d A_d^T + h \left(A_e + P_1^d A_{ed}^V + P_2^d A_{ed}^T \right) \right\}$$

PWIA : $A_d^T \propto \left(2R_0(p)R_2(p) + \sqrt{\frac{1}{2}}R_2(p)^2 \right) / \left(R_0(p)^2 + R_2(p)^2 \right)$



PRL 82 (1999) 687



Conclusions and outlook

- ▷ Complex nuclei
 - Breakdown of factorization
 - Dynamical relativity
 - Q^2 -dependence of spectroscopic factors
 - Optical-potential vs. Glauber approaches to FSI
- ▷ ${}^3\text{He}$
 - Unpolarized \approx OK except very high p_m
 - $(e,e'd)$ vs. $(e,e'p)$ correspondence (g.s. WF components)
 - Intricate interplay of FSI, MEC, IC
 - Triple polarization
- ▷ Deuteron
 - Unresolved discrepancies in XS, L/T responses
 - Simultaneous description of C'_x and p_y
 - Upcoming BLAST data on A_{ed}^V and A_d^T
 - Pion and resonance physics with polarization

Dynamical relativity (“spinor distortion”) in A(e, e'p)

$$\begin{aligned} [\vec{\alpha} \cdot \vec{p} + \beta(M + S_b) + (V_b - E_b)] \Phi &= 0 && \text{bound} \\ [\vec{\alpha} \cdot \vec{p} + \beta(M + S_c) + (V_c - E_c)] \Psi &= 0 && \text{continuum} \end{aligned}$$

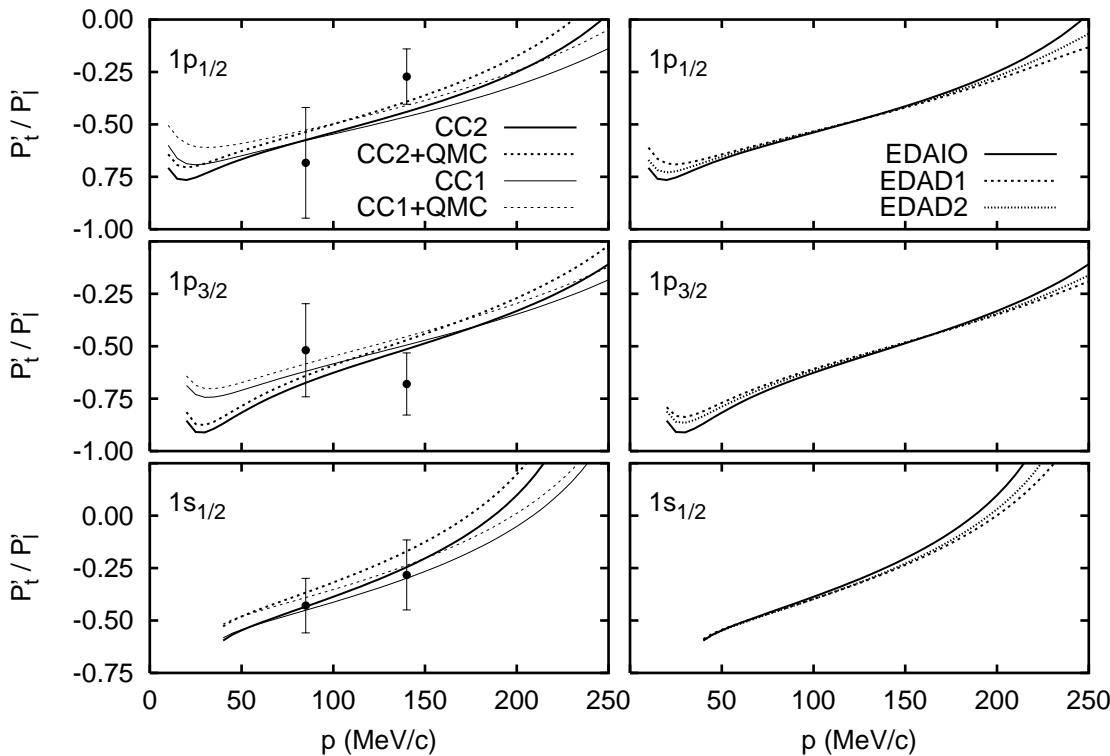
$$\begin{aligned} \Rightarrow & \left[\nabla^2 + E^2 - M^2 - 2E \left(\textcolor{red}{U}^C + \textcolor{red}{U}^{LS} \vec{L} \cdot \vec{\sigma} \right) \right] \xi = 0 \\ \textcolor{red}{U}^C &= V + \frac{M}{E}S + \frac{S^2 - V^2}{2E} + \frac{1}{2E} \left[-\frac{1}{2Dr^2} \frac{d}{dr} (r^2 D') + \frac{3}{4} \left(\frac{D'}{D} \right)^2 \right] \\ \textcolor{red}{U}^{LS} &= -\frac{1}{2Er} \frac{D'}{D} \\ D(\textcolor{blue}{r}) &= 1 + \frac{S(r) - V(r)}{E + M} \end{aligned}$$

$$\Phi = \Omega_b \xi_b, \quad \Psi = \Omega_c \xi_c, \quad \Omega(\vec{p}, r) = \left(\frac{1}{\frac{\vec{\sigma} \cdot \vec{p}}{(E+M)D(\textcolor{blue}{r})}} \right) \textcolor{red}{D}^{1/2}(\textcolor{blue}{r})$$

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$^{16}\text{O}(\vec{e}, e' \vec{p})$: polarization transfer

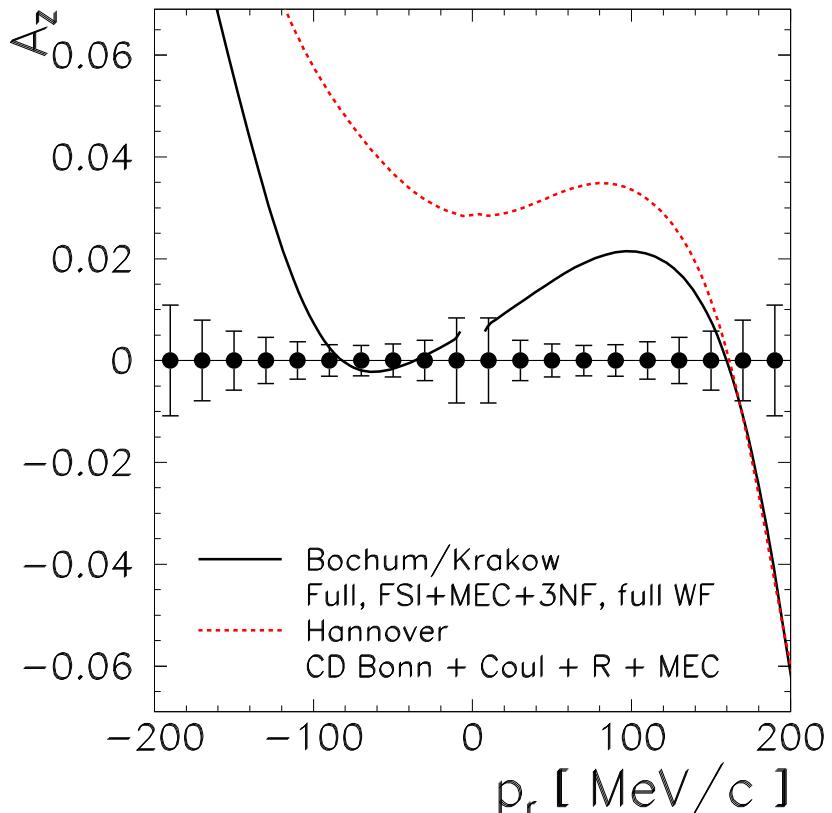
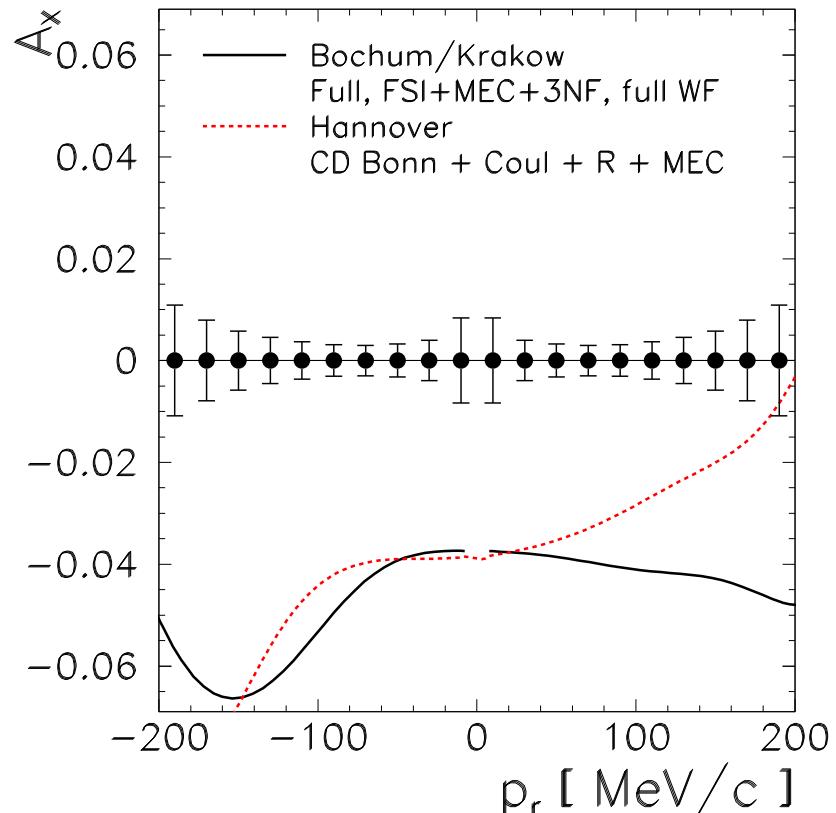
- quest for modification of proton FF inside the nuclear medium
- pol. observables least sensitive to most nuclear structure uncertainties



- ▷ dynamical relativity
distortion of bound and scattered Dirac spinors
- ▷ kinematical rel. effects
relativistic kinematics
relativistic N current operator
- ▷ realistic description of FSI
relativistic optical potentials
- ▷ medium-modified FFs
density dependence

$$G_{E,M}(Q^2, \rho) = G_{E,M}^{\text{GK}}(Q^2) \frac{G_{E,M}^{\text{QMC}}(Q^2, \rho)}{G_{E,M}^{\text{QMC}}(0, \rho)}$$

PRC 69 (2004) 034604



- discrepancies in theories everyone believes in
- expected to run in 2008/09