# **Graal Apparatus**























FIG. 1: Upper part: The correlation  $\Delta\theta$  vs.  $(\Delta\phi-180^{\circ})$  in three-dimensional view(a) and its projection on two dimensions (b); Lower part: The bidimensional gaussian fit (c) and its projection on two dimensions (d). 070904 EINN 2007



FIG. 2: The three-dimensional correlation  $E_{\eta}^{calc}/E_{\eta}^{meas}$  vs.  $(M_X - M_N)$  (a) (see text for explanation) and its projection on two dimensions (b).



FIG. 5: Left: Coplanarity between the  $\eta$  and the free proton (solid line) and the quasi-free proton (dotted line): the smearing between the two distributions is due to the Fermi motion; Right: The Fermi momentum distribution before (solid line) and after (dashed line) the application of the bidimensional cuts (see text for details).



FIG. 3: The  $\eta$  invariant mass without cuts (solid line) and with the kinematical cuts (dotted line) for a central proton (a) and neutron (b) (in logarithmic scale), for a forward proton (c) and neutron (d). U70904 EINN





FIG. 6: Beam asymmetry  $\Sigma$  in  $\eta$  photoproduction on the quasi-free proton (open squares) in the deuteron and on the free proton (full circles)[2]. The energy value outside and inside parenthesis indicate the mean value of the bin for quasi-free and free protons respectively. In dotted lines are illustrated the predictions of Maid2001 [7] for the free proton, in solid and dashed lines those for the quasi-free proton of Maid2001 [7] and the reggeized model [3], respectively (see text for details).

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FIG. 7: Energy dependence of the differential cross section for  $\gamma p \rightarrow \eta p$  calculated at  $\theta = 50^{\circ}$ . The unpolarized cross section (5), the polarized cross section (6) and the beam asymmetry  $\Sigma$  are presented in the three pairs of curves for the free proton (dashed curves) and the quasi free proton (solid curves) respectively.



FIG. 8: Beam asymmetry  $\Sigma$  in  $\eta$  photoproduction on the quasi-free neutron in eleven energy bins, plotted as a function of the  $\theta_{\eta}^{cm}$ . In each plot, the mean  $\gamma$  energy of the bin is also indicated. In solid and dashed lines are illustrated the predictions for neutrons of Maid2001 and of the reggeized model respectively (see text for details).

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FIG. 9: Comparison between the beam asymmetry  $\Sigma$  in  $\eta$  photoproduction on the quasi-free proton (open squares) and the quasi-free neutron (full triangles) in the eleven energy bins, plotted as a function of the  $\theta_{\eta}^{cm}$ . See text for details. 070904 EINN 2007 18



FIG. 10: Comparison between the beam asymmetry  $\Sigma$  in  $\eta$  photoproduction on the quasi-free proton (open squares) and the quasi-free neutron (full triangles) in seven angular bins, plotted as a function of the  $\gamma$  energy (intervals of  $\simeq 25$  MeV width.

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# $\gamma + n + (p) \rightarrow \eta + n + (p)$ Yield(E<sub> $\gamma$ </sub>)

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 $\gamma + n + (p) \rightarrow \eta + n + (p)$  Yield(W)







FIG. 2: Total cross section of the reaction  $\gamma p \rightarrow \eta \pi^0 p$ . The dots are the experimental data of this work. The open circles are from reference [9]. The results of the model of Ref. [6, 7] are given with their uncertainty by a hatched band of the figure. The uncertainty originates from the one on the  $\gamma p \Delta(1700)$  coupling which was taken from the PDG [10]

 $\mathbf{O}_{\mathsf{tot}}$ 

 $\gamma + p \rightarrow \pi^0 + \eta + p$ IM



FIG. 3: For the reaction  $\gamma p \to \eta \pi^0 p$ , spectra of invariant mass of  $(p\pi^0)$ ,  $(p\eta)$  and  $(\eta\pi^0)$  groups of the final state, presented in 3 different columns. The various rows, labeled (a),...,(e) correspond to beam energies given at the bottom of the figure. In empty circles are the experimental results, in thin line for  $\gamma p \to \eta \pi^0 p$  with a 3-body phase space in the final state, in dashed line for  $\gamma p \to \eta \Delta$  with  $\Delta \to \pi^0 p$ , in dotted lines for  $\gamma p \to \pi^0 S_{11}$  with  $S_{11} \to \eta p$ . The theoretical curves, given by thick lines, are the central values of the results of the model of Ref. [6, 7].

> It looks like:  $\gamma + p \rightarrow \eta + \Delta$  $\Delta \rightarrow \pi^{\circ} + p$



FIG. 4: Beam asymmetry of the reaction  $\gamma p \rightarrow \eta \pi^0 p$ . The theoretical results are calculated with the model of Ref. [6, 7]

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## **Compton Scattering Kinematics**

2. The maximum energy lost by the electrons after an elastic scattering with a laser photon is given by the maximum energy acquired by the photon:

$$E_{el}^{0} - E_{el}^{scatt} = E_{\gamma \max} = \frac{4\gamma^{2}E_{laser}}{1 + \frac{4\gamma E_{laser}}{m_{o}}} \approx 4\gamma^{2}E_{laser}$$

This energy loss is measured by the displacement **d** of the scattered electrons from the primary electron beam after the first magnetic dipole. For the ESRF electron energy of 6.03 GeV and a UV laser line of 3.53 eV, the energy loss is 1.487 GeV and corresponds to an electron displacement at the position of the Graal tagging detector:  $d \approx 52.3$  mm.

The microstrips of the Graal tagging detector measure the displacement d of the scattered electrons from the main orbit and therefore the energy lost by the electrons (and acquired by the gamma-rays):

 $E_{\gamma} \propto d$ 

## **Compton Scattering Kinematics**



# **Graal Tagging Microstrips**

Schematic description of the tagging detector in more details. Vertical cut: the electrons fly out of the screen.





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### **Daily Compton Edges Distributions**



Experimental data plotted as a function of (solar) hour, showing their daily variation. The dispersion of data around the average, taken arbitrarily at zero, is expressed in fractions of microstrip (300 micrometers width or about 7 MeV for one microstrip).



Same as the previous figure, but each point is the average over one hour. The dotted lines show the refill time of the machine corresponding to a possible change in the temperature of the tagging detector or the position of the beam. The average is expressed in microstrip fractions (0.01 =  $3\mu$ m).

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## **Graal Beam Orientation on the Earth**





## **Compton Edge Positions vs CMB Dipole**



Experimental data plotted as a function of the azimuth (above); below, the variation of the angle between the beam and the CMB dipole decomposed to azimuth (dotted) and declination (dashed) angles is shown.

Assuming an error of

2 · 10<sup>-4</sup>

in our determination of the position of the Compton Edge, we could arrive to an estimated upper limit on the asymmetry of the velocity of light of:

$$\Delta\beta \approx \frac{1}{2} \left(\frac{1}{\gamma^2}\right) \frac{\Delta E_{\gamma}}{E_{\gamma}} \approx 0.4 \cdot 10^{-8} \frac{\Delta d}{d} \approx 0.4 \cdot 10^{-8} \cdot 2 \cdot 10^{-4} \approx 10^{-12}$$

Considering that we have analyzed old data and we have not been able to reconstruct completely the status of the system - accelerator + tagging detector - during our runs we have published the more conservative number:

3 · 10<sup>-12</sup>

## An Optimistic View of the Future

In conclusion if optimistically we assume a systematic error of 2.5  $\mu$ m in the distance between the position of the Compton Edge and the electron beam, we have:

$$\frac{\left(\Delta d\right)_{sys}}{d} \approx \frac{2.5\,\mu m}{52.3\,mm} \approx 5\cdot 10^{-5} \approx \frac{\left(\Delta E_{\gamma}\right)_{stat}}{E_{\gamma}}$$

and we can hope to be able to verify the isotropy of the velocity of light with respect to some absolute reference frame with a precision of:

$$\Delta\beta \approx \frac{1}{2} \left(\frac{1}{\gamma^2}\right) \frac{\Delta E_{\gamma}}{E_{\gamma}} \approx 0.4 \cdot 10^{-8} \frac{\Delta E_{\gamma}}{E_{\gamma}} \approx 0.4 \cdot 10^{-8} \cdot 5 \cdot 10^{-5} \approx 2 \cdot 10^{-13}$$