

# Nucleon Form Factors

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Proton and Neutron FFs in  
space-like domain  
(focus on new development)

EINN-2007

Milos

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# Outline

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- Introduction - EM FFs,
  - » New parameterizations
- Recent experimental results -
  - » low  $Q^2$
  - » high  $Q^2$
- Theory links observations and ideas -
  - » Charge density
  - » GPDs fit from FFs
- Experimental outlook - GEP-III, ..., 12 GeV
- Summary

# Reviews and Analysis

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## Reviews

Ch. Hyde-Wright and C. de Jager, *Ann. Rev. Nucl. Part.Sci.* 54, 217 (2004)  
H. Gao, *Int. J. Mod. Phys A*20, 1595 (2005)  
Ch. Perdrisat, V. Punjabi, and M. Vanderhaeghen, [hep-ph/06012014](#)  
J. Arrington, C.D. Roberts, and J.M. Zanotti, [nucl-th/0611050](#)

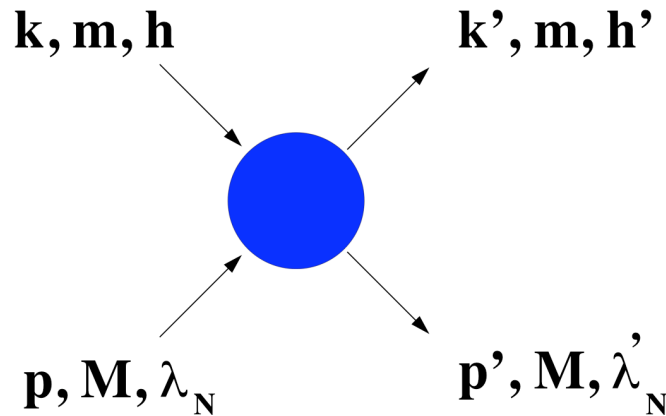
## Parameterizations

P. Bosted, *PRC* 64, 409 (1995) - Fit of pre-JLab data  
E. Brash et al. *PRC* 65, 051001(R) 2002 - Fit with JLab GEP/GMP  
J.J. Kelly, *PRC* 66, 065203 (2002) - Breit frame densities  
[BABB](#), [arXiv:hep-ex 0708.1946](#) - Fit with local duality constraints  
J. Arrington, W.Melnitchouk, J.A.Tjon, [nucl-ex/0707.1861](#) - 2-gamma

## GPD

M.Diehl et al., *Eur.Phys.J.* C39 (2005) 1-39, [GPDs from FFs data](#)  
M.Guidal et al., *PRD* 72, 054013 (2005) , [FFs from GPDs](#)

# Lepton-Nucleon scattering



$$l(k, h) + N(p, \lambda_N) \rightarrow l(k', h') + N(p', \lambda'_N)$$

$h, h', \lambda_N,$  and  $\lambda'_N$  are helicities

$$P = \frac{p+p'}{2}, \quad K = \frac{k+k'}{2}, \quad q = k - k' = p' - p$$

$$s = (p + k)^2, \quad t = q^2 = -Q^2, \quad u = (p - k')^2$$

$$T_{\lambda'_N, \lambda_N}^{h', h} \equiv \langle k', h'; p', \lambda'_N | T | k, h; p, \lambda_N \rangle$$

Total 16 amplitudes.

Parity invariance  $\rightarrow$  number of independent helicity amplitudes from 16 to 8.

Time reversal invariance  $\rightarrow$  to 6.

When neglect the lepton mass  $\rightarrow$  to 3.

Three complex amplitudes

$$T_{+,+}^{+,+}; \quad T_{-,-}^{+,-}; \quad T_{-,+}^{+,+} = T_{+,-}^{+,-}$$

which are functions of  $(s - u)$  and  $t$ .

# Electro-Magnetic Form Factors

One-photon approximation,  $\alpha_{em} = 1/137$ , hadron current

$$\mathcal{J}_{hadronic}^\mu = ie\bar{N}(p') \left[ \gamma^\mu F_1(Q^2) + \frac{i\sigma^{\mu\nu}q_\nu}{2M} F_2(Q^2) \right] N(p) \quad \text{Rosenbluth (1950)}$$

Full expression for  $\mathcal{M}$  has three complex functions,  $F_1, F_2, F_3$  Guichon & Vanderhaeghen

$$\mathcal{M} = \frac{4\pi\alpha}{Q^2} \bar{u}' \gamma_\mu u \cdot \bar{N}' \left( \tilde{F}_1 \gamma^\mu - \tilde{F}_2 [\gamma^\mu, \gamma^\nu] \frac{q_\nu}{4M} + \tilde{F}_3 K_\nu \gamma^\nu \frac{P^\mu}{M^2} \right) N \quad \text{Afanasev et al.}$$

$$\tilde{G}_M = \tilde{F}_1 + \tilde{F}_2 \quad \tilde{G}_E = \tilde{F}_1 - \tau \tilde{F}_2$$

$\tilde{F}_i$  are functions of  $(s - u)$  and  $t$

Blunden et al.

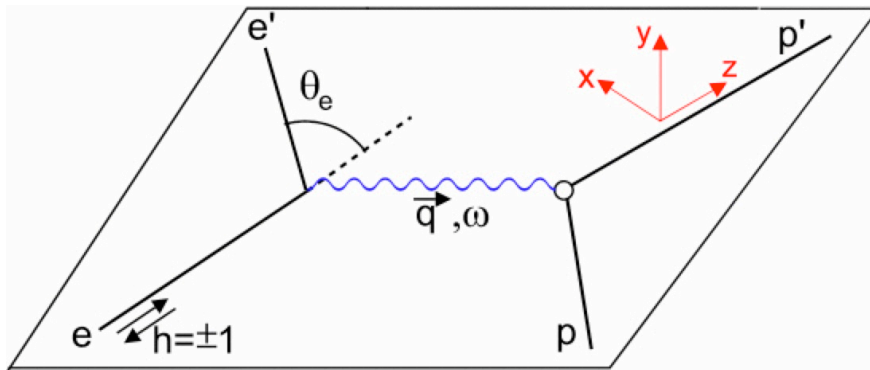
$$d\sigma = d\sigma_{NS} \left\{ \varepsilon \left( \tilde{G}_E + \frac{s-u}{4M^2} \tilde{F}_3 \right)^2 + \tau \left( \tilde{G}_M + \varepsilon \frac{s-u}{4M^2} \tilde{F}_3 \right)^2 \right\}$$

old  $G_{E,M}$  are real functions of  $t=-Q^2$

$$\sigma_R = \varepsilon G_E^2 + \tau G_M^2 + 2\tau G_M \operatorname{Re} \left( \delta \tilde{G}_M + \varepsilon \frac{s-u}{M^2} \tilde{F}_3 \right) + 2\varepsilon G_E \operatorname{Re} \left( \delta \tilde{G}_E + \frac{s-u}{M^2} \tilde{F}_3 \right)$$

Extra terms contribute less than few % to  $\sigma_R$

# Double Polarized Observables



$$N(\vec{e}, e' \vec{N})$$

Akhiezer et al., (1958)

Arnold et al., (1981)

Guichon &

Vanderhaeghen (2003)

Two-photon correction,  $\delta \sim 0.02$   
at typical values  $\varepsilon = 0.3$  –:–0.8

$$P_x = -2\sqrt{\tau(\tau + 1)}G_E^p G_M^p \tan \frac{\theta_{e'}}{2} / I_0$$

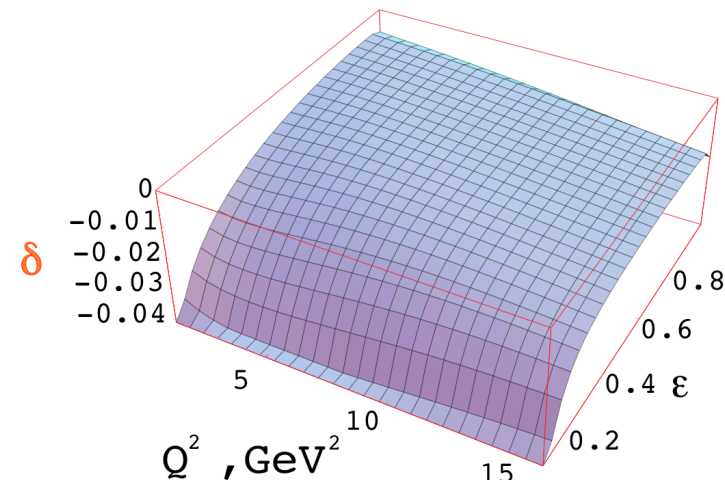
$$P_z = \frac{E_e + E_{e'}}{M_p} \sqrt{\tau(\tau + 1)} (G_M^p)^2 \tan^2 \frac{\theta_{e'}}{2} / I_0$$

$$I_0 \propto \epsilon (G_E^p)^2 + \tau (G_M^p)^2$$

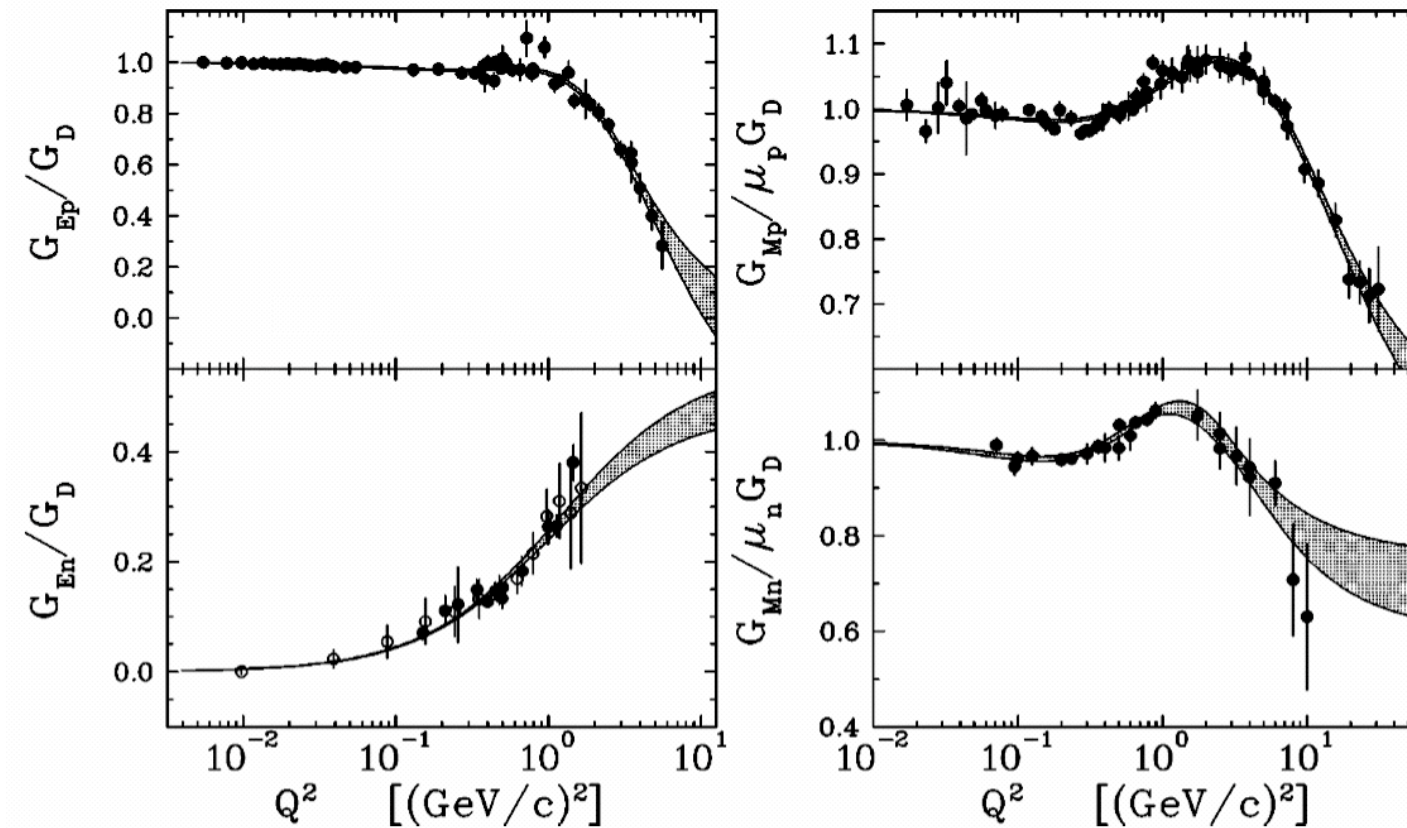
$$\frac{G_E^p}{G_M^p} \Big|_{1-\gamma} = -\frac{P_x}{P_z} \frac{E_e + E_{e'}}{2M_p} \tan(\theta_e/2)$$

$$\mu \frac{G_E^p}{G_M^p} \Big|_{1,2-\gamma} = \mu \frac{G_E^p}{G_M^p} \Big|_{1-\gamma} + \delta$$

Similar result for polarized target case



# Kelly's Parameterization

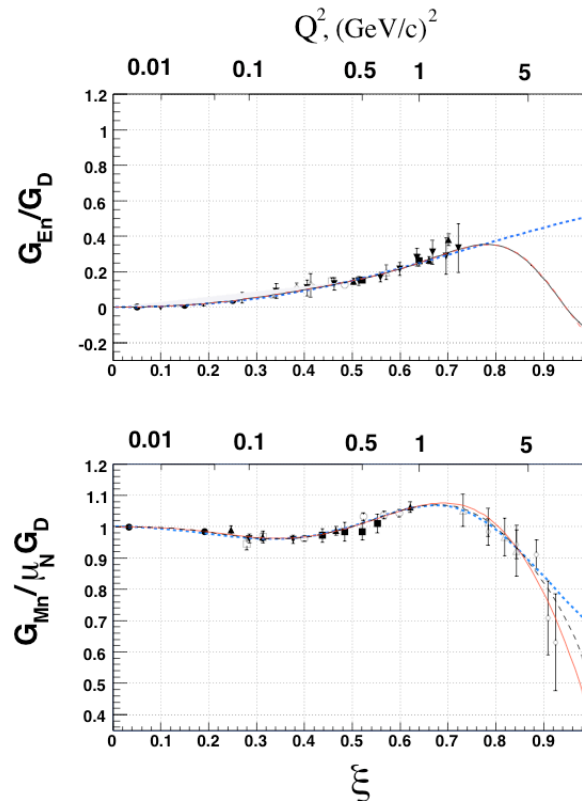
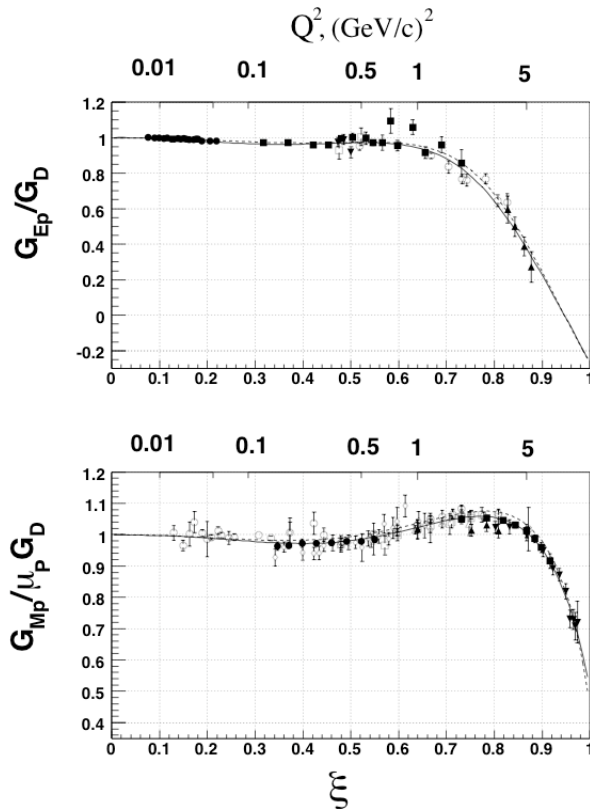


J. Kelly,  
PRC 70,  
068202  
(2004)

$$G(Q^2) = \sum_{k=0}^{n=1} a_k \tau^k / (1 + \sum_{k=1}^{n+2=3} b_k \tau^k)$$

scaling constraint:  $Q \rightarrow \infty, G \sim Q^{-4}$

# Duality constrained parameterization



Bodek,  
Avvakumov,  
Bradford, Budd  
arXiv:hep-ex  
0708.1946

$\xi$  is Nachtmann  
scaling variable

Two constraints  
QCD motivated

Kelly's

$(d/u) = 0, 0.2$

$$\xi^{p,n} = 2 / (1 + \sqrt{1 + \tau_{p,n}^{-1}})$$

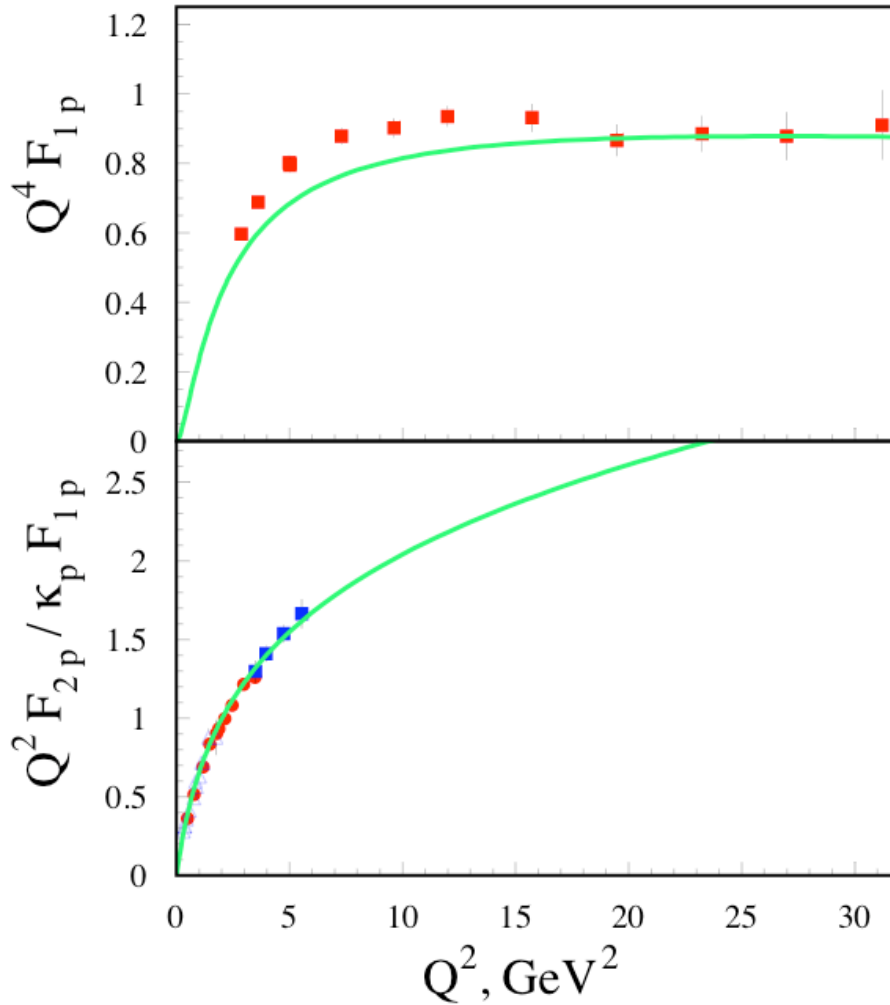
constrains: at  $\xi \rightarrow 1$

$$G_{BAB} = A(\xi) \times G_{Kelly}(Q^2)$$

$$1) \frac{G_{Mn}^2}{G_{Mp}^2} = \frac{1+4(d/u)}{4+(d/u)} \quad 2) \frac{G_{En}^2}{G_{Mn}^2} = \frac{G_{Ep}^2}{G_{Mp}^2}$$



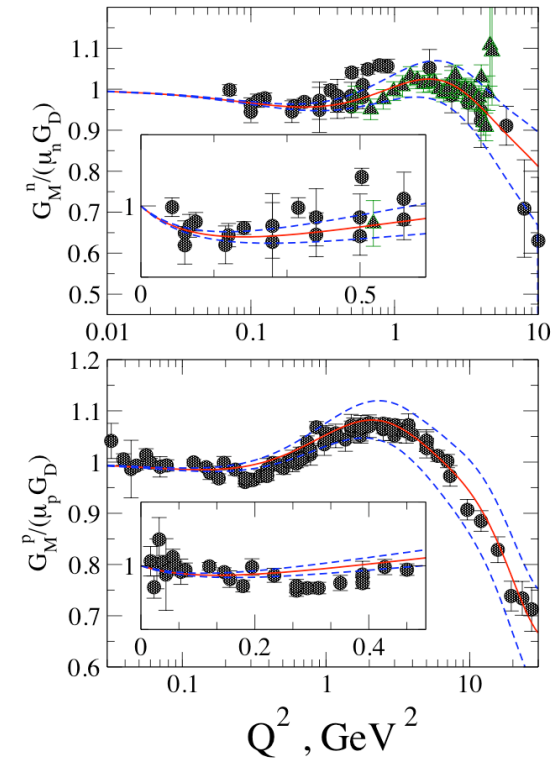
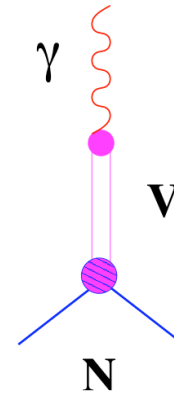
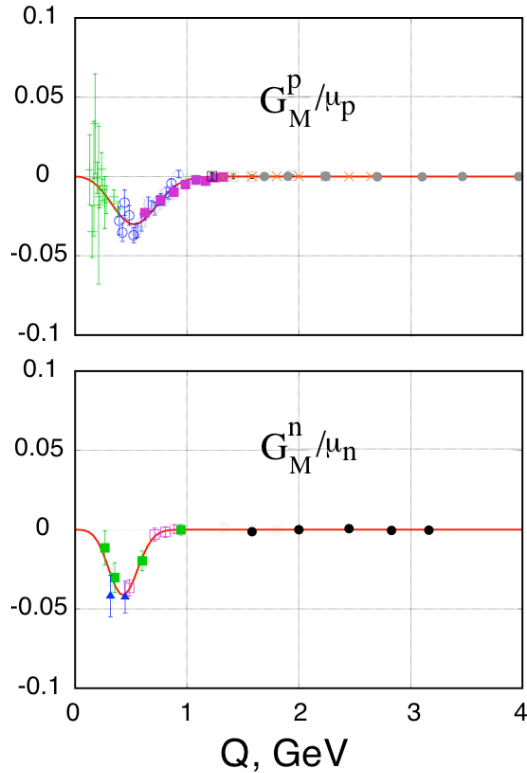
# FFs from GPDs



Guidal et al.,  
PRD 72, 054013,  
(2005)

FFs from GPDs and  
test of the scaling  
behavior of the  
proton FFs

# Form Factors at low $Q^2$



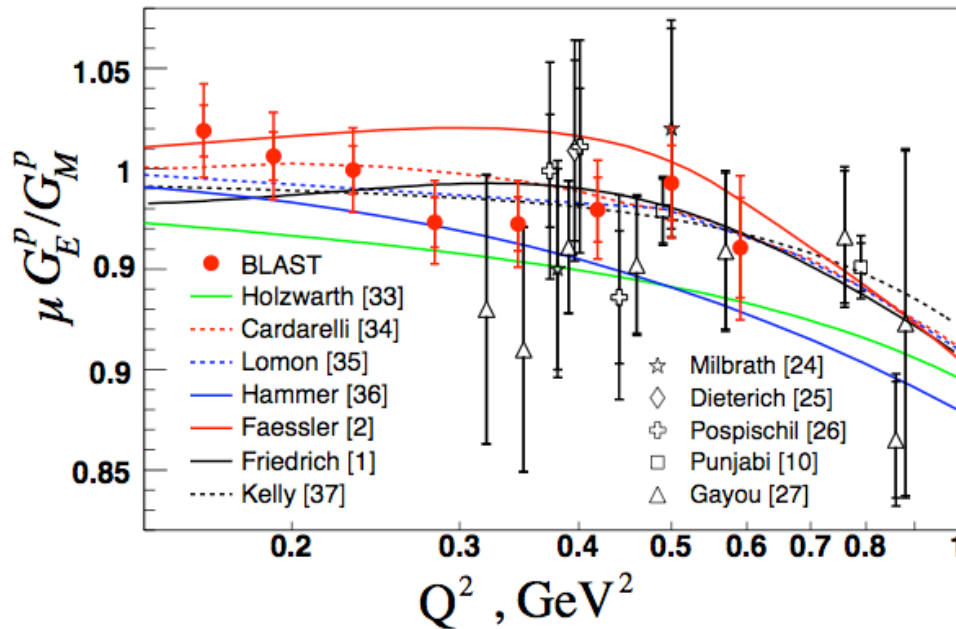
Friedrich&Walcher, EPJ (2003)  
have found a bump in all four FFs  
relatively to a two-dipoles fit

Belushkin et al, PRC (2007) get a good  
description of most data with **dispersion  
analysis** including meson continua

It will be very exciting to find deviation from DA fit => revived interest  
in precision experiments at  $Q^2$  range below 1  $\text{GeV}^2$

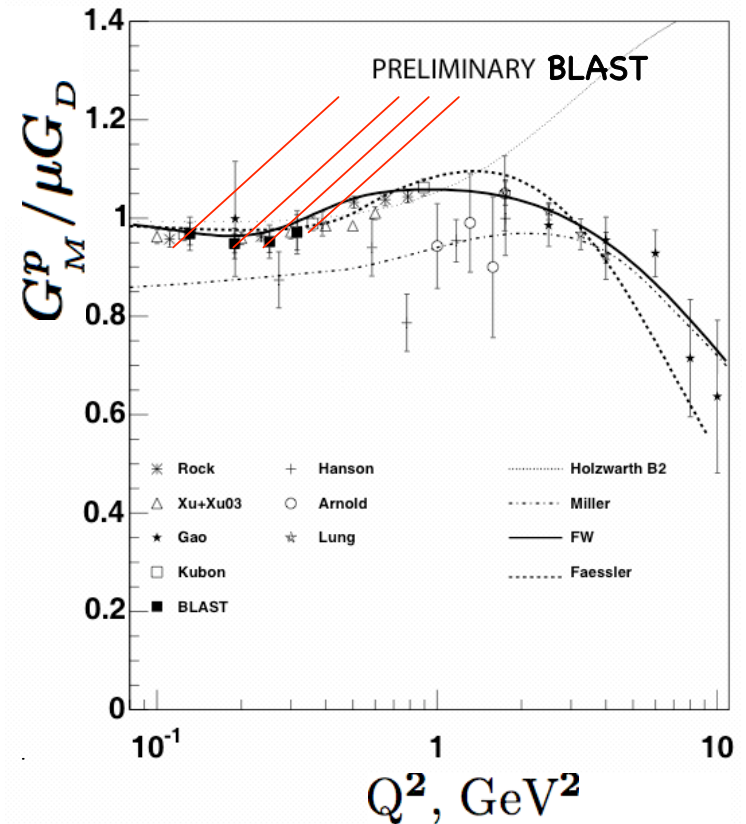
# Low $Q^2$ FFs from BLAST

Crawford et al., PRL 98, 052301 (2007)



Consistent with unity in 0.2-:-0.6  $\text{GeV}^2$

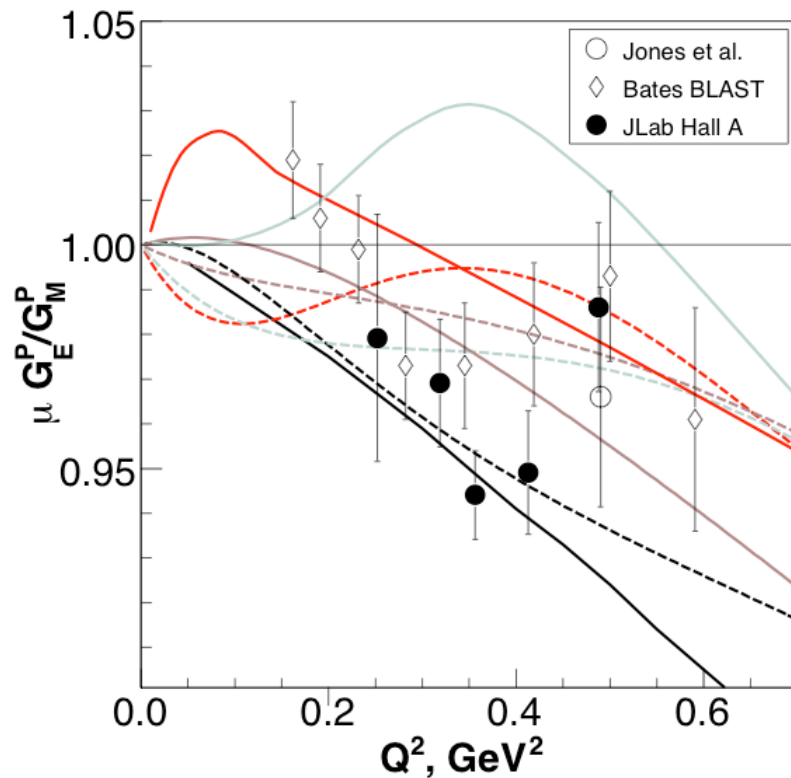
Alarcon, Eur.Phys.J. A32,477 (2007)



Show a dip at 0.2  $\text{GeV}^2$

# Low $Q^2$ GEP/GMP

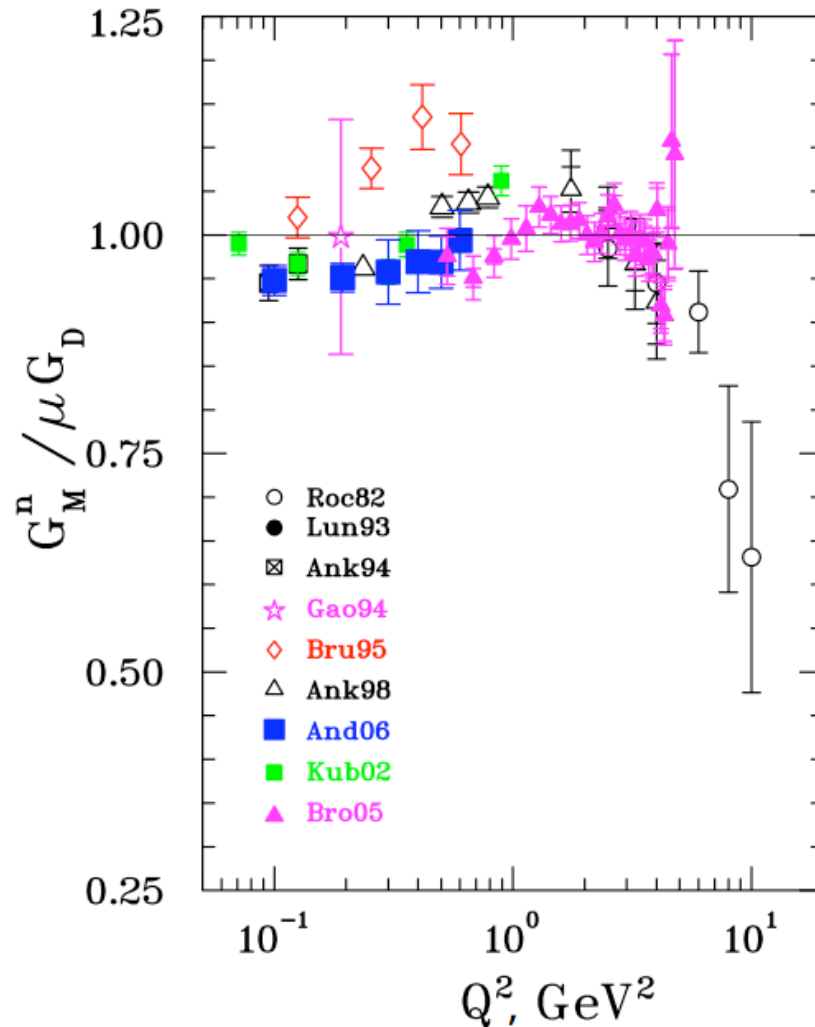
New polarization transfer data  
from JLab Hall A:  
G.Ron at al., arXiv:0706.0128



- J. Kelly Fit
- J. Arrington Fit
- Arrington & Sick Fit
- Friedrich & Walcher Fit
- G.A. Miller LFCBM
- Boffi et al. PFCCQM
- Belushkin, Hammer & Meißner VMD
- Faessler et al. LCQM

A ratio less than unity in range from 0.2 to 0.5  $\text{GeV}^2$

# Low $Q^2$ GMN inconsistency



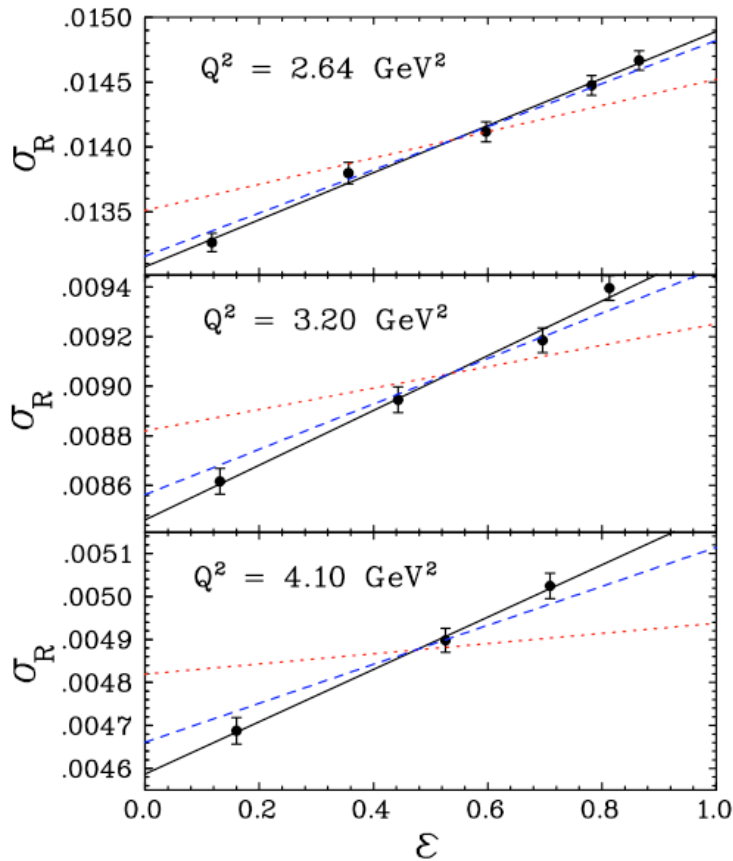
A systematic difference of several % between results (■ ■ ▲) in  $Q^2$  range  $0.4$  :-  $0.8 \text{ GeV}^2$ . A final analysis and paper from CLAS is coming soon.

Reminder that at least two independent experiments are always needed.

# Super ratio Rosenbluth

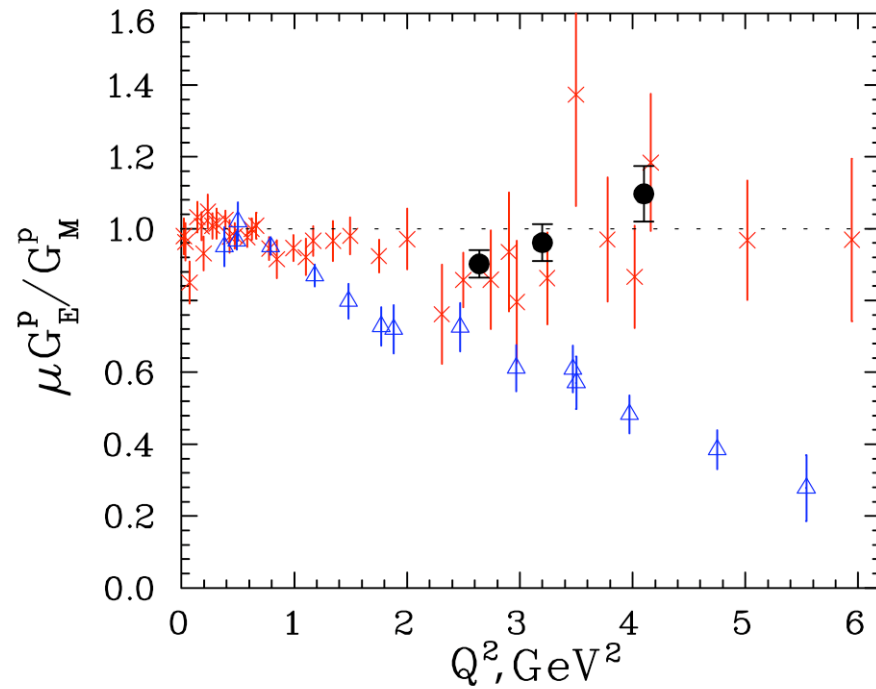
Arrington & Segel

$$\sigma_{\varepsilon 1} / \sigma_{\varepsilon 2} / \sigma_{\varepsilon 3} \dots$$



$$H(e, p)e'$$

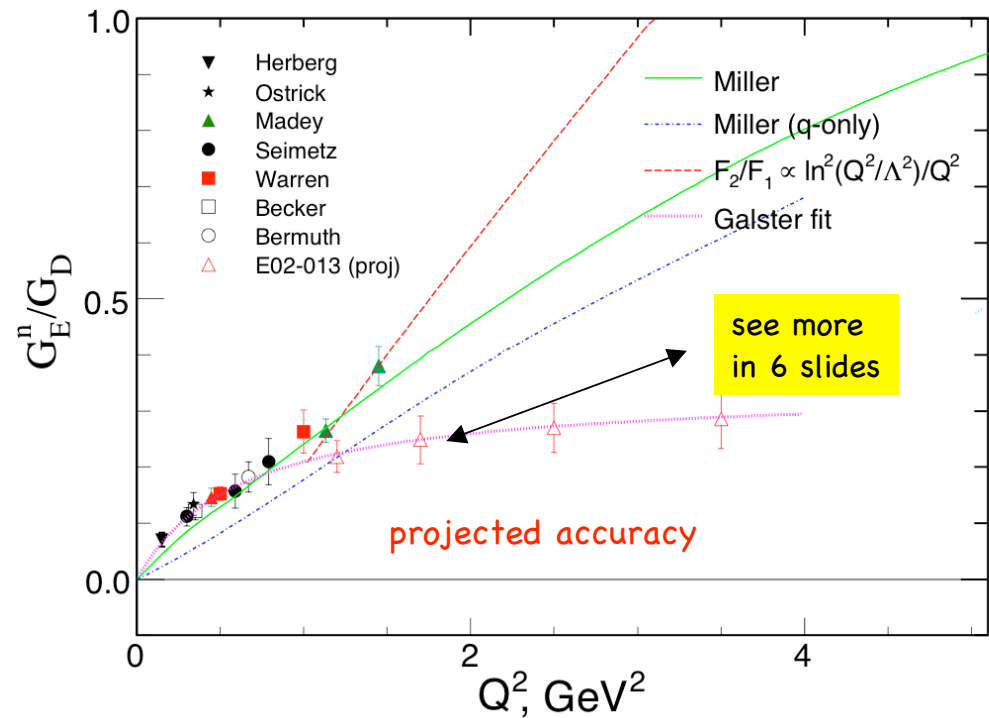
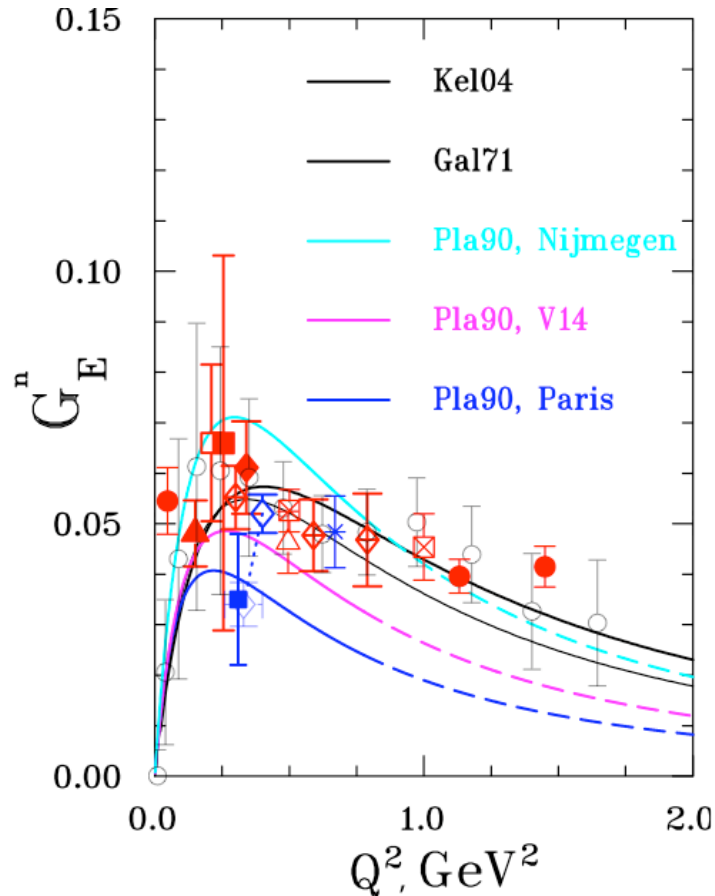
Qattan et al., PRL  
94, 142301 (2005)



New L/T experimental results consistent with old set

# Hall A GEN (E02-013)

G.Cates, N.Liyanage, and BW



Factor 2 increase of  $Q^2$  require  
x30 larger Figure-of-Merit

# High $Q^2$ GEN

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- ✓ Since 1984, when Blankleider&Woloshin suggested  ${}^3\vec{H}e(\vec{e}, e'n)$ , several experiments of this type have been performed at NIKHEF-K and Mainz (A1, A3) for  $Q^2$  up to  $0.7 \text{ GeV}^2$ , a big success in part due to **a new accurate 3-body calculation possible at low  $Q^2$**  (Glockle et al.)
- ✓ At  $Q^2$  **above 1-2  $\text{GeV}^2$**  Glauber method becomes sufficiently accurate (Sarksian)
- ✓ Electron-polarized neutron luminosity and high polarization of  ${}^3\text{He}$  target made **measurement about 10 times** more effective than with  $\text{ND}_3$ . In combination with a large acceptance electron spectrometer the total enhancement is **more than 100**, which allows to reach  $3.5 \text{ GeV}^2$

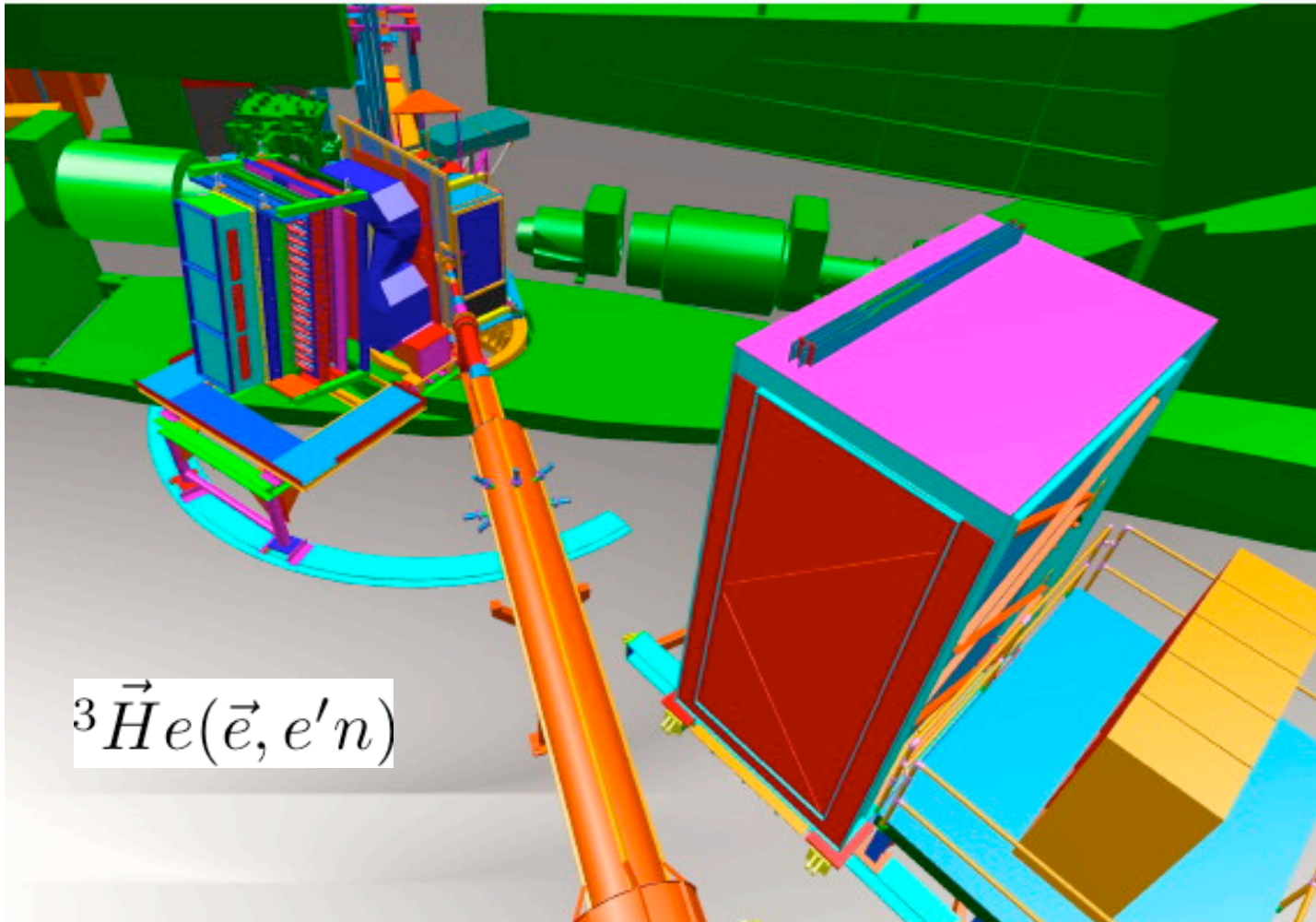
Require super

- Polarized target
- Electron spectrometer
- Neutron detector



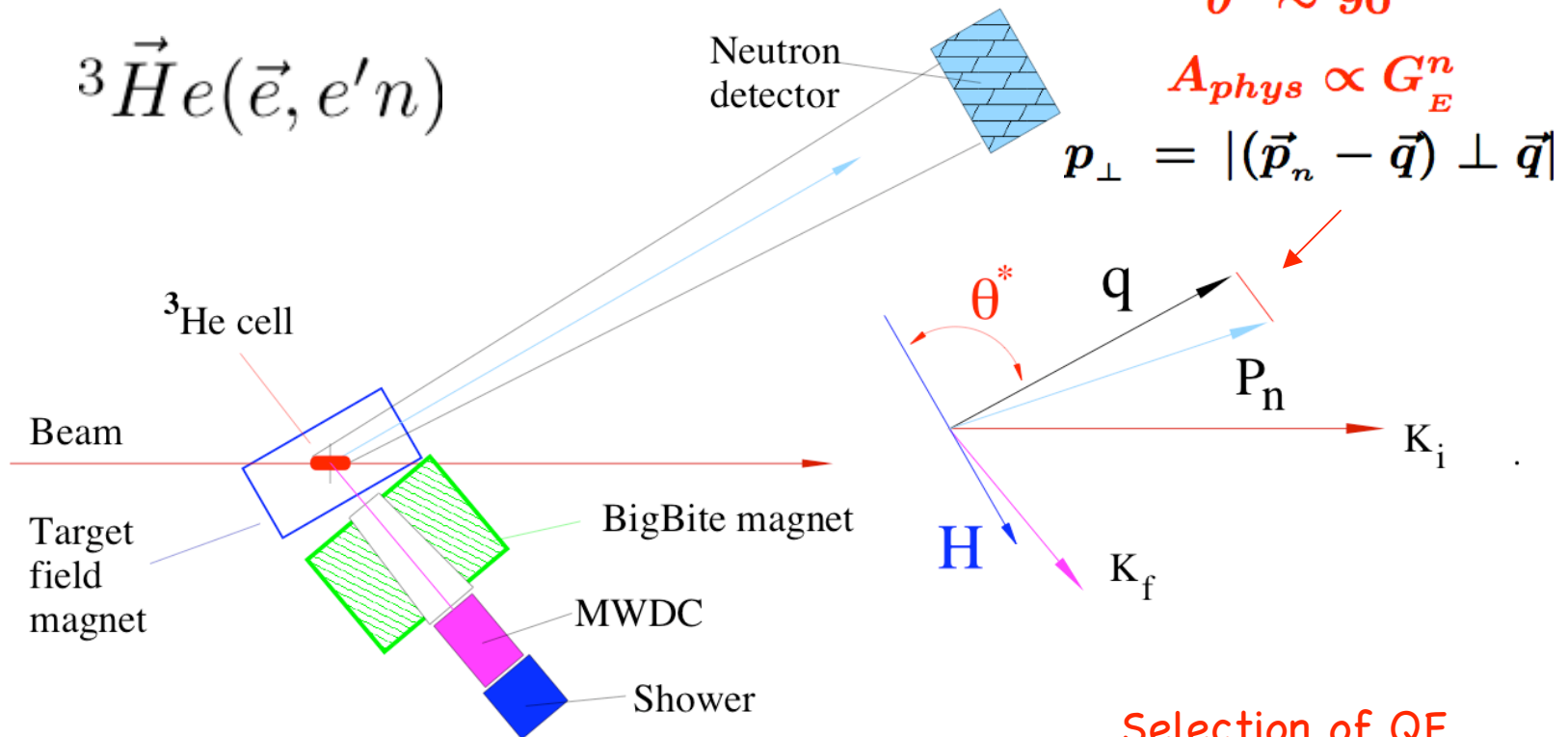
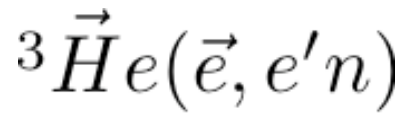
# Hall A GEN (E02-013)

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# E02-013 scheme

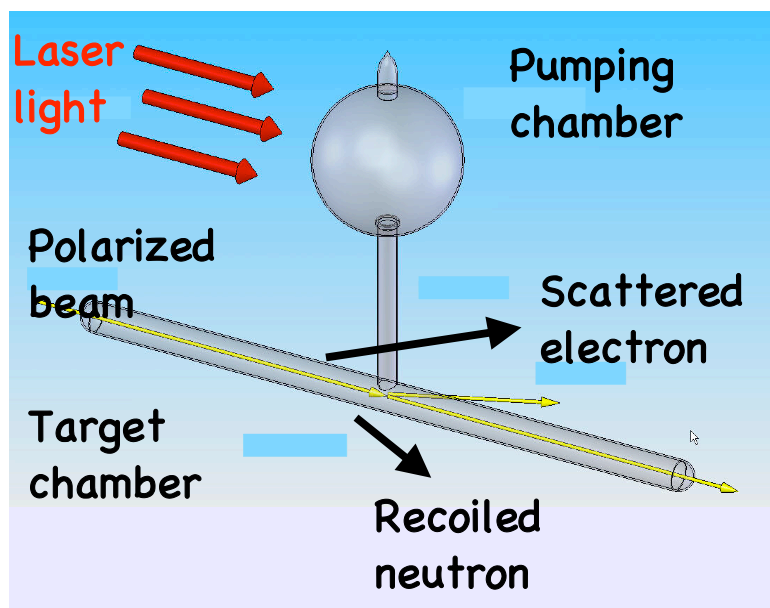
$$A_{phys} = -C1_{kin} \cdot G_E^n G_M^n \sin \theta^* \cos \phi^* - C2_{kin} \cdot (G_M^n)^2 \cos \theta^*$$



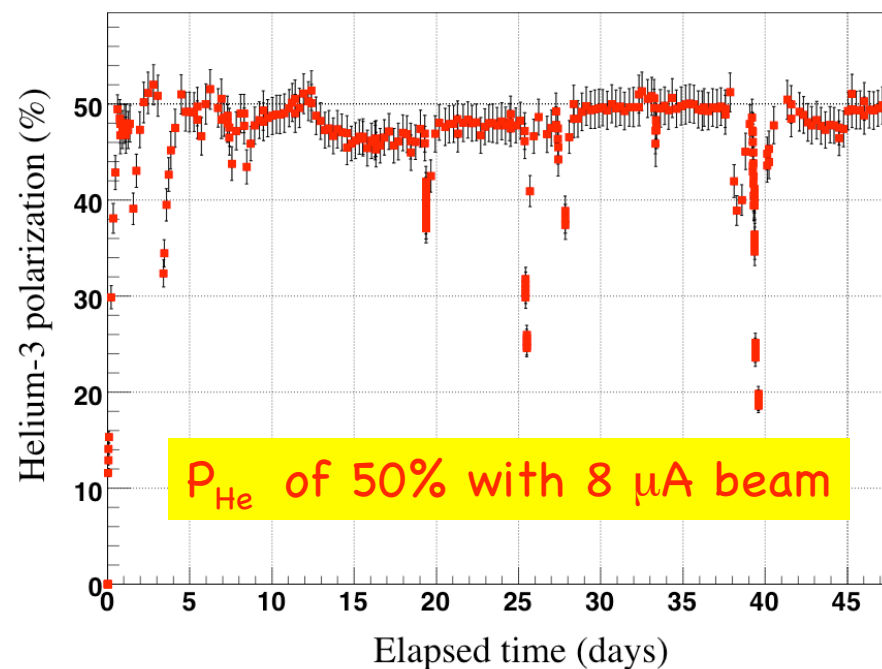
Selection of QE  
by cut  $P_{\perp} < 150$  MeV

# Polarized target

$${}^3\text{He} = p + p + n$$
$$S + S' + P \text{ waves}$$
$$P_n = 0.86 P_{\text{He}}$$



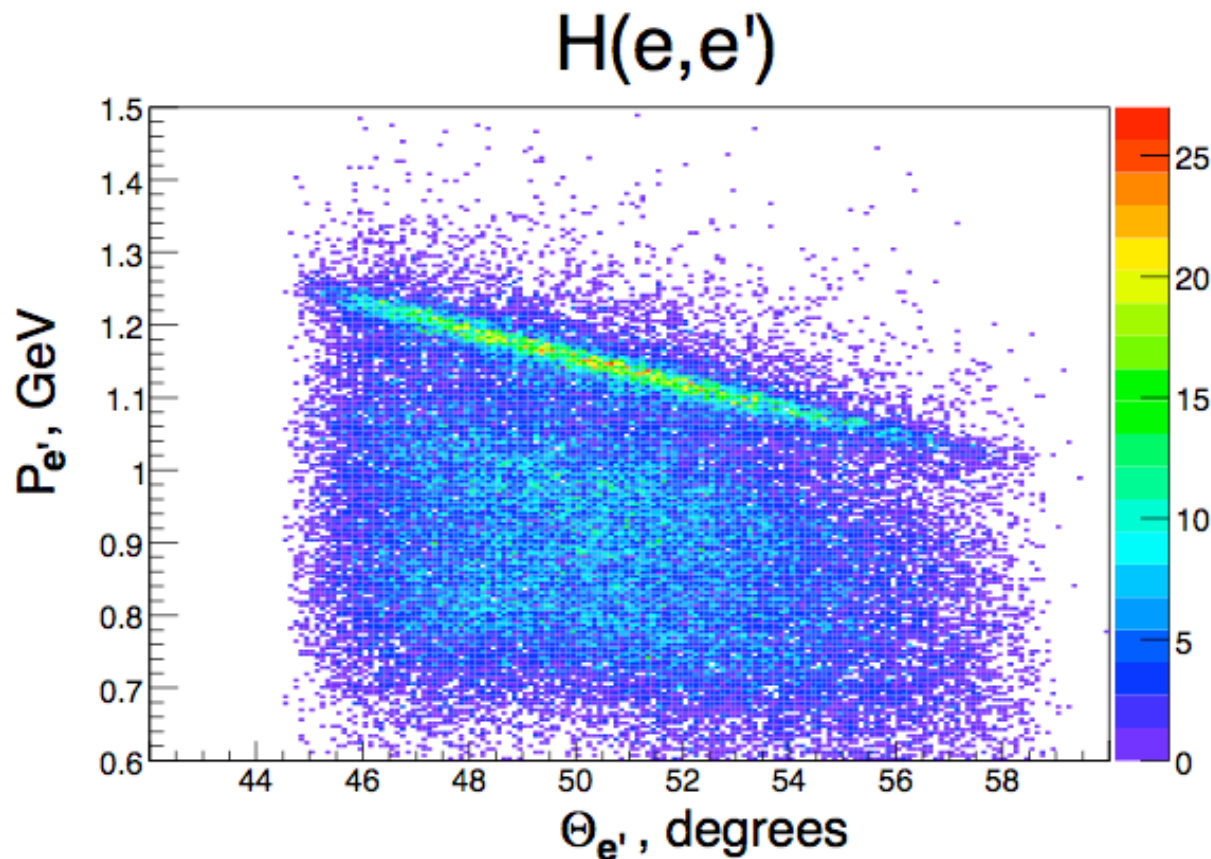
Polarization vs time for target cell "Edna"



Rb + K mixture has shortened spin-up time to 5-8 hours. The hybrid method of optical pumping was used here for the first time in the nuclear target.

# Electron Spectrometer

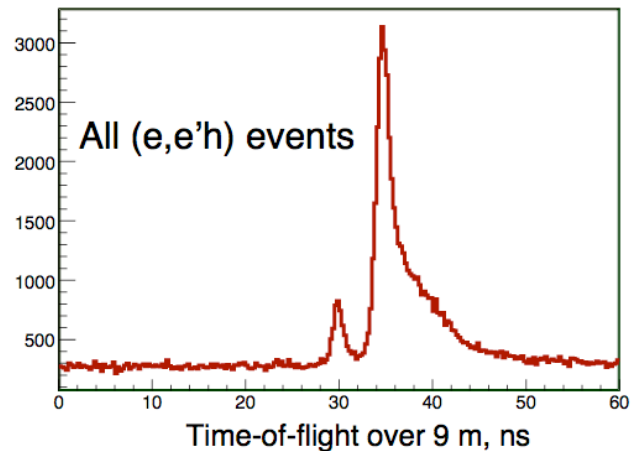
Useful  $\Delta Q^2/Q^2 \sim 0.1$  with  $\max \Omega$  leads to a large aspect ratio, limited just of  $30^\circ$  for the polar. target. BigBite was designed at NIKHEF for aspect ratio  $\Delta\theta/\Delta\phi = 1/5$ . Spectrometer has solid angle up to **95 msr**.



With luminosity of JLab polarized target,  $10^{37} \text{ cm}^{-2}/\text{s}$ , the open geometry - a dipole spectrometer - works well when all MWDCs located behind the magnet.

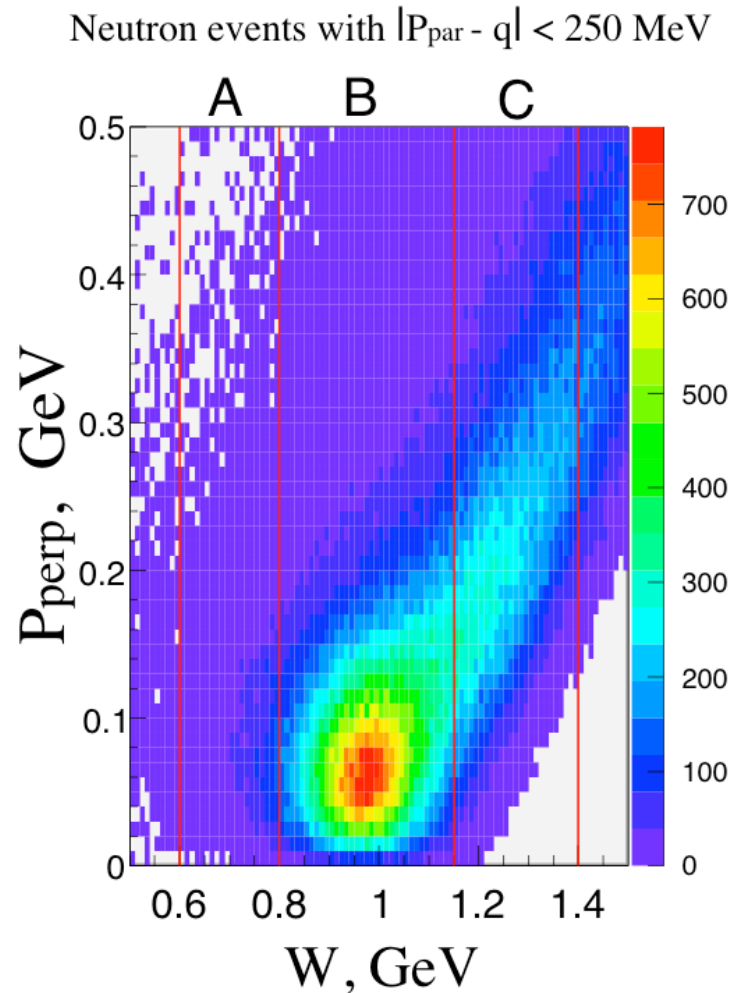
# Neutron Detector

- Match BigBite solid angle for QE kinematics
- Flight distance ~ 10 m
- Operation at  $3 \cdot 10^{37} \text{ cm}^2/\text{s}$
- $1.6 \times 5 \text{ m}^2$  active area
- 6-7 layers (~ 250 bars)
- 2 veto layers (~ 200)
- **0.38 ns time resolution**



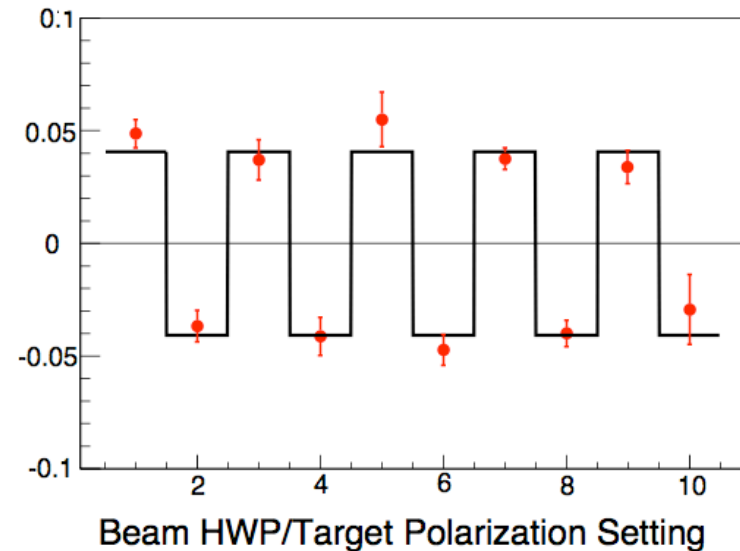


# Data analysis



Selection of QE ( $e, e'n$ ) events

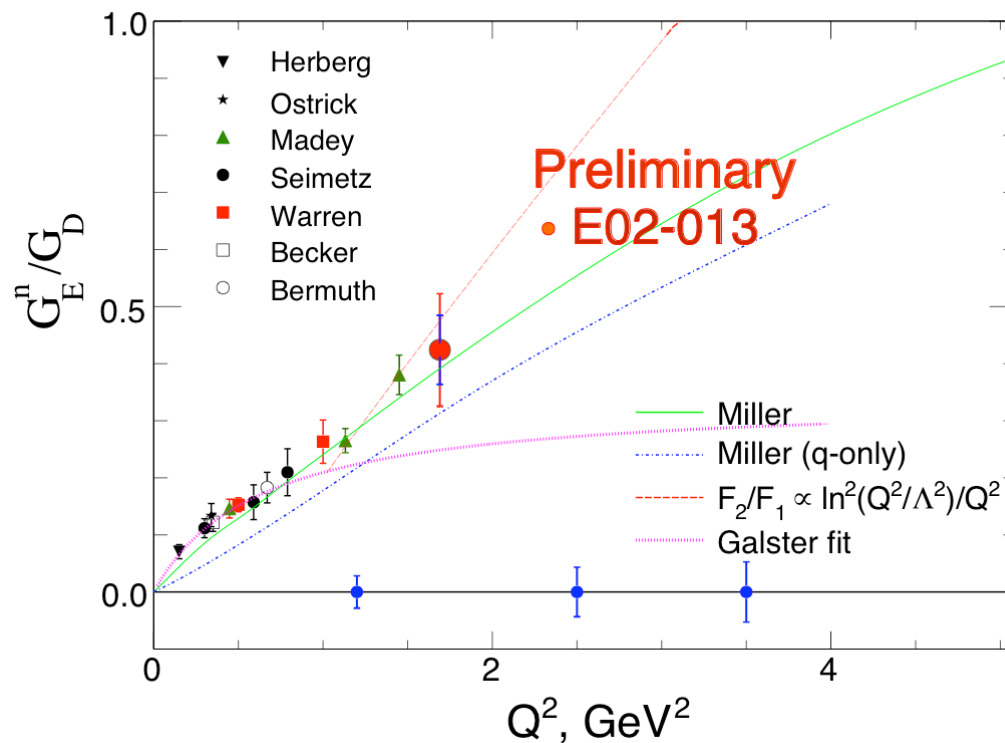
Observed Asymmetry for Quasi-elastic Neutrons



Asymmetry than corrected for

1. p-n misidentification
2. accidental events
3.  $A_{||}$  contribution
4. FSI for  $e, e'n$  process
5. Target, beam polarizations

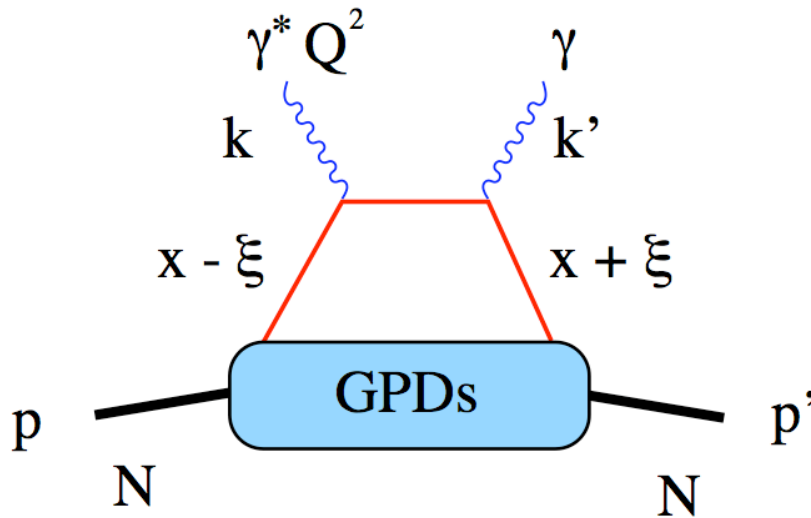
# First physics result from Hall A GEN



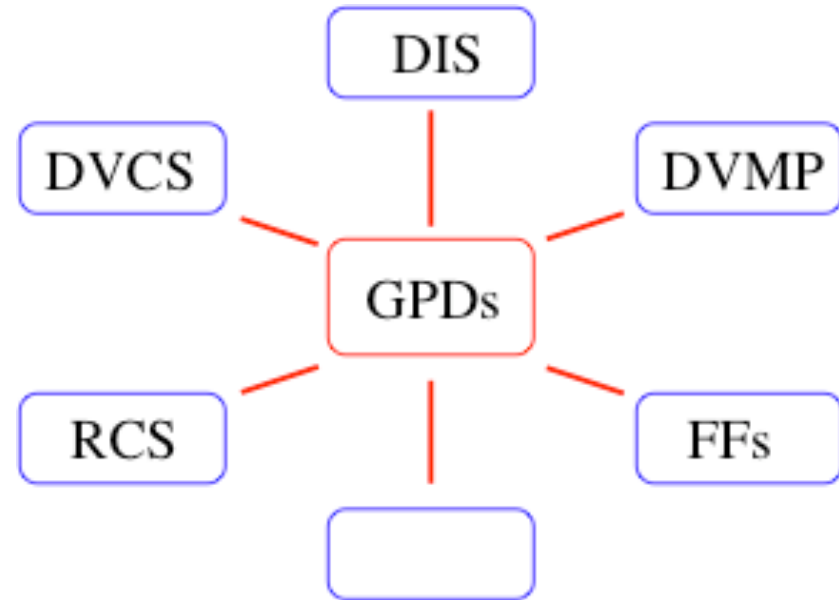
- Result is well above Galster.
- Nuclear corrections include neutron polarization and estimate of Glauber (~5%).
- Present error (~20%) dominated by preliminary “neutron dilution factor”, and is expected to be ~7% stat. and 8% syst. with further analysis.
- 3.4 GeV<sup>2</sup> result to be released in October DNP meeting at Newport News.

# GPDs of nucleon

Müller (94), Ji (97), Radyushkin (97)



where  $\xi = (p_q^+ - p_q'^+)/ (p_q^+ + p_q'^+)$



Quark dynamics of nucleon encoded in GPD functions

$H(x, \xi, t)$ ,  $\tilde{H}(x, \xi, t)$  hadron helicity-conserving; vector and axial-vector  
 $E(x, \xi, t)$ , and  $\tilde{E}(x, \xi, t)$  helicity-flipping; tensor and pseudo-scalar



# GPDs information

Reduction formulas at  $\xi = t = 0$   
for DIS and  $\xi = 0$  for FFs

$$H^q(x, \xi = 0, t = 0) = q(x)$$

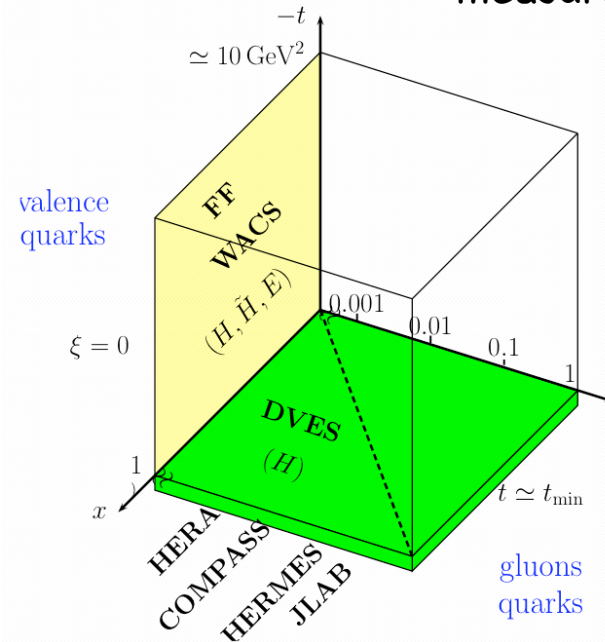
$$\tilde{H}^q(x, \xi = 0, t = 0) = \Delta q(x)$$

$$\int_{-1}^{+1} dx H^q(x, 0, Q^2) = F_1^q(Q^2)$$

$$\int_{-1}^{+1} dx E^q(x, 0, Q^2) = F_2^q(Q^2)$$

P.Kroll, Excl.-07

a lot to  
measure



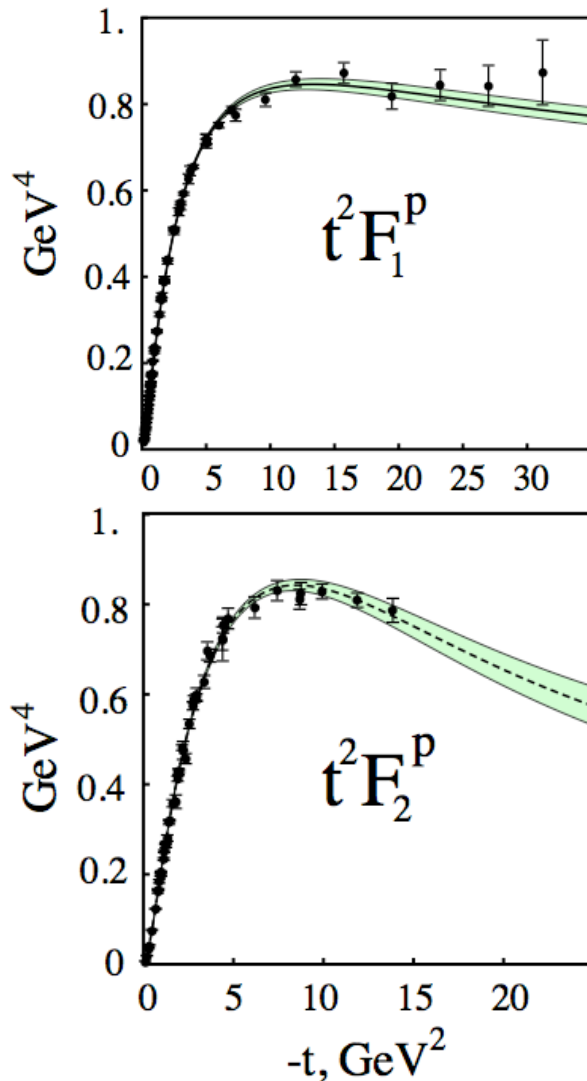
Ji's sum rule for quark orbital momentum

$$\langle L_v^q \rangle = \frac{1}{2} \int_0^1 dx [x E_v^q(x, \xi = 0, t = 0) + x q_v(x) - \Delta q_v(x)]$$

DVCS will access low  $t$ , large  $Q^2$  kinematics

FFs presently are the main source for  $E_v^q$

# Model of GPD and Form Factors



Diehl et al (2005),  
Guidal et al (2005)

use **all** available data on  $G_M^p, G_M^n, G_E^p, G_E^n, F_A$

and CTEQ6 parametrization of  $q(x), \Delta q(x)$

in order to determine  $H_v^{u,d}, \tilde{H}_v^{u,d}, E_v^{u,d}$

ANSATZ:  $H_v^q(x, t) = q_v(x) \exp [f_q(x)t]$

$$f_q = [\alpha' \log(1/x) + B_q] (1-x)^{n+1} + A_q x (1-x)^n$$

fixed  $\alpha' = 0.9 \text{ GeV}^{-2}, n = 1, 2;$

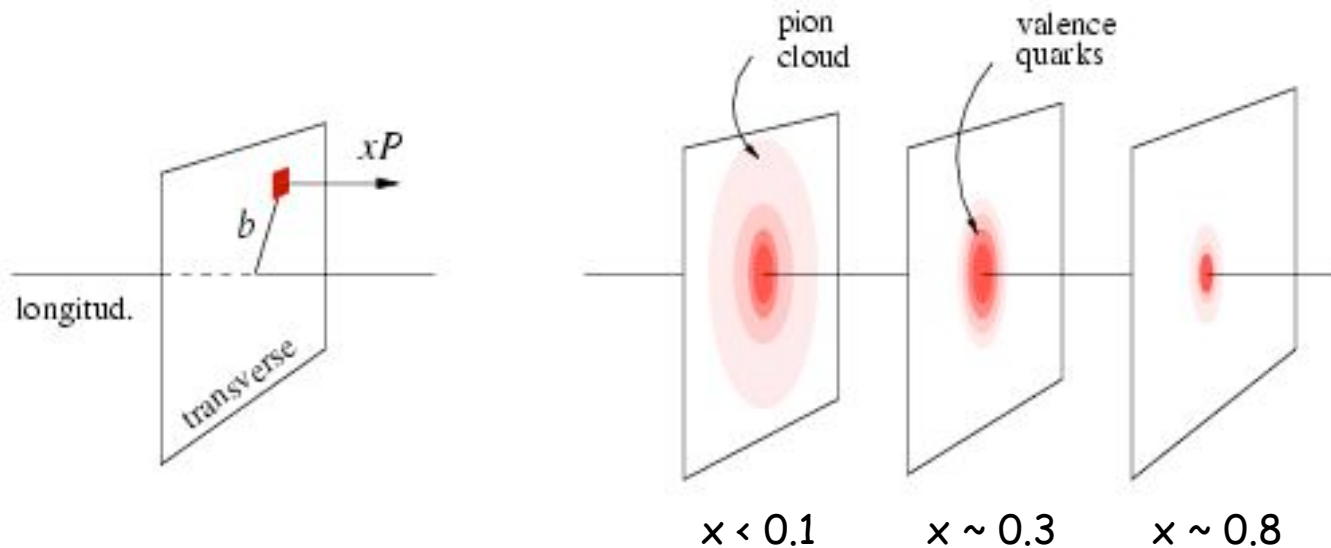
only  $A_q$ , and  $B_q$  parameters are fitted

# GPDs and impact parameter

Transverse momentum invariance allows frame independent Fourier transform from  $H(x, \xi, t)$  to  $q(x, \xi, b)$  with impact parameter  $b$  defined relative to center of momentum

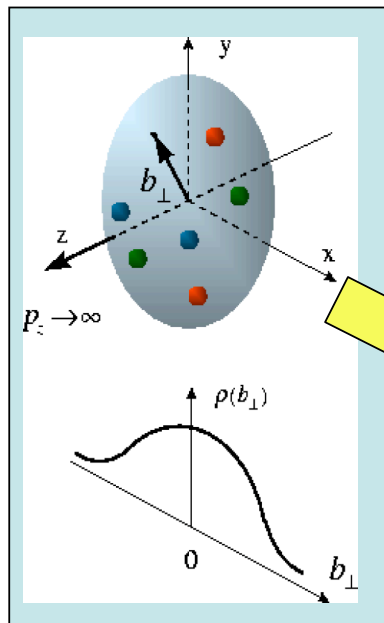
Burkardt, Int. J. Mod. Phys. A 18, 173 (2003)

Fourier transform in momentum transfer

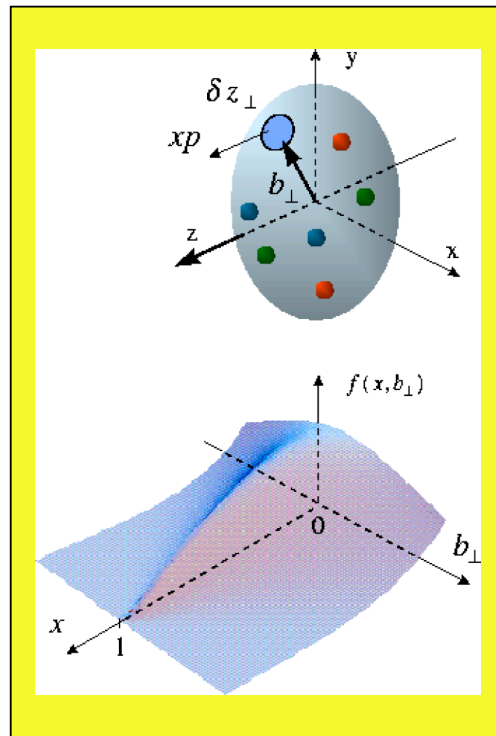


gives transverse size of quark (parton) with longitud. momentum fraction  $x$

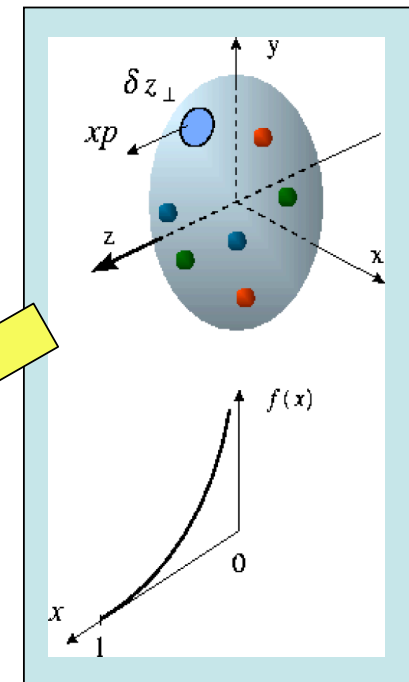
# 3-d picture of the nucleon



Proton form factors,  
transverse charge &  
current densities



Correlated quark momentum  
and helicity distributions in  
transverse space - GPDs



Structure functions,  
quark longitudinal  
momentum & helicity  
distributions

# Nucleon Density from GPD

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$$F_1(t) = \sum_q e_q \int dx H_q(x, t)$$

Muller, Ji, Radyushkin

$$q(x, \mathbf{b}) = \int \frac{d^2q}{(2\pi)^2} e^{i \mathbf{q} \cdot \mathbf{b}} H_q(x, t = -q^2)$$

M. Burkardt

$$\rho(\mathbf{b}) \equiv \sum_q e_q \int dx q(x, \mathbf{b}) = \int d^2q F_1(q^2) e^{i \mathbf{q} \cdot \mathbf{b}}$$

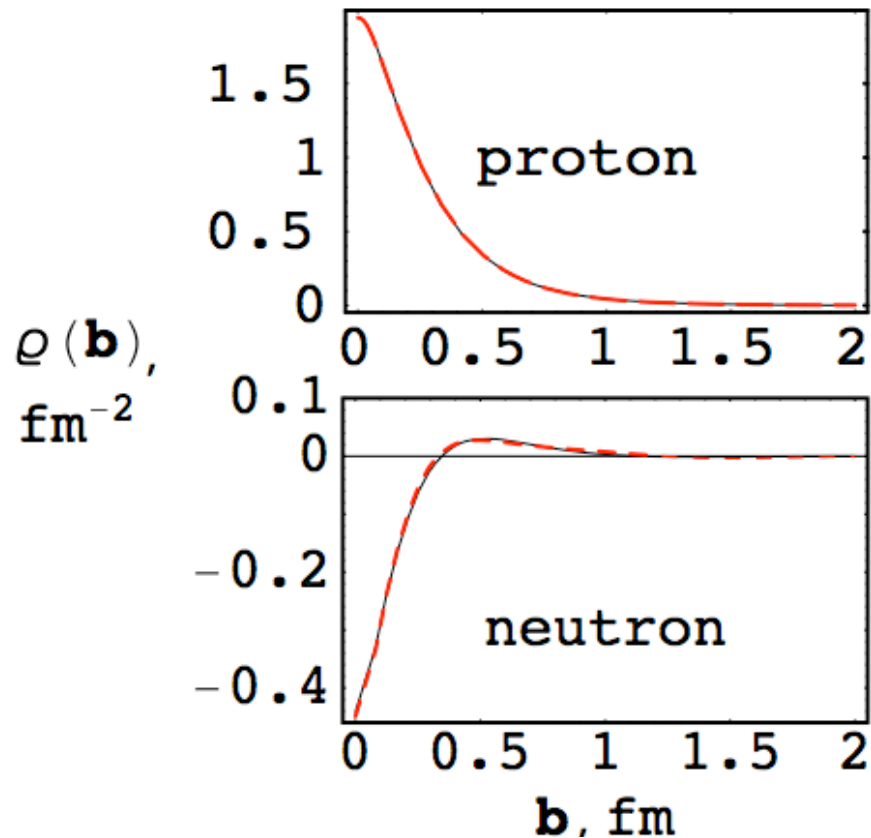
$$\rho(\mathbf{b}) = \int_0^\infty \frac{Q \cdot dQ}{2\pi} J_0(Qb) \frac{G_E(Q^2) + \tau G_M(Q^2)}{1 + \tau}$$

G. Miller,  
arXiv:0705.2409

center of momentum  $\mathbf{R}_\perp = \sum_i \mathbf{x}_i \cdot \mathbf{r}_{\perp, i}$

$\mathbf{b}$  is defined relative to  $\mathbf{R}_\perp$

# Neutron is negative inside



G.Miller, arXiv:0705.2409

using FFs from

Kelly's fit

BBBA's fit

Does it contradict intuition ?

Static picture: a neutron is  
a proton in the center plus  $\pi^-$

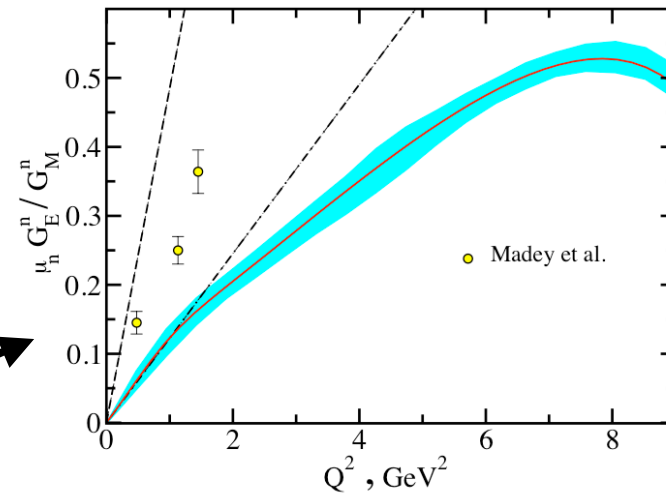
Intuitive picture of static  
constituent quarks is not  
applicable for large  $Q^2$   
where quark DFs play role

- Negative density at low  $\mathbf{b}$  in a neutron => d quarks dominate
- High  $Q^2$  elastic process in Feynman mechanism requires a large  $x$  quark, so d quarks dominate at large  $x$ , in agreement with DIS

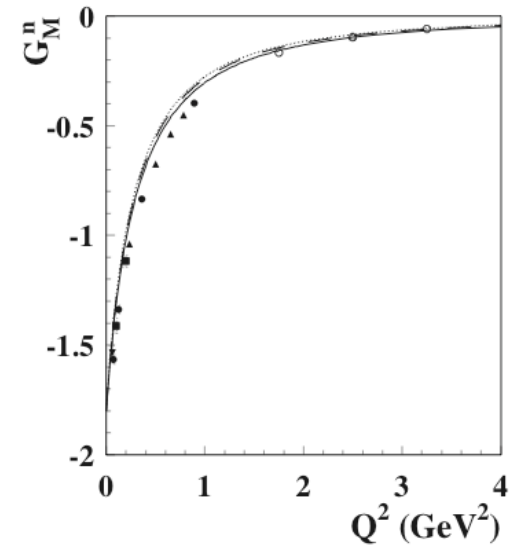
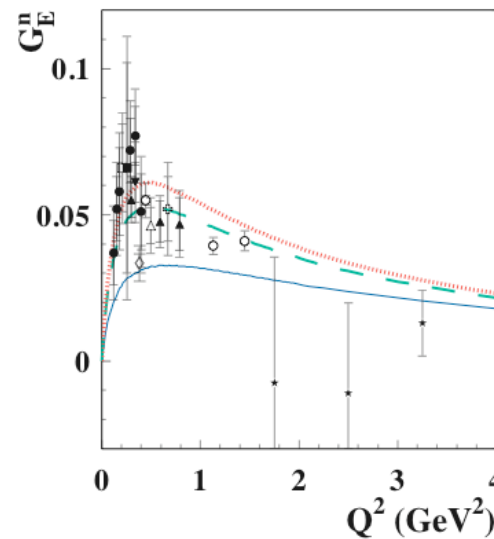
# Nucleon models

See recent review Arrington, Roberts, Zanotti: [nucl-th/0611050](https://arxiv.org/abs/nucl-th/0611050)

- Miller, since 2002, Spin densities in LF-Cloudy-Bag-Model lead to **non-spherical shape of the proton**.
- Bhagwat et al., arXiv:nucl-th/0610080, the Dyson-Schwinger equations used to calculate FFs



- Pasquini&Boffi, arXiv:0707.2897 developed nucleon-meson LFQCM and calculated many features of electroweak structure of the nucleon. GEN is sensitive to the **% of S' component**, find negative charge in the neutron at small impact parameter **b**



# GEP-III: $G_E^p/G_M^p$ for $8.6 \text{ GeV}^2$

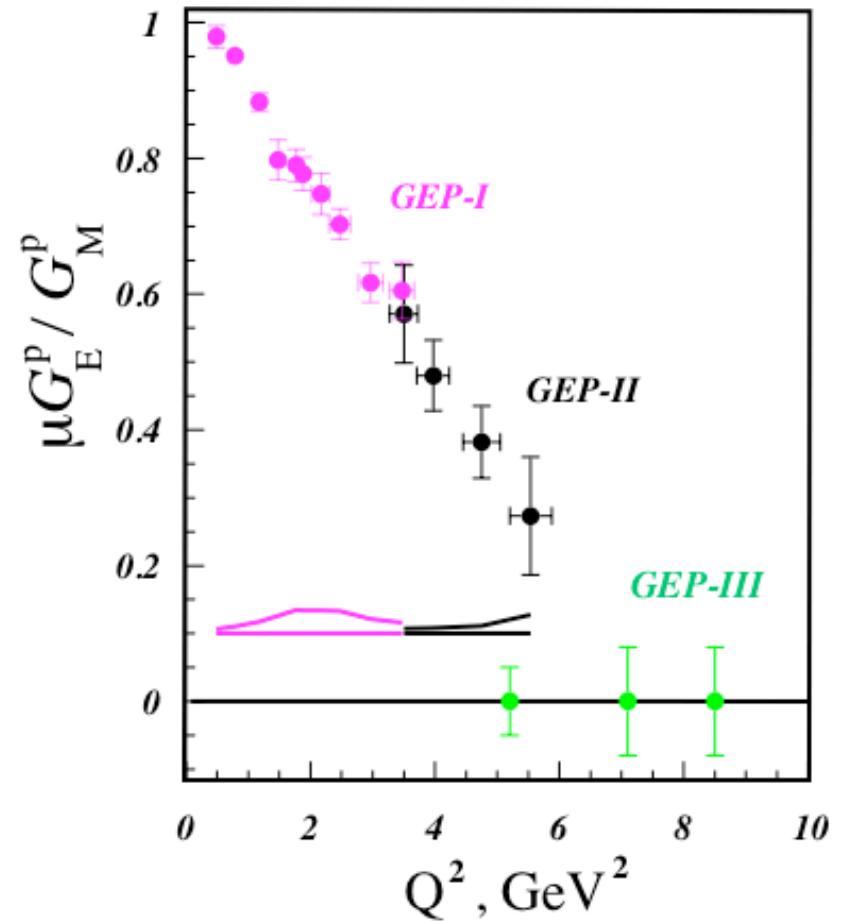
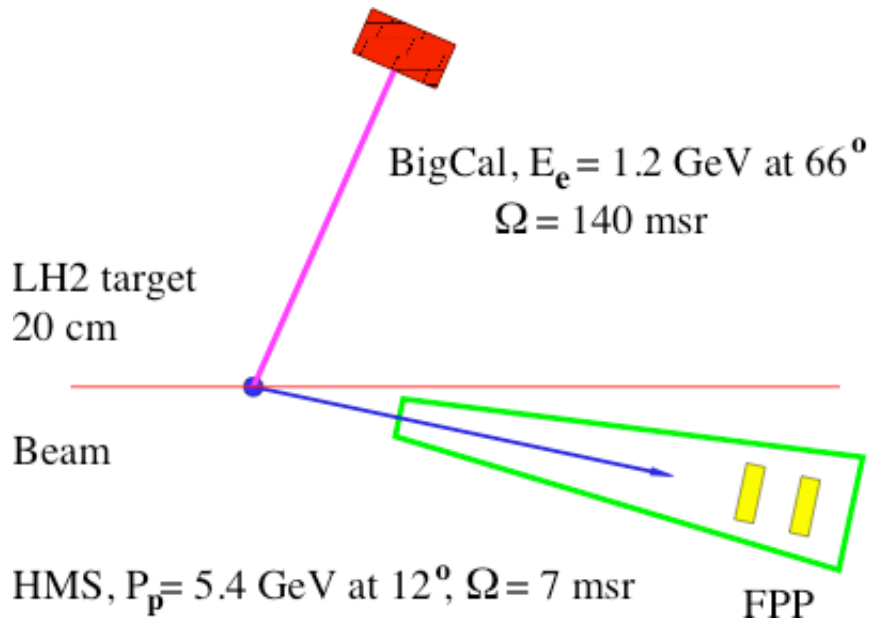
Brash, Jones, Perdrisat, Punjabi

Polarization transfer in  $H(\vec{e}, e'\vec{p})$

- New detector for scattered electron
- New polarimeter for recoiled proton

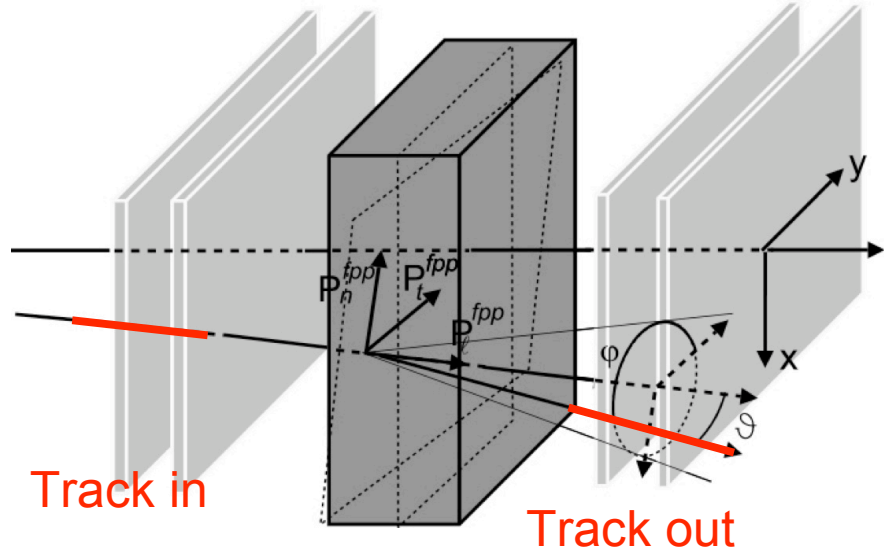
Commissioning will start next month

High  $Q^2$  data will be taken by 5/08





# Focal Plane Polarimeter

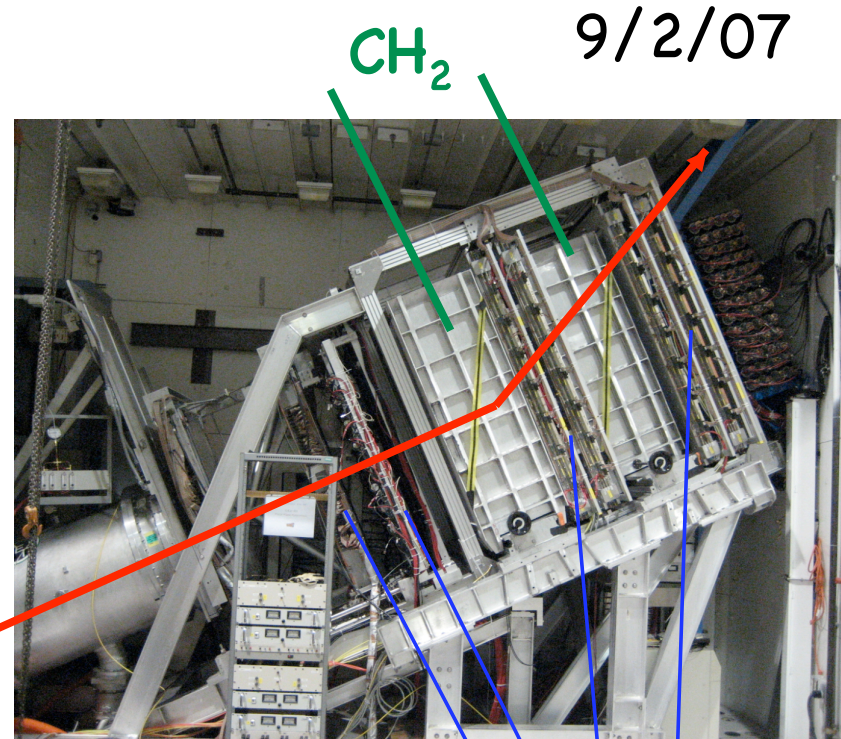


$$f^{\pm}(\vartheta, \varphi) = \frac{\epsilon(\vartheta, \varphi)}{2\pi} \left[ 1 \pm A_y (P_x^{fpp} \sin \varphi - P_y^{fpp} \cos \varphi) \right]$$

where  $\pm$  refers to electron beam helicity

$$A = \frac{f^+ - f^-}{f^+ + f^-} = A_y \left( P_x^{fpp} \sin \varphi - P_y^{fpp} \cos \varphi \right)$$

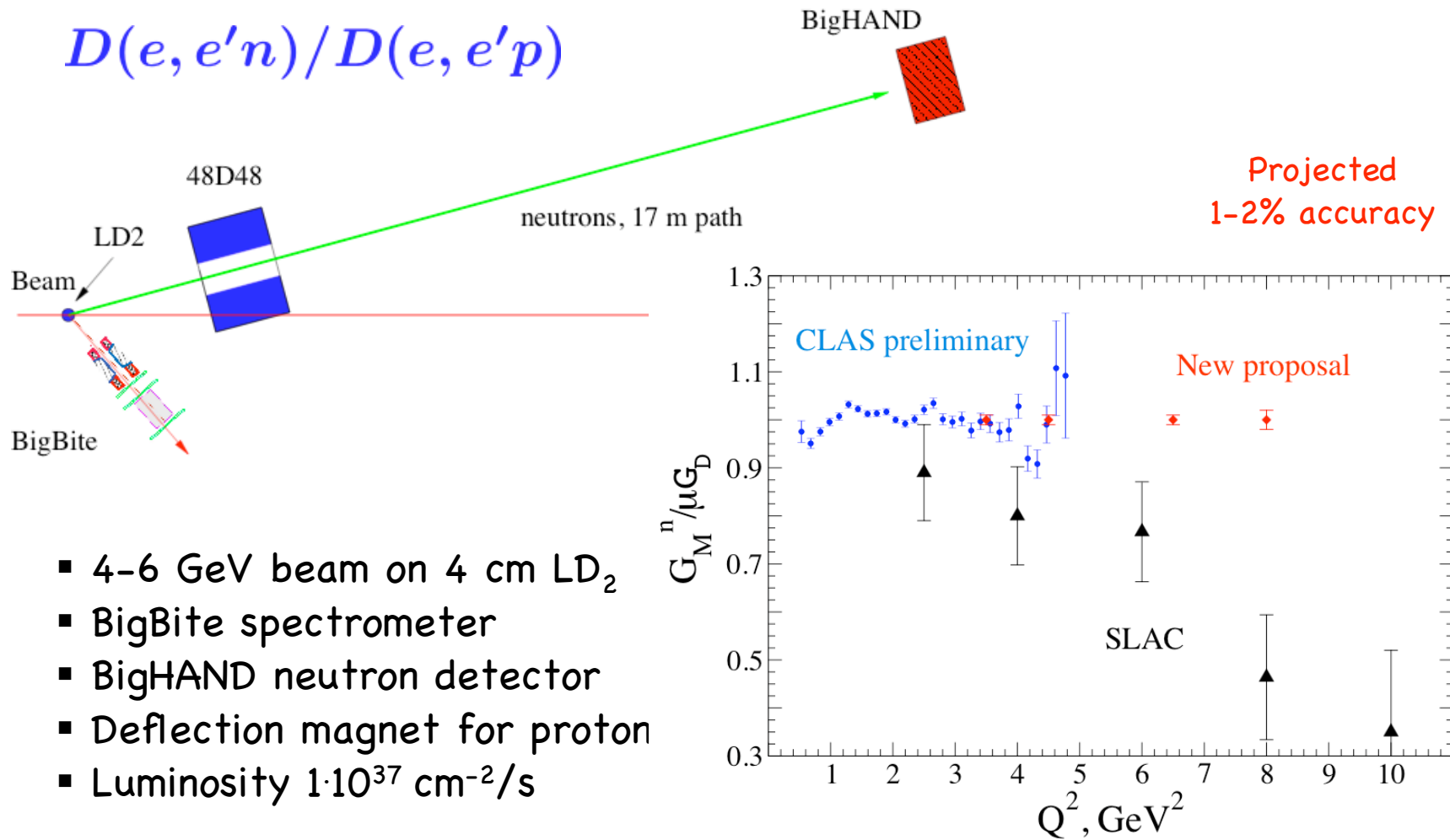
$$\mu_p \frac{G_E^p}{G_M^p} = -\mu_p \frac{E_e + E'_e}{2M_p} \tan \frac{\theta_e}{2} \left( \frac{P_x^{fpp}}{P_y^{fpp}} \sin \chi_{\theta} + \gamma_p (\mu_p - 1) \Delta \phi \right)$$



MWDCs

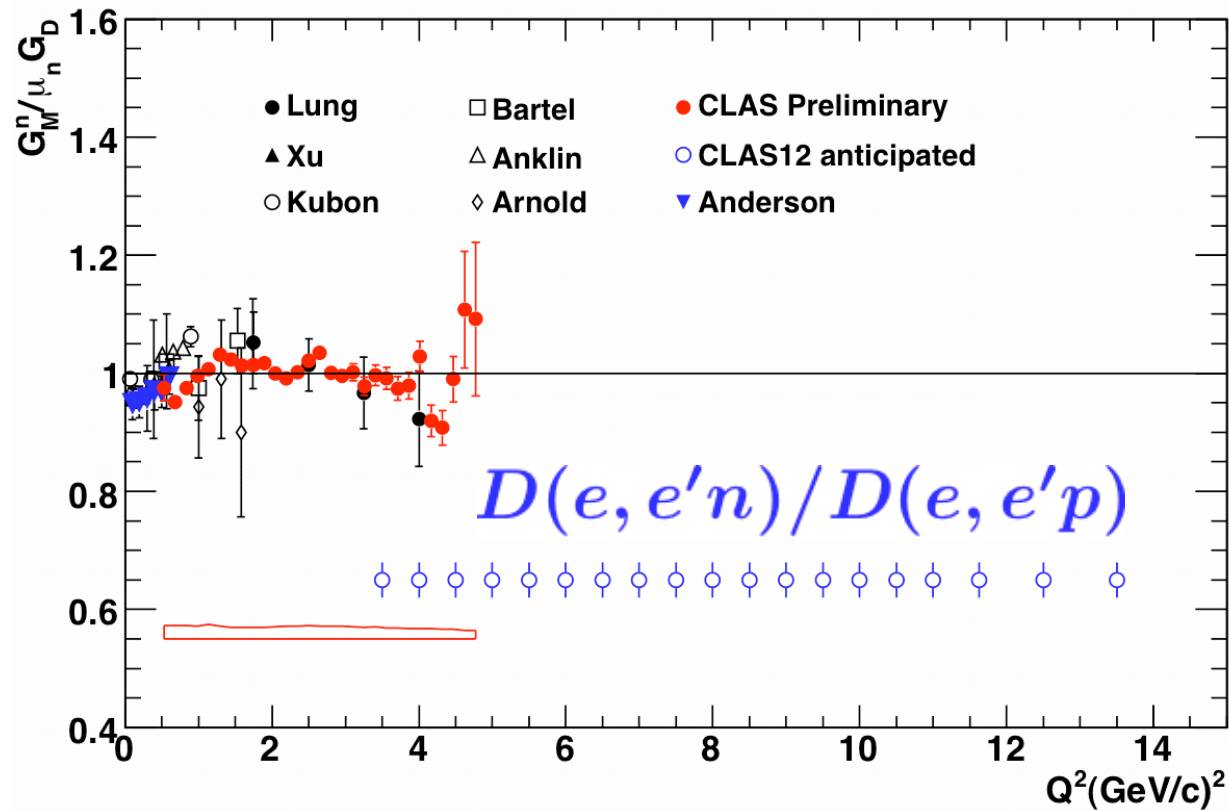
# Near future: GMN-8 (for PAC33)

B.Quinn and BW



# GMN-14 with CLAS++

Gilfoyle, Brooks, Vineyard, Hafidi, Lachniet

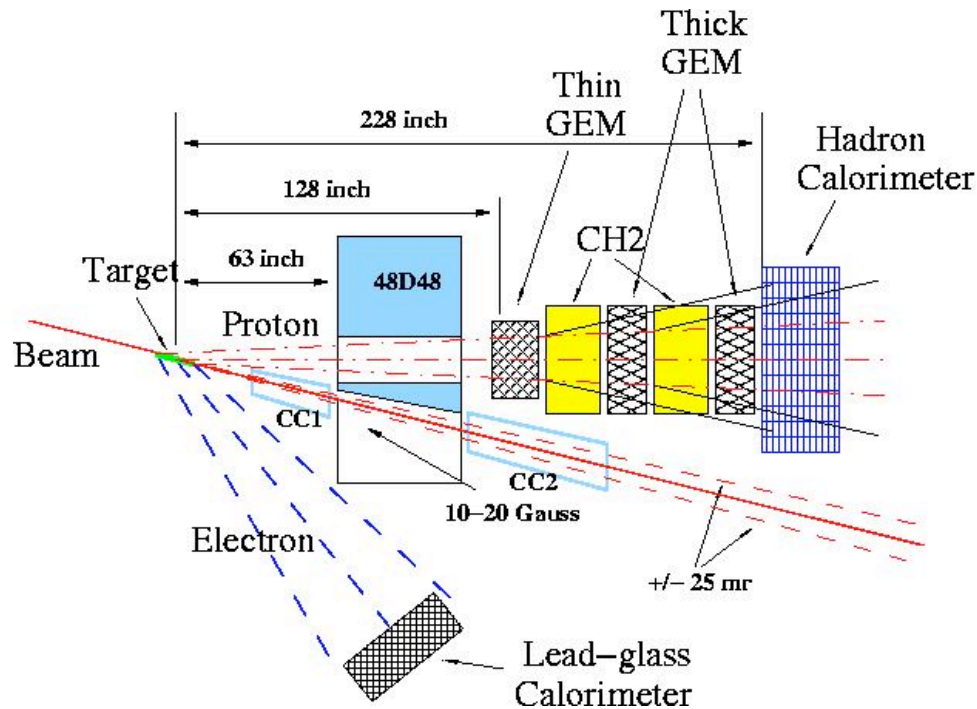


approved by PAC32  
for 12 GeV program

# GEP-15: $G_E^p/G_M^p$ up to $15 \text{ GeV}^2$

Perdrisat, Pentchev, Cisbani, Punjabi, BW

$$H(\vec{e}, e' \vec{p})$$



Beam:  $75 \mu\text{A}$ , 85% polarization

Target is 40 cm liquid  $\text{H}_2$

Electron arm at  $37^\circ$ , covers

$$Q^2 = 12.5 \text{ to } 16 \text{ GeV}^2$$

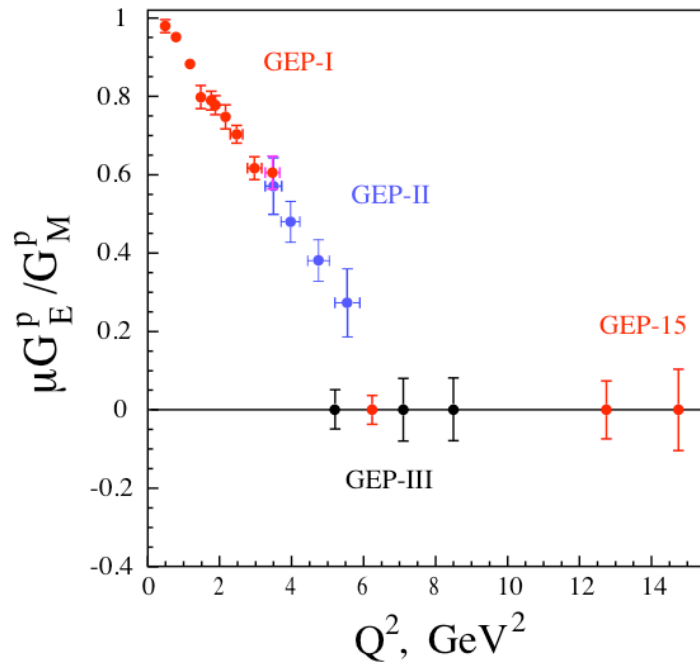
Proton arm at  $14^\circ$ ,  $\Omega \sim 35 \text{ msr}$

58 days of production time  
resulting accuracy:

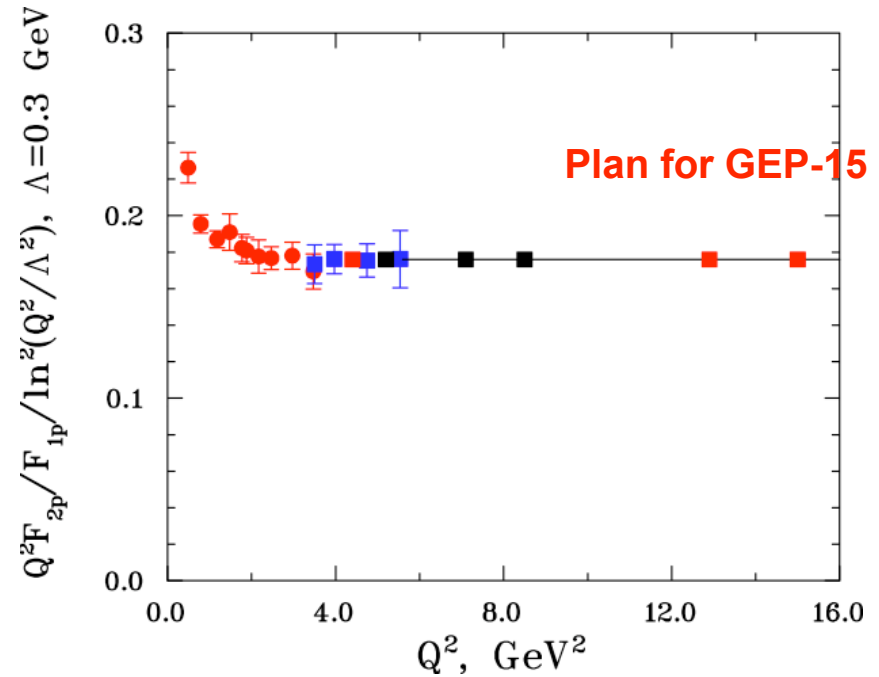
$$\Delta(\mu G_E^p/G_M^p) = \pm 0.10$$

approved by PAC32  
for 12 GeV program

# GEP-15: Projected accuracy



$$\Delta(\mu G_E^p / G_M^p) = \pm 0.10$$



$\Delta(F_2/F_1)/(F_2/F_1)$  accuracy will be **3%**

compare to  $\frac{\ln^2(Q^2=10/\Lambda^2)}{\ln^2(Q^2=15/\Lambda^2)} = 0.85$

# GMP-18: New measurement of $G_M^p$

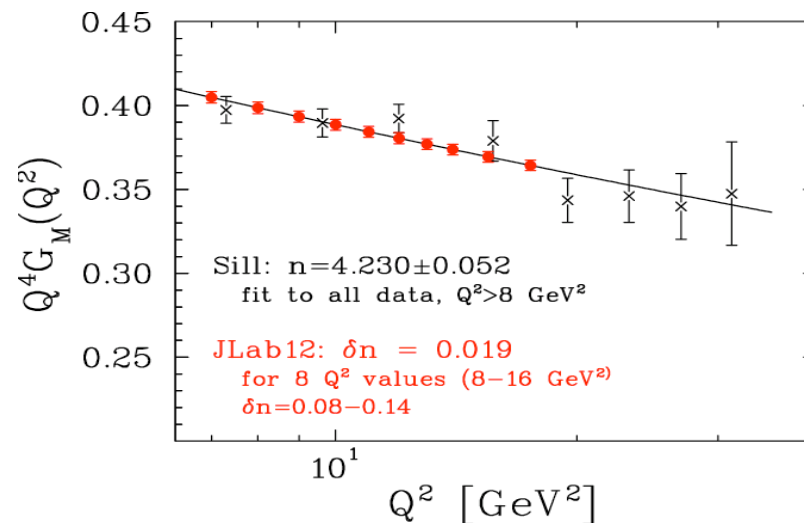
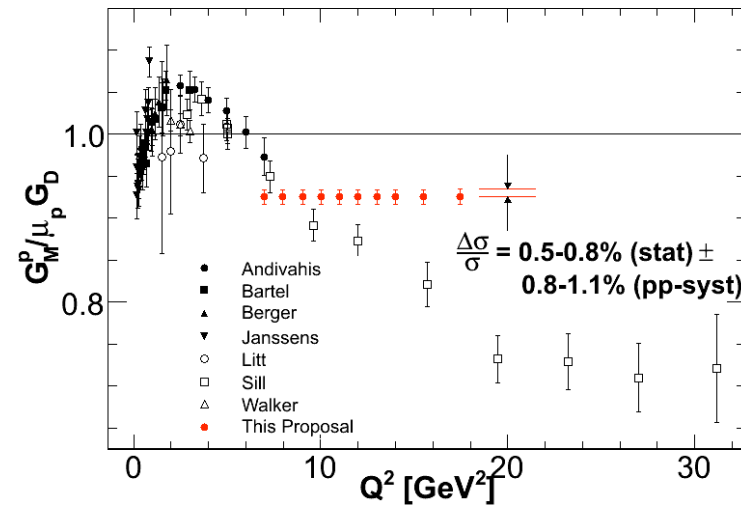
Moffit, Gilad, Arrington, BW

The cross section of  $H(e,e')p$ .

By using two existing Hall A High Resolution Spectrometers with several new ideas for improved control of systematic.

With 11 GeV beam in 31-day run.

approved by PAC32  
for 12 GeV program



# Summary

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- ❑ Experiment and theory have created an improved basis for the understanding of the nucleon
- ❑ Future experiments will provide precision FFs data for  $Q^2$  up to 7/14/15/18  $\text{GeV}^2$
- ❑ The GPD approach, as expected, sheds light on the nucleon structure
- ❑ Lattice QCD results for FFs are simmering

It is an exciting time for nucleon FF physics, when we know a lot about FFs and are searching for QCD-based interpretation

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backup  
slides  
after this



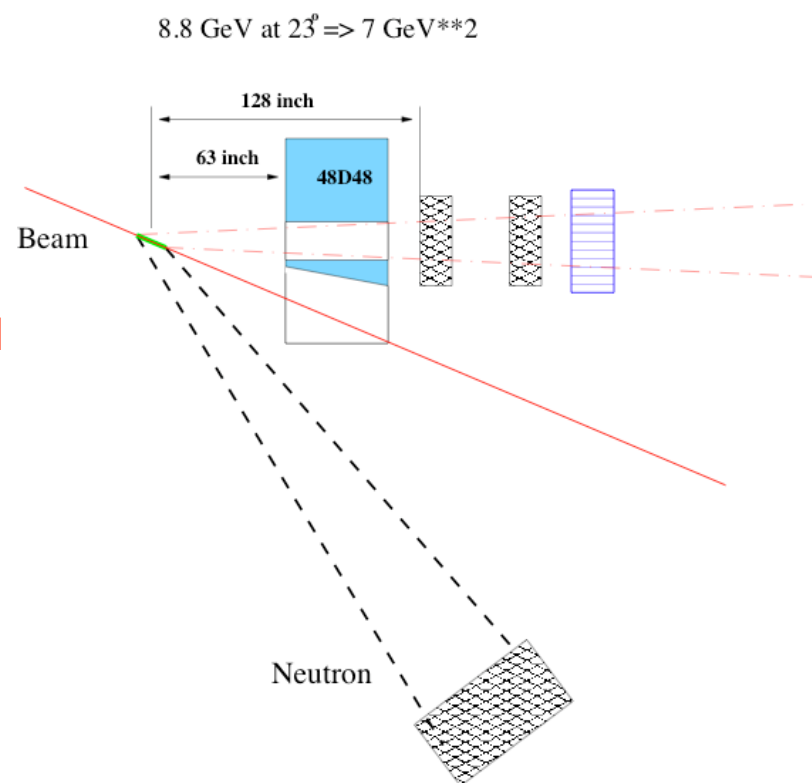
# Perspective: $G_E^n$ up to $7 \text{ GeV}^2$

A new plan for GEN-7 includes:

$${}^3\vec{H}e(\vec{e}, e'n)$$

- 8.8 GeV 85% polarized beam => triple FOM
- Resolution  $\sigma_p/p$  for electron - BNL magnet, GEM tracker => 3 times higher resolution
- He-3 cell in vacuum => lower background rate in neutron arm by a factor of 3
- Hybrid He-3 cell with narrow pumping laser line => 70% polarization

$G_E^n$  at  $7 \text{ GeV}^2$  with uncertainty 15% of Miller's value in 30-days run



# GEP-15: Proton Arm

- Magnet: 48D48 - 46 cm gap, 3 Tm field integral, 100 ton
- **solid angle is 35 msr for GEP**, could be  $\sim 70$  msr at larger angle
- **GEM chambers** for tracking with  $70 \mu\text{m}$  resolution
- **momentum resolution is 0.5% for 8.5 GeV/c proton**
- **angular resolution is 0.3 mrad**
- trigger threshold is 4 GeV from **hadron calorimeter**

