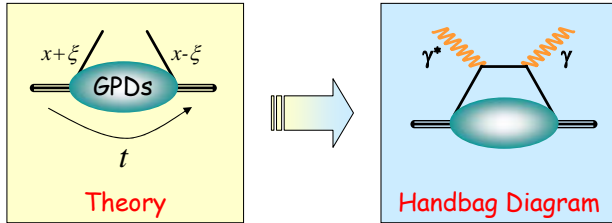




# GPDs from Theory to Experiment



## 2. The GPDs enter the DVCS amplitude as an integral over $x$ :

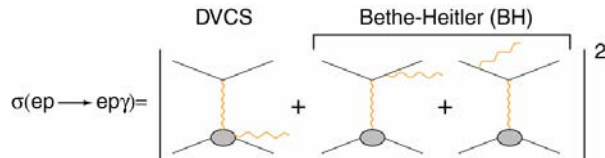
- GPDs appear in the **real part** through a PP integral over  $x$
- GPDs appear in the **imaginary part** but at the line  $x=\xi$

$$T^{DVCS} = \int_{-1}^{+1} \frac{GPD(x, \xi, t)}{x - \xi + i\varepsilon} dx + \dots$$

$$= P \int_{-1}^{+1} \frac{GPD(x, \xi, t)}{x - \xi} dx - i\pi GPD(x = \xi, \xi, t) + \dots$$

# Experimental observables linked to GPDs

Experimentally, DVCS is undistinguishable with Bethe-Heitler



However, we know FF at low  $t$  and **BH is fully calculable**

Using a polarized beam on an unpolarized target, 2 observables can be measured:

$$\frac{d^4\sigma}{dx_B dQ^2 dt d\varphi} \approx |T^{BH}|^2 + 2T^{BH} \cdot \text{Re}(T^{DVCS}) + |T^{DVCS}|^2$$

$$\frac{d^4\vec{\sigma} - d^4\overleftarrow{\sigma}}{dx_B dQ^2 dt d\varphi} \approx 2T^{BH} \cdot \text{Im}(T^{DVCS}) + \left[ |T^{DVCS \rightarrow}|^2 - |T^{DVCS \leftarrow}|^2 \right]$$

At low energy,  
 $|T^{DVCS}|^2$  supposed small

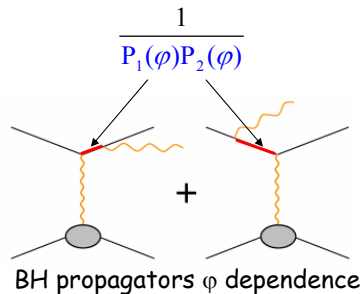
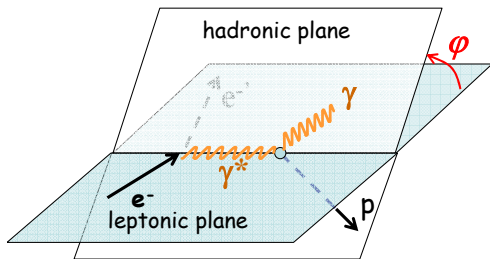
# Into the harmonic structure of DVCS

$$\frac{d^4\sigma}{dx_B dQ^2 dt d\varphi} = \frac{1}{P_1(\varphi)P_2(\varphi)} \Gamma_1(x_B, Q^2, t) \left\{ c_0^{BH} + c_1^{BH} \cos \varphi + c_2^{BH} \cos 2\varphi \right\} + \frac{1}{P_1(\varphi)P_2(\varphi)} \Gamma_2(x_B, Q^2, t) \left\{ c'_0 + c'_1 \cos \varphi + c'_2 \cos 2\varphi + c'_3 \cos 3\varphi \right\}$$

$$|T^{BH}|^2$$

$$\frac{d^4\vec{\sigma} - d^4\overleftarrow{\sigma}}{dx_B dQ^2 dt d\varphi} = \frac{\Gamma(x_B, Q^2, t)}{P_1(\varphi)P_2(\varphi)} \left\{ s'_1 \sin \varphi + s'_2 \sin 2\varphi \right\}$$

Interference term



# Deep Exclusive experiments

Published or Finalized 2001			Published or close to be		Analysis in progress	2008 – 2009	2010?	2013?
HERMES 27 GeV	CLAS 4-5 GeV	CLAS 5.75 GeV	Hall A 6 GeV	CLAS 6 GeV	HERMES	JLab	COMPASS	JLab@ 12GeV
<u>DVCS –</u> <u>BSA + BCA</u>	<b>DVCS</b> <b>BSA</b>	<b>DVCS</b> <b>DDVCS</b> <b>ADVCS</b> <b>D2VCS</b>  <b>Polarized</b> <b>DVCS</b>	<u>DVCS</u> <u>proton</u> <u>neutron</u>	<b>DVCS</b> <b>Proton</b>	<u>DVCS</u> <u>BSA+BCA</u>  <i>With recoil</i> <i>detector</i>	<u>Hall A</u> <u>DVCS<sup>2</sup></u>  CLAS BSA+TSA	<b>DVCS</b> <b><math>\sigma</math>+BCA</b>  <i>With recoil</i> <i>detector</i>	EVERYTHING, with more statistics than ever before
+ <u>nuclei</u> <u>BSA</u> <u>d-BCA</u>								
+ <u>Polarized</u> <u>target</u> <u>DVCS</u>								
<u><math>ep \rightarrow epp</math></u> <u><math>\sigma_L</math> + DSA</u>		<u><math>ep \rightarrow epp_L</math></u>	<u><math>ep \rightarrow ep\pi^0</math></u>	<u><math>ep \rightarrow ep\pi^0/\eta</math></u>				
		<u><math>ep \rightarrow ep\omega_L</math></u>						
		<u><math>ep \rightarrow ep\pi^0/\eta</math></u>						
		<u><math>ep \rightarrow en\pi^+</math></u>						
		<u><math>ep \rightarrow ep\Phi</math></u>						

More information for underlined items by: A. Camsonne, W. Dieter Nowak, A. Sandacz

## Special case of the asymmetry

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The asymmetry can be written as:

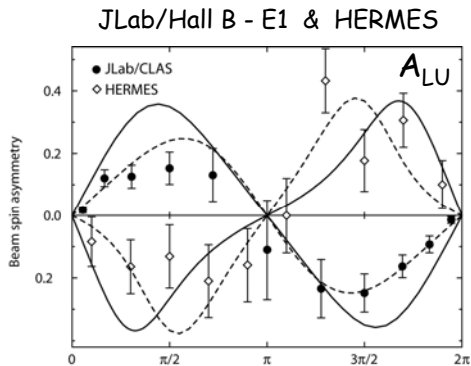
$$\frac{d^4 \vec{\sigma} - d^4 \overleftarrow{\sigma}}{d^4 \vec{\sigma} + d^4 \overleftarrow{\sigma}} = \Gamma_A(x_B, Q^2, t) \frac{s_1^I \sin \varphi + s_2^I \sin 2\varphi}{c_0^I + c_0^{BH} + (c_1^I + c_1^{BH}) \cos \varphi + \dots}$$

Pro: easier experimentally, smaller RC, smaller systematics

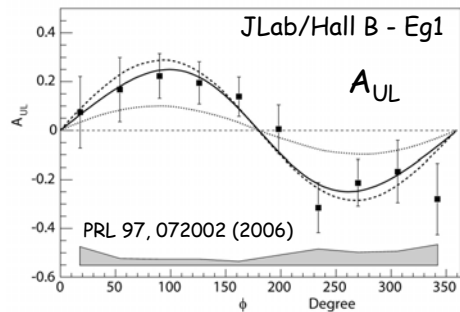
Con: direct extraction of GPDs is model- (or hypothesis-) dependent  
(denominator complicated and unknown)

It was naturally the first observable extracted from non-dedicated experiments...

# Published non-dedicated results on $A_{LU}$ and $A_{UL}$



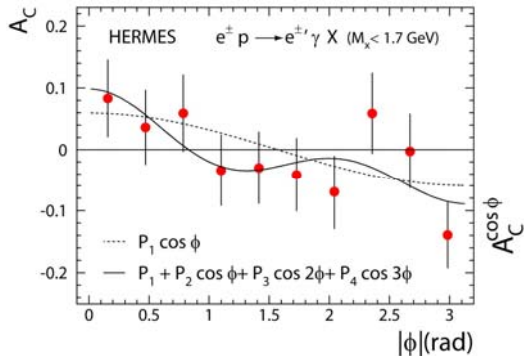
CLAS: PRL 87, 182002 (2001)  
HERMES: PRL 87, 182001 (2001)



Both results show, with a limited statistics, a  $\sin \phi$  behavior  
(necessary condition for handbag dominance)

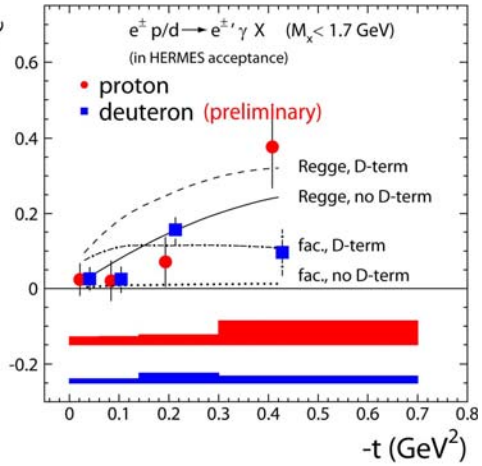
In the  $A_{LU}$  result, models (VGG) tend to over-estimate the data

# HERMES published (p) or unpublished (d) results on the BCA



Sensitive to the real part of the DVCS.BH interference

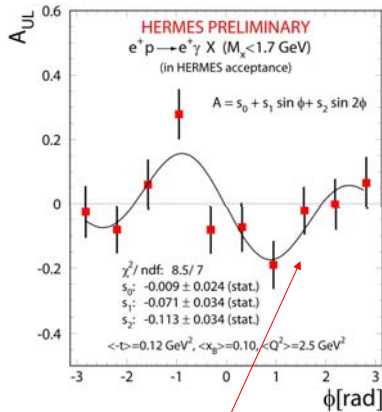
Low statistics but a cosine wave is observed.  
Still hard to compare to models





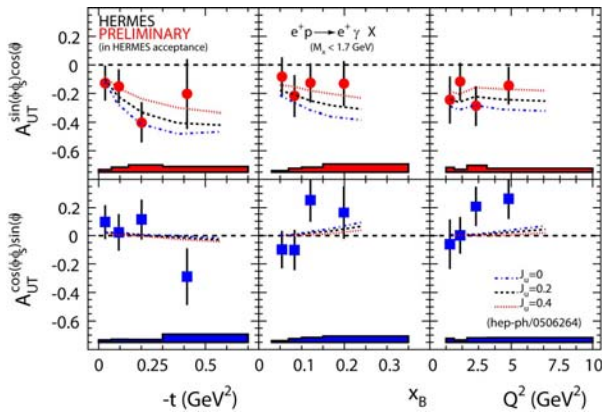
# HERMES preliminary results for $A_{UL}$ and $A_{UT}$

More in W. Dieter Nowak's talk



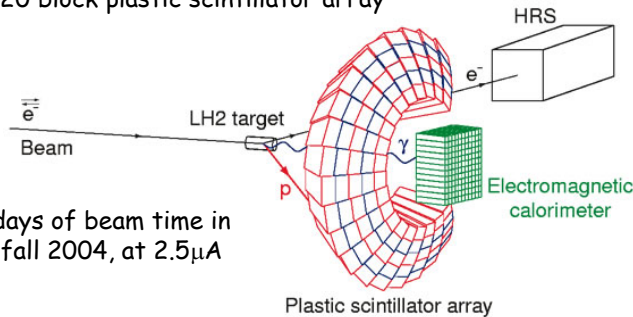
Strong  $\sin 2\phi$  is observed, unlike in the CLAS data

Model-dependent constraint on  $J_u - J_d$  (more on this later)



# E00-110 experimental setup and performances

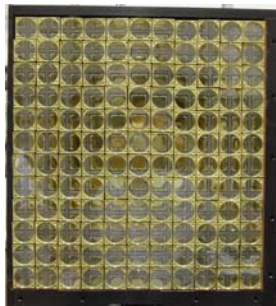
- 75% polarized 2.5 $\mu$ A electron beam
- 15cm LH2 target
- Left Hall A HRS with electron package
- 11x12 block PbF2 electromagnetic calorimeter
- 5x20 block plastic scintillator array



50 days of beam time in the fall 2004, at 2.5 $\mu$ A

$$\int Lu \cdot dt = 13294 \text{ fb}^{-1}$$

Kin	$Q^2$ (GeV $^2$ )	$x_B$	$\theta_{\gamma^*}$ (deg.)	$W$ (GeV)
1	1.5	0.36	22.3	1.9
2	1.9	0.36	18.3	2.0
3	2.3	0.36	14.8	2.2

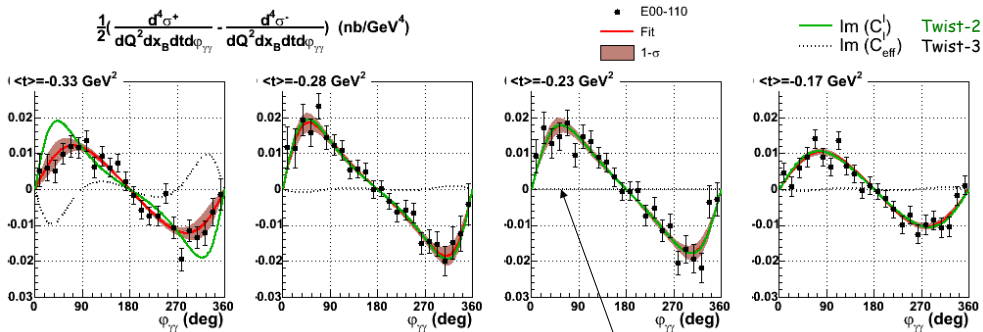


# Difference of cross-sections

PRL97, 262002 (2006)

$$\langle Q^2 \rangle = 2.3 \text{ GeV}^2$$

$$\langle x_B \rangle = 0.36$$

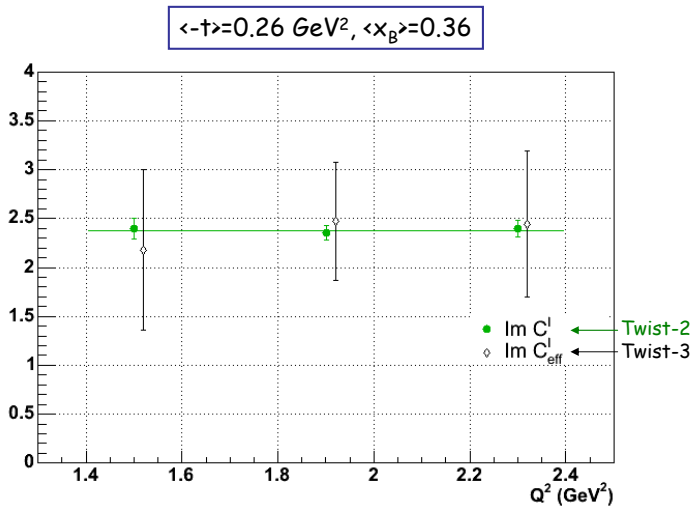


Corrected for real+virtual RC  
 Corrected for efficiency  
 Corrected for acceptance  
 Corrected for resolution effects  
 Checked elastic cross-section @ ~1%

Extracted Twist-3  
 contribution small!

New work by P. Guichon !

## $Q^2$ dependence and test of scaling



No  $Q^2$  dependence using BMK separation:  
strong indication for scaling behavior and handbag dominance

Twist 4+ contributions are smaller than 10%

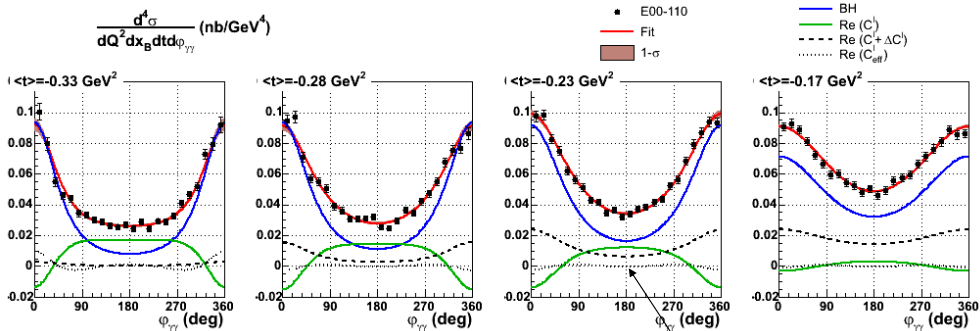
# Total cross-section

PRL97, 262002 (2006)

See also A. Camsonne's talk

$$\langle Q^2 \rangle = 2.3 \text{ GeV}^2$$

$$\langle x_B \rangle = 0.36$$



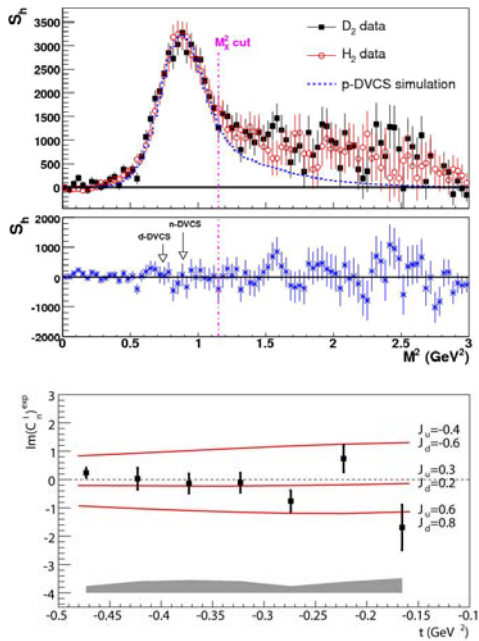
Corrected for real+virtual RC  
 Corrected for efficiency  
 Corrected for acceptance  
 Corrected for resolution effects

Extracted Twist-3  
 contribution small !

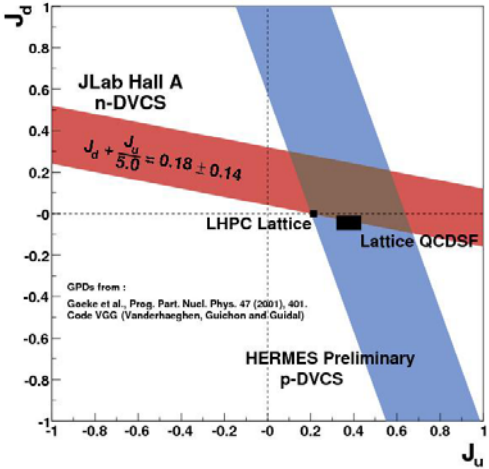
but impossible to disentangle DVCS<sup>2</sup>  
 from the interference term !

# DVCS on the neutron in JLab/Hall A: E03-106

LD<sub>2</sub> target  
 24000 fb-1  
 $x_B=0.36, Q^2=1.9 \text{ GeV}^2$

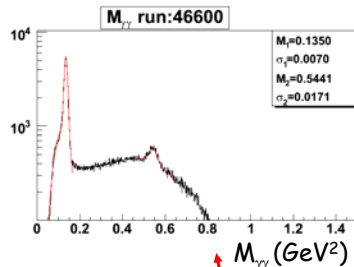


MODEL-DEPENDENT  
 J<sub>u</sub>-J<sub>d</sub> extraction

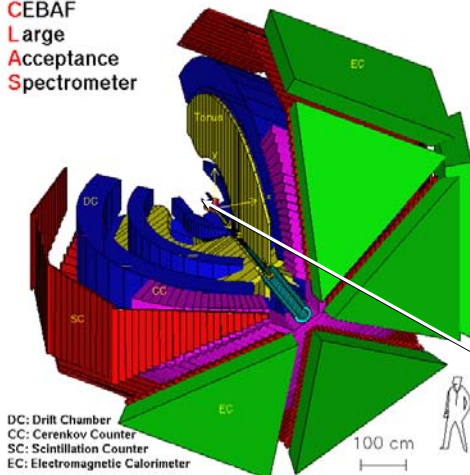


# E1-DVCS with CLAS : a dedicated DVCS experiment in Hall B

Beam energy:  $\sim 5.8$  GeV  
Beam Polarization: 75-85%  
Integ. Luminosity:  $45 \text{ fb}^{-1}$



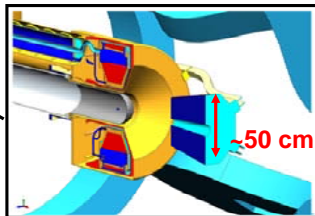
CEBAF  
Large  
Acceptance  
Spectrometer



DC: Drift Chamber  
CC: Cerenkov Counter  
SC: Scintillation Counter  
EC: Electromagnetic Calorimeter

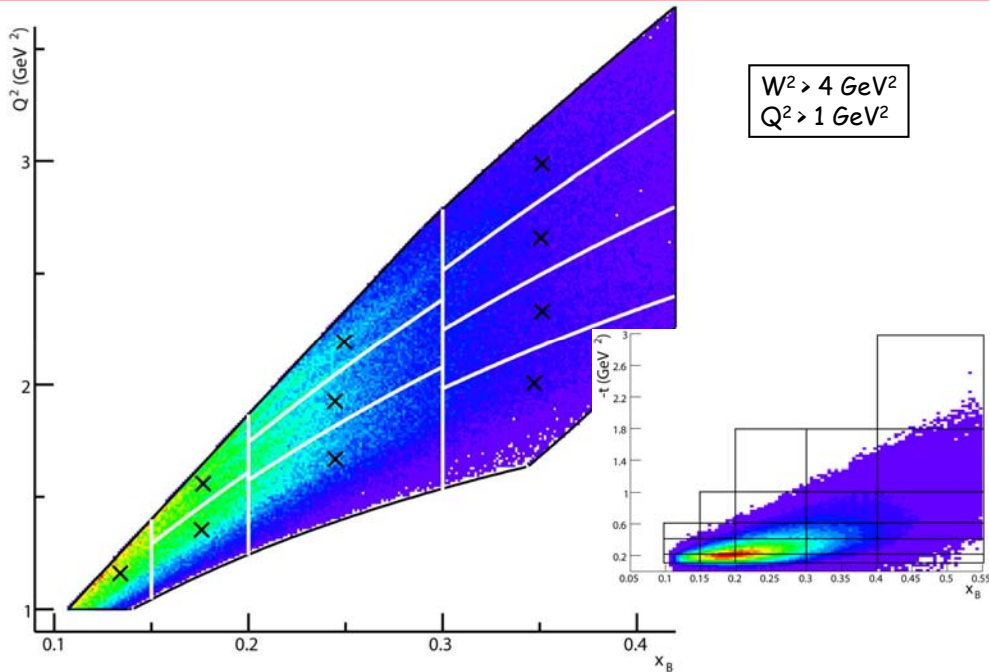
100 cm

Inner Calorimeter  
+ Moller shielding solenoid



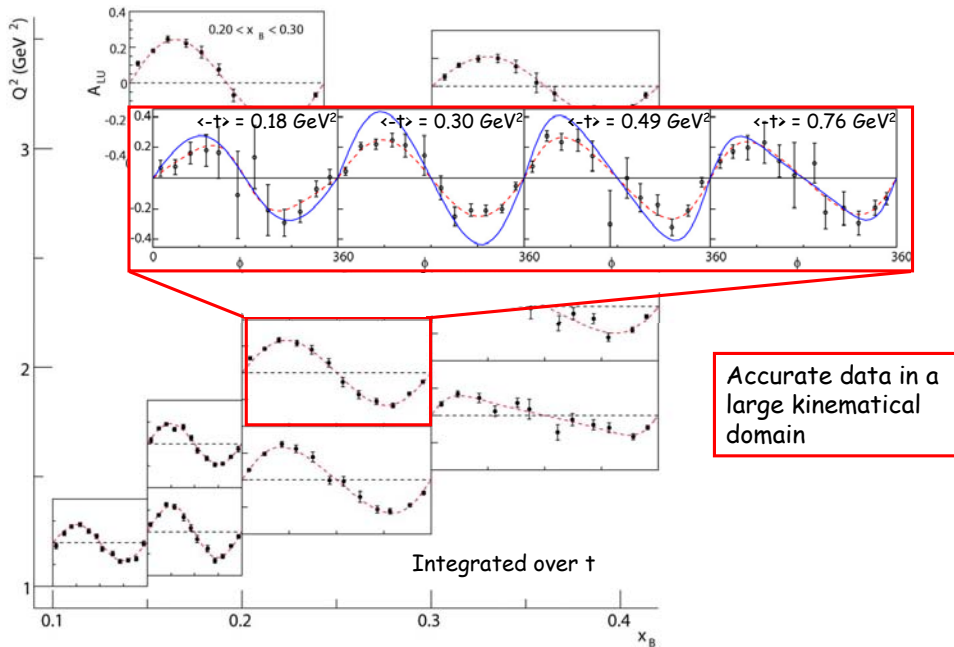
$\sim 50$  cm

# E1-DVCS kinematical coverage and binning

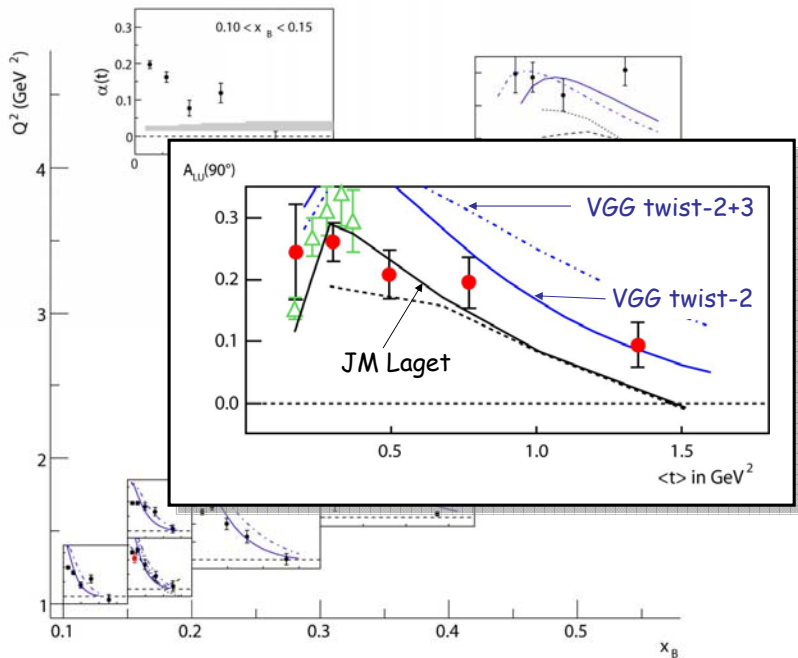




# E1-DVCS : Asymmetry as a function of $x_B$ and $Q^2$

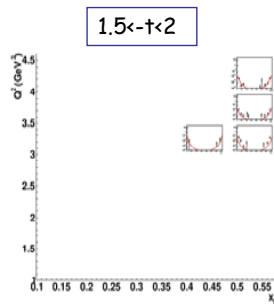
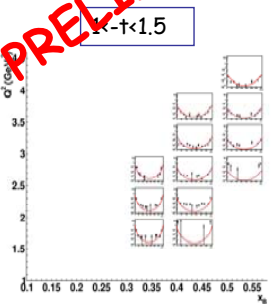
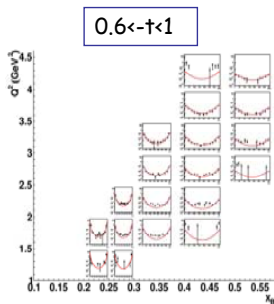
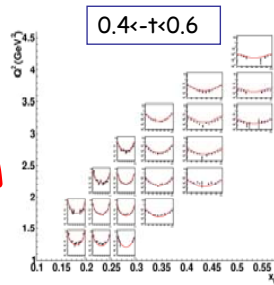
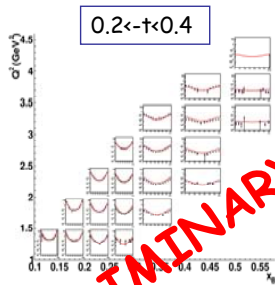
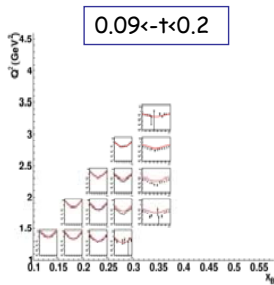


# E1-DVCS : $A_{LU}(90^\circ)$ as a function of $|t|$ + models



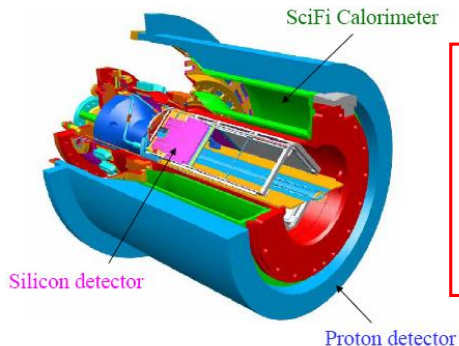
# E1-DVCS : Cross-sections over a wide kinematical range

PhD Thesis H.S. Jo

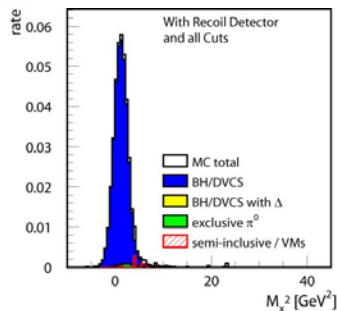
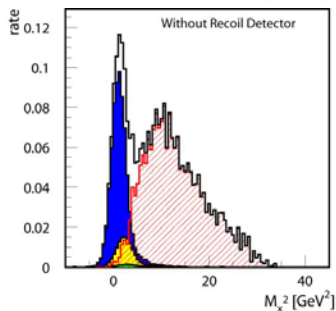


PRELIMINARY

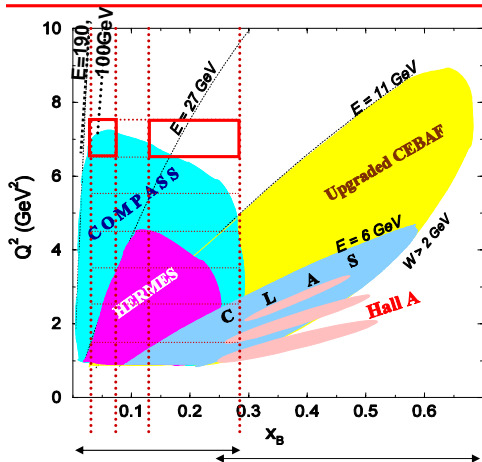
# The HERMES recoil detector: towards « true » exclusivity



- Inside a 1T superconducting solenoid
- Detection of the recoil proton: momentum measurement and PID (p/ $\pi$ )  
=> improves  $t$  resolution, enhances signal &  $\pi^0$  background rejection
- Collected statistics:
  - e<sup>-</sup> 2006: H<sub>2</sub> 5k DVCS, D<sub>2</sub> 1k DVCS
  - e<sup>-</sup>/<sub>+</sub> 2007: H<sub>2</sub> 42k DVCS, D<sub>2</sub> 10k DVCS



# COMPASS at the horizon 2010



## COMPASS at CERN-SPS

High energy muon beam  
100/190 GeV

$\mu^+$  or  $\mu^-$   
change once per day  
 $\text{polar}(\mu^+) = -0.80$   
 $\text{polar}(\mu^-) = +0.80$

$2.10^8 \mu$  per SPS cycle

valence quarks  
and sea quarks  
and gluons

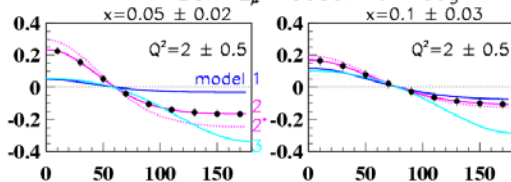
COMPASS  
2010?

valence quarks

JLab 12 GeV  
2013?

2 bins, for 6 months of data at 25% efficiency

BCA  $E_\mu = 100 \text{ GeV}$   $\vartheta = 1 \text{ deg}$



## Summary

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After 7 years of worldwide experimental efforts on DVCS:

-It is rather firmly established that leading twist contribution seems to dominate unlike in the meson case, and very similarly to DIS.

-GPD models such as *VGG* do not reproduce well most observables and it is now the ideal time to get the theory/phenomenology up to speed with the experimental side.

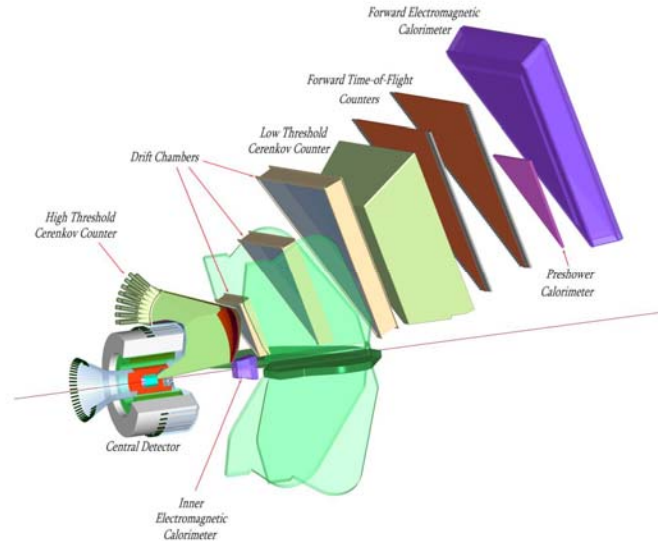
-It is only by a strong and common effort between theorists (modelists) and experimentalists that we will gain information about the GPDs themselves : disentangling GPDs from observables is a considerable challenge.

-The future 6 GeV and 12 GeV data from JLab, along with data coming from the analysis of recoil-detector-enabled data from HERMES and COMPASS will bring an amazing quantity of clean data in the next 8 years. We need to refine our analysis tools in preparation for this challenge !

# Experimental Setup and proposed experiments at 11 GeV

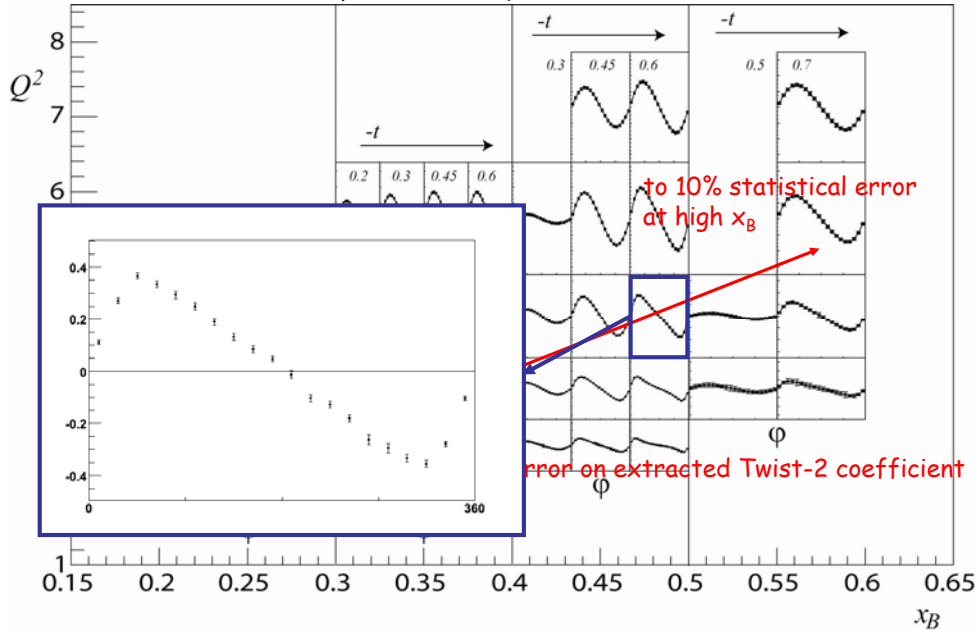
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Use of base CLAS12 equipment, including Inner Calorimeter (IC)



# Beam Spin Asymmetry

IC in standard position - 80 days -  $10^{35}$  Lum - VGG model





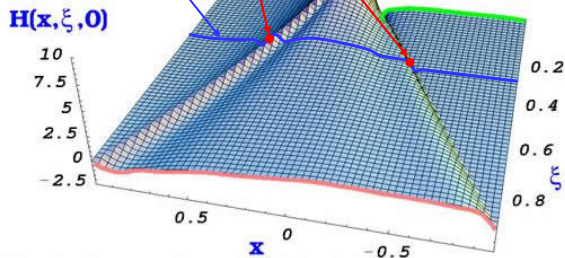
# Observables and their relationship to GPDs

$$T^{DVCS} = \int_{-1}^{+1} \frac{GPD(x, \xi, t)}{x - \xi + i\epsilon} dx + \dots$$

The **cross-section difference** accesses the imaginary part of DVCS and therefore GPDs at  $x = \xi$

$$= P \int_{-1}^{+1} \frac{GPD(x, \xi, t)}{x - \xi} dx - i\pi GPD(x = \xi, \xi, t) + \dots$$

The **total cross-section** accesses the real part of DVCS and therefore an integral of GPDs over  $x$



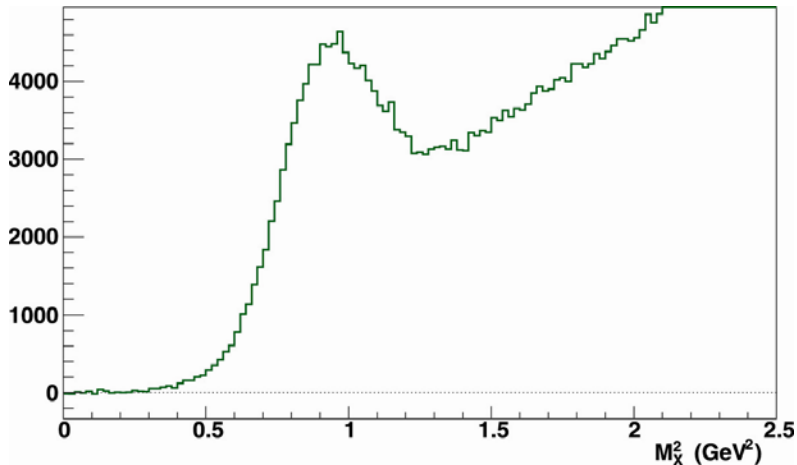
# Analysis - Looking for DVCS events

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HRS: Cerenkov, vertex, flat-acceptance cut with R-functions

Calo: 1 cluster in coincidence in the calorimeter above 1 GeV

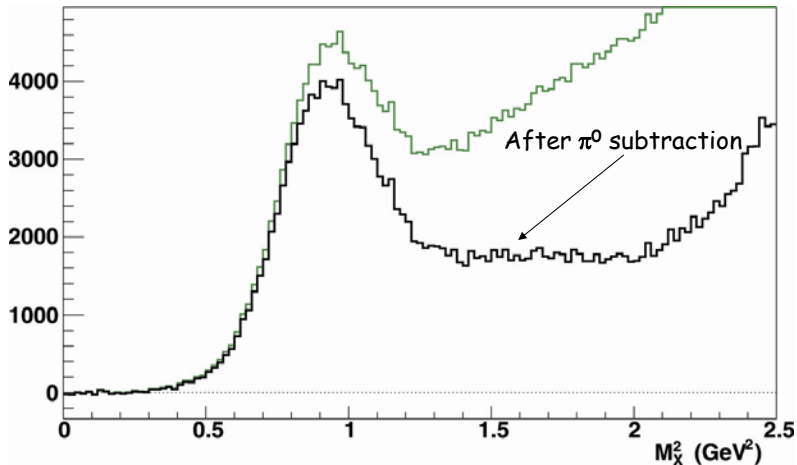
With both: subtract accidentals, build missing mass of  $(e,\gamma)$  system



# Analysis - $\pi^0$ subtraction effect on missing mass spectrum

---

Using  $\pi^0 \rightarrow 2\gamma$  events in the calorimeter, the  $\pi^0$  contribution is subtracted bin by bin



# Analysis - Exclusivity check using Proton Array and MC

---

Using Proton-Array, we compare the missing mass spectrum of the triple and double-coincidence events.

The missing mass spectrum using the Monte-Carlo gives the same position and width. Using the cut shown on the Fig., the contamination from inelastic channels is estimated to be under 3%.

